States of quantum systems and their liftings

Joachim Kupsch[1]
Fachbereich Physik, Universität Kaiserslautern
D-67653 Kaiserslautern, Germany

Oleg G. Smolyanov[2] and Nadejda A. Sidorova[3]
Faculty of Mechanics and Mathematics,
Moscow State University, 119899 Moscow, Russia

Abstract

Let H_1, H_2 be complex Hilbert spaces, H be their Hilbert tensor product and let tr_2 be the operator of taking the partial trace of trace class operators in H with respect to the space H_2. The operation tr_2 maps states in H (i.e. positive trace class operators in H with trace equal to one) into states in H_1. In this paper we give the full description of mappings that are linear right inverse to tr_2. More precisely, we prove that any affine mapping $F(W)$ of the convex set of states in H_1 into the states in H that is right inverse to tr_2 is given by $W \mapsto W \otimes D$ for some state D in H_2.

In addition we investigate a representation of the quantum mechanical state space by probability measures on the set of pure states and a representation – used in the theory of stochastic Schrödinger equations – by probability measures on the Hilbert space. We prove that there are no affine mappings from the state space of quantum mechanics into these spaces of probability measures.

1 INTRODUCTION

In quantum mechanics the states of a physical system are given by the statistical operators or density matrices in the Hilbert space associated to this system. The state of a subsystem is uniquely calculated as the reduced statistical operator by the partial trace. But it seems that the inverse problem: to define a linear mapping from the set of states of a subsystem to the set of states of an enlarged system such that the reduced state coincides with the original state, has not been studied systematically in the literature. In this article we want to investigate this lifting problem of states and the adjoint problem of reducing observables in some detail.

In the sequel all Hilbert spaces are assumed to be complex (and separable). For any Hilbert space we denote by $\mathcal{L}(H)$ the (complex) vector space of all linear bounded operators in H; by $\mathcal{L}^a(H)$ we denote the real vector subspace of $\mathcal{L}(H)$ consisting of all self-adjoint operators from $\mathcal{L}(H)$, by $\mathcal{L}^+(H)$ we denote the cone of positive operators within $\mathcal{L}(H)$ (and hence within $\mathcal{L}^a(H)$). The (complex) vector space of all trace class operators in H is denoted by $\mathcal{L}_1(H)$. In addition we use the following notations: $\mathcal{L}_1^+(H) = \mathcal{L}^+(H) \cap \mathcal{L}_1(H)$, $\mathcal{L}_1^a(H) = \mathcal{L}_1^+(H) \cap \mathcal{L}^a(H)$, and $\mathcal{D}(H)$ is the convex set of all operators

[1]e-mail: kupsch@physik.uni-kl.de
[2]e-mail: smolyan@mail.ru
[3]e-mail: nadja@sidorova.mccme.rssi.ru

from $\mathcal{L}_1^+(H)$ having trace equal to one. If H is the Hilbert space associated to a physical system, then the elements of $\mathcal{L}^a(H)$ represent the (bounded) observables of the system, the elements of $\mathcal{D}(H)$ represent (mixed and pure) states, and the closed subset $\mathcal{P}(H) \subset \mathcal{D}(H)$ of rank one projection operators represents the pure states.

If $\mathcal{S}$ and $\mathcal{E}$ are physical systems with Hilbert spaces H_S and H_E, then the Hilbert space of the composite system – denoted by $\mathcal{S}\times\mathcal{E}$ – of these systems is the Hilbert tensor product of Hilbert spaces H_S and H_E, i.e. $H = H_S \otimes H_E$. The scalar product in H is written as $\langle\ ,\ \rangle_H$; the corresponding notations are used for scalar products in H_S and H_E. Hence $\mathcal{S}$ is a subsystem of the quantum system $\mathcal{S}\times\mathcal{E}$, and the system $\mathcal{E}$ can be interpreted as an environment of H_S. For any state $W \in \mathcal{D}(H)$ of the total system $\mathcal{S}\times\mathcal{E}$ the state of the system $\mathcal{S}$ – called the reduced state – is given by the partial trace $\mathrm{tr}_{H_E} W \in \mathcal{D}(H_S)$. This partial trace is uniquely defined for all $W \in \mathcal{L}_1(H)$ as the operator $\mathrm{tr}_{H_E} W \in \mathcal{L}_1(H_S)$ which satisfies the identity $\langle \mathrm{tr}_{H_E} W x_1, x_2\rangle_{H_S} = \sum_j \langle W(x_1 \otimes e_j^E), (x_2 \otimes e_j^E)\rangle_H$ for an orthonormal basis $\left\{e_j^E\right\}$ of H_E and all $x_1, x_2 \in H_S$. The mapping $W \to \mathrm{tr}_{H_E} W$, $\mathcal{L}_1(H) \to \mathcal{L}_1(H_S)$, is obviously linear and continuous.

