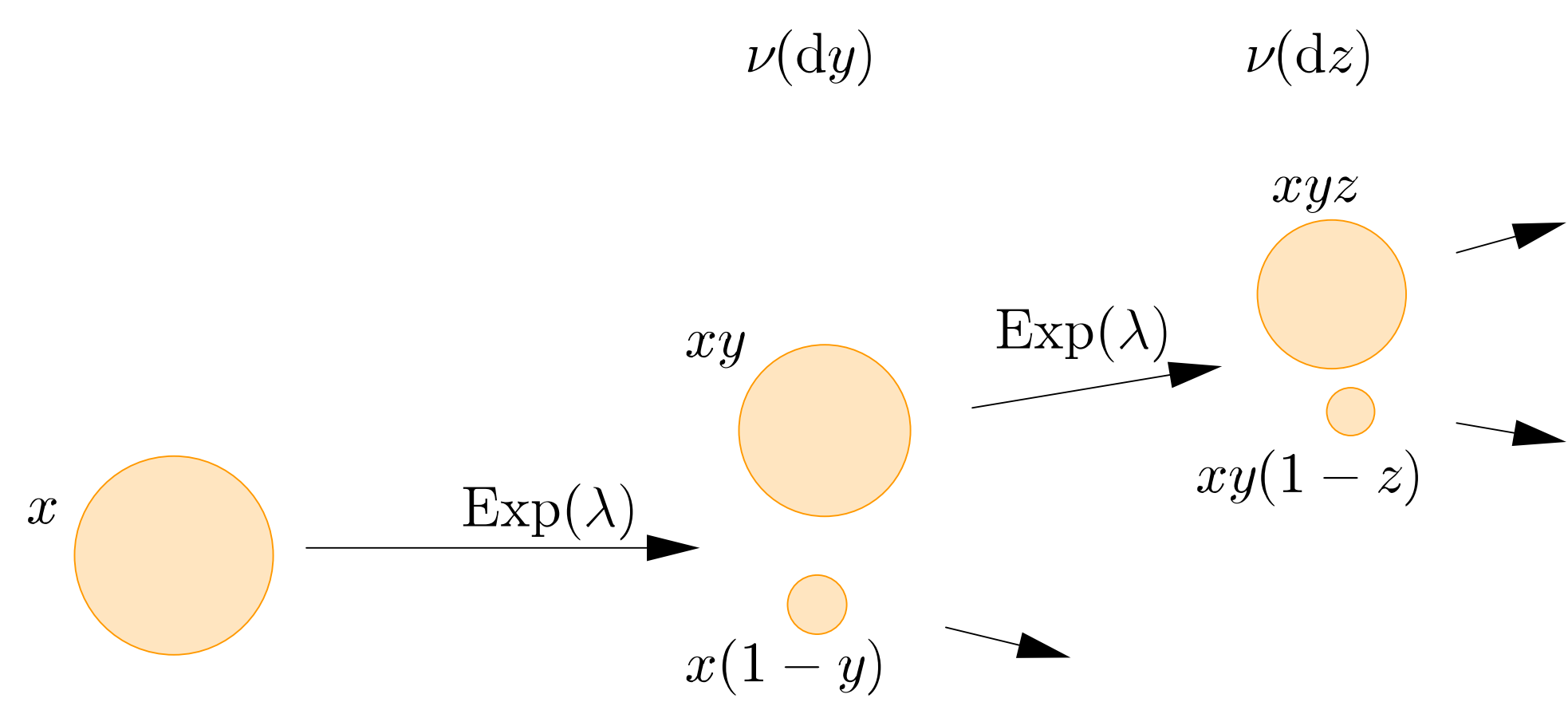


# MODELS FOR FRAGMENTATION WITH GROWTH

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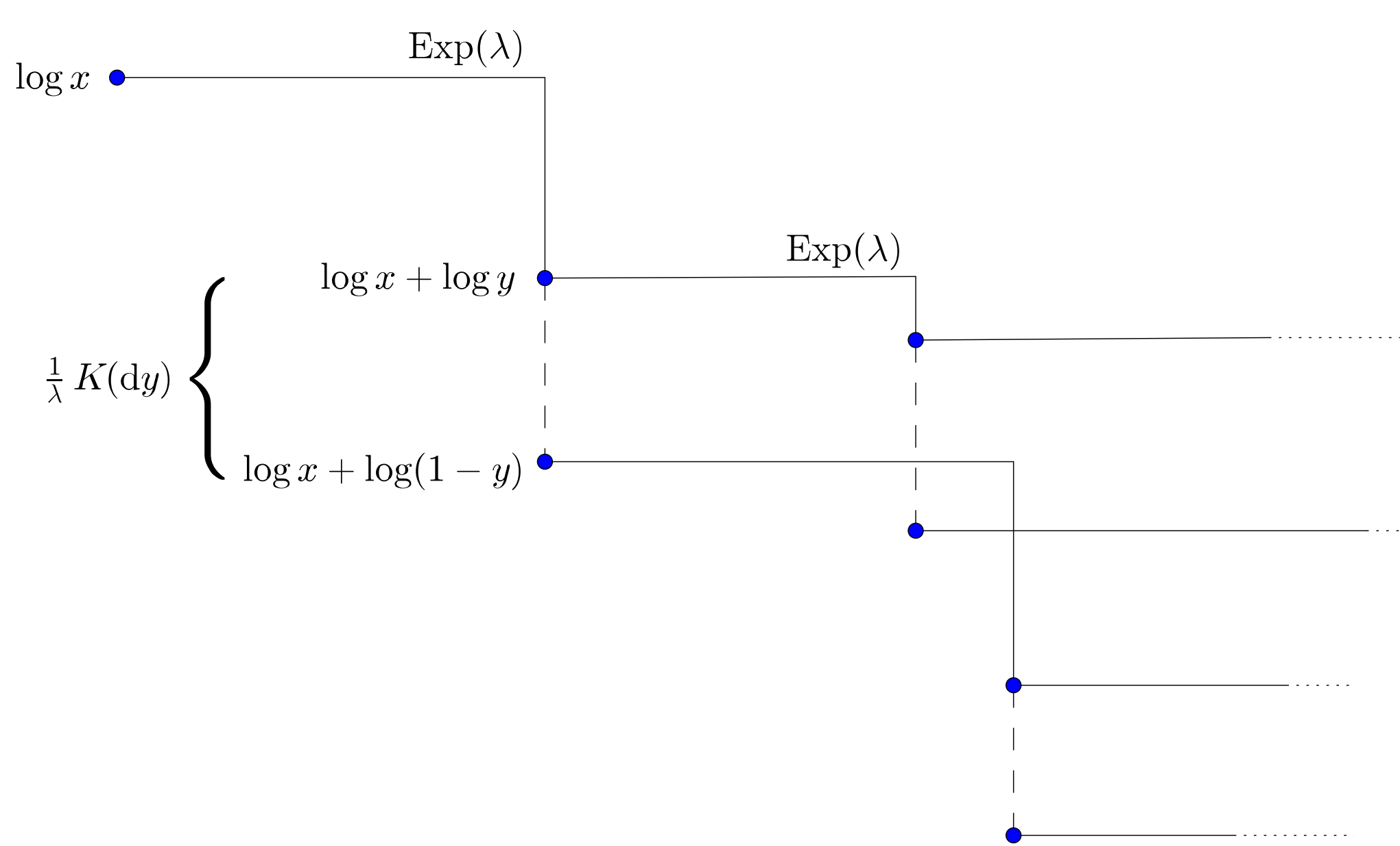
## A simple model of fragmentation



We begin with a single particle of size  $x$ . It waits for an exponential time of rate  $\lambda$ , and then splits into two daughter particles with sizes  $xy$  and  $x(1-y)$ , where  $y$  is chosen from some probability distribution  $\nu$  with support in  $[\frac{1}{2}, 1]$ . These act independently of each other and of the past, each waiting an exponential time and fragmenting according to the same rule.

## Evolution of the particle sizes

To study this model, let  $\mathcal{Z}(t) = \sum_{\text{particles } u} \delta_{\log(\text{size of } u)} \mathbb{1}_{\{u \text{ present at time } t\}}$ . We plot the evolution of these particle sizes; let  $K(dy) = \lambda \nu(dy)$ .

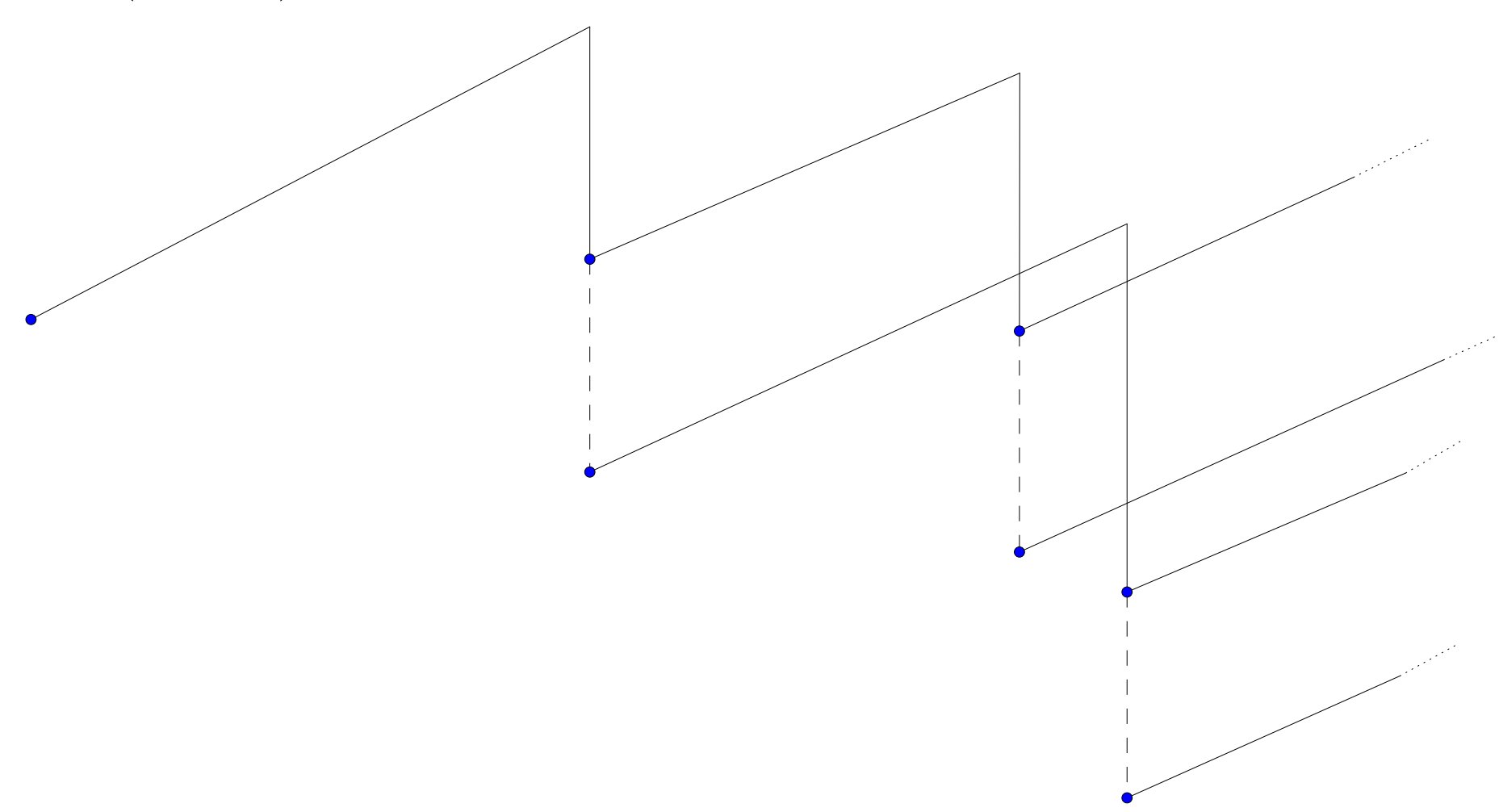


This graphic has a useful interpretation. By *identifying each particle with its largest child*, what we see above is a particle whose size remains constant for some exponential time, and then jumps by  $\log y$ . At the same time, a new particle is immigrated below at position  $\log(1-y)$  relative to the pre-jump position. Each particle then independently runs the same process.

We see that  $\mathcal{Z}$  is precisely a compound Poisson process, with Lévy measure  $\Pi = K \circ \log^{-1}$ , together with *immigration at its jumps*.

## A more elaborate model

This suggests a generalisation. Let  $X$  be a Lévy process with only downward jumps, and denote its Lévy measure by  $\Pi$ . A more general model of fragmentation,  $\mathcal{Z}$ , can be regarded as follows: we start with a single particle, which runs the process  $X$ ; and at every jump of  $X$  having size  $\log y$ , we immigrate an independent copy of  $\mathcal{Z}$  at relative position  $\log(1-y)$ .



One can show that such a process does exist, and is locally finite (each compact set contains only finitely many particles.)

Let  $\mathcal{Y} = \mathcal{Z} \circ \exp^{-1} = \sum_{\text{particles } u} \delta_{\text{size of } u} \mathbb{1}_{\{u \text{ present at time } t\}}$ . This is the *compensated fragmentation process*.

The measure  $K := \Pi \circ \exp^{-1}$  is called the *dislocation measure*, and it satisfies  $\int (1-y)^2 K(dy) < \infty$ .

## A simple fragmentation equation

If  $K$  is a finite measure, the following equation describes the mean behaviour of the ‘simple model’, left.

$$\partial_t \langle \mu_t, f \rangle = \left\langle \mu_t, x \mapsto \int_{[\frac{1}{2}, 1]} \{f(xy) + f(x(1-y)) - f(x)\} K(dy) \right\rangle, \quad f \in C_c^\infty(0, \infty),$$

$$\mu_0 = \delta_1$$

This equation still makes sense when  $K$  is infinite but satisfies  $\int (1-y) K(dy) < \infty$ , and has been well-studied in the context of exchangeable fragmentations.

## The growth-fragmentation equation

We study the following generalisation, which is well-defined for all  $K$  such that  $\int (1-y)^2 K(dy) < \infty$ .

$$\partial_t \langle \mu_t, f \rangle = \left\langle \mu_t, x \mapsto axf'(x) + \int_{[\frac{1}{2}, 1]} \{f(xy) + f(x(1-y)) - f(x) + (1-y)xf'(x)\} K(dy) \right\rangle \quad (0\text{-GF})$$

$$\mu_0 = \delta_1$$

There are two ‘growth’ terms now present. The first is the term involving  $a$ ; this represents a deterministic growth (or erosion, if  $a < 0$ ) of all particles simultaneously. The second is the extra term in the integral; this represents a sort of compensation for the infinite rate of ‘small fragmentations’, where the larger child is close in size to the parent, and the smaller child has size close to zero; it is required for the weakening of the integral condition on  $K$ .

## Theorem 1

Let

$$\kappa(q) = aq + \int_{[\frac{1}{2}, 1]} \{y^q + (1-y)^q - 1 + q(1-y)\} K(dy).$$

The unique solution of (0-GF) is given by

$$\langle \mu_t, f \rangle = \mathbb{E}_{\delta_1} \left[ \sum_u f(\mathcal{Y}_u(t)) \right] = e^{t\kappa(\omega)} \mathbb{E}_1[\xi(t)^{-\omega} f(\xi(t))],$$

where  $\mathcal{Y}$  is the compensated fragmentation process described on the left, and  $\xi$  is the exponential of a Lévy process with only downward jumps, having Laplace exponent  $\kappa(\cdot + \omega) - \kappa(\omega)$ ; the parameter  $\omega \in \text{dom } \kappa$  can be freely chosen.

This theorem corresponds exactly with what was already shown by Haas for the case  $\int (1-y) K(dy) < \infty$ , and the methods are very similar. However, there are some surprises in store for the self-similar equation, which we now describe.

## The self-similar growth-fragmentation equation

Let  $\alpha \in \mathbb{R}$ . The self-similar version of (0-GF) is as follows:

$$\partial_t \langle \mu_t, f \rangle = \left\langle \mu_t, x \mapsto x^\alpha \left[ axf'(x) + \int_{[\frac{1}{2}, 1]} \{f(xy) + f(x(1-y)) - f(x) + (1-y)xf'(x)\} K(dy) \right] \right\rangle. \quad (\alpha\text{-GF})$$

The presence of the term  $x^\alpha$  modifies the overall speed of evolution of a particle of size  $x$ . Thus, if  $\alpha < 0$ , then small particles (with size close to zero) grow and fragment rapidly, while larger particles have their behaviour slowed down; the opposite holds when  $\alpha > 0$ .

## Theorem 2

We assume that certain *Malthusian hypotheses* are satisfied. This means that there exists some  $\omega \in \text{dom } \kappa$  such that  $\kappa(\omega) = 0$ , and either  $\alpha < 0$  and  $\kappa'(\omega) > 0$ , or else  $\alpha > 0$  and  $\kappa'(\omega) < 0$ . Under these conditions, we have the following results.

- There exists a solution  $(\mu_t)$  to  $(\alpha\text{-GF})$ , such that  $\langle \mu_t, x \mapsto x^\omega \rangle \equiv 1$  and  $\mu_0 = \delta_1$ .  
– It is given by  $\langle \mu_t, f \rangle = \mathbb{E}_1[X(t)^{-\omega} f(X_t)]$ .
- There exists *another* solution  $(\gamma_t)$ , such that  $\langle \gamma_t, x \mapsto x^\omega \rangle \equiv 1$  for  $t > 0$  but  $\gamma_0 = 0$ .  
– If  $\alpha < 0$ , it is given by  $\langle \gamma_t, f \rangle = \mathbb{E}_0[X(t)^{-\omega} f(X_t)]$ .  
– If  $\alpha > 0$ , it is given by  $\langle \gamma_t, f \rangle = \mathbb{E}_{+\infty}[X(t)^{-\omega} f(X_t)]$ .

Here,  $X$  is a positive, self-similar Markov process with index of self-similarity  $-\alpha$ , which is driven (through the Lamperti transform) by a Lévy process with Laplace exponent  $\kappa(\cdot + \omega)$ ; the parameter  $\omega$  is the Malthusian exponent appearing above.

This indicates that the solutions can exhibit *spontaneous generation of mass*, and precludes uniqueness.

## Open questions

Some possible topics: existence of a minimal solution — study of the nonlinear semigroup and KPP equation — process variant of ‘starting from zero mass’ — study of biased mass functions ‘ $\langle \mu_t, x \mapsto x^q \rangle$ ’...