By the partial trace we can calculate the Schrödinger dynamics of the subsystem $\mathcal{S}$ – the so called reduced dynamics – from the Schrödinger dynamics of the whole system $\mathcal{S}\times\mathcal{E}$. But in general, this dynamics does not depend linearly on the initial state of the subsystem, see Remark 2. In order to obtain the linear dependence one has to find a linear solution for the lifting problem, which can be formulated as follows. For any state $W_S \in \mathcal{L}_1(H_S)$ to find a state $F(W_S) \in \mathcal{L}_1(H)$ such that $\mathrm{tr}_{H_E} F(W_S) = W_S$; such a mapping F_S is called the lifting.

The simplest solution of this problem is given by the mapping $F_D : \mathcal{L}_1(H_S) \longrightarrow \mathcal{L}_1(H)$, $W \mapsto W \otimes D$, where D is an element of $\mathcal{L}_1(H_E)$, which is usually called a reference state. This choice is well known from the theory of open systems, see e.g. [1],[2],[3].

The main theorem of the paper – Theorem 1 of the next section – implies that actually any linear lifting coincides with F_D, for some D.

Remark 1 *The vector space $\mathcal{L}^a(H)$ of bounded observables can be identified with the space of continuous affine linear functionals on the state space $\mathcal{D}(H)$ equipped with the topology induced by the trace norm $\|.\|_1$ of $\mathcal{L}_1(H) \supset \mathcal{D}(H)$, see e.g. [4]. Affine linearity means that such a functional $f : \mathcal{D}(H) \longrightarrow \mathbb{R}$ respects the mixing property: $f(\alpha W_1 + \beta W_2) = \alpha f(W_1) + \beta f(W_2)$ for $0 \le \alpha, \beta \le 1$ with $\alpha + \beta = 1$, and $W_1, W_2 \in \mathcal{D}(H)$. In fact, any such a functional can be uniquely extended to a continuous $\mathbb{C}$-linear functional $\bar{f} : \mathcal{L}_1(H) \longrightarrow \mathbb{C}$, see e.g. [5]. Since $\mathcal{L}(H)$ is dual to $\mathcal{L}_1(H)$, with the duality pairing*

$$\mathcal{L}(H) \times \mathcal{L}_1(H) \longrightarrow \mathbb{C} : (A, W) \mapsto \langle A, W\rangle \equiv \mathrm{tr}_H AW, \tag{1}$$

there exists $A_f \in \mathcal{L}(H)$ such that for any $W \in \mathcal{L}_1(H)$ the identity $\overline{f}(W) = \mathrm{tr}_H W A_f$ is true.

On the other side, according to Gleason's theorem [6], the state space $\mathcal{D}(H)$ can be identified with the set of linear functionals $\omega : \mathcal{L}(H) \longrightarrow \mathbb{C}$ having the following properties:

(1) if $A \in \mathcal{L}^+(H)$ then $\omega(A) \ge 0$;

(2) $\omega(Id) = 1$;

(3) $\omega(\sum_j P_j) = \sum_j \omega(P_j)$ for any finite or countable family of mutually orthogonal projectors.

For any ω which satisfies these constraints there exists an element $W_\omega \in \mathcal{D}(H)$ such that $\omega(A) = \mathrm{tr}_H W_\omega A$ is true for all $A \in \mathcal{L}(H)$. The natural norm of the state space is $\sup_{\|A\|=1} |\omega(A)|$ which coincides with the trace norm of W_ω.

Remark 2 *The time evolution of a composite system with Hilbert space $H = H_S \otimes H_E$ in the Schrödinger picture is given by a family Φ_t, $t \in \mathbb{R}$ or $t \in \mathbb{R}_+$, of continuous affine linear mappings $\Phi_t : \mathcal{D}(H) \longrightarrow \mathcal{D}(H)$. We normalize these evolutions by $\Phi_0(W) = W$. The affine linear mappings Φ_t can be extended to $\mathbb{C}$-linear mappings on $\mathcal{L}_1(H)$, again denoted by Φ_t. In the usual case of a Hamiltonian (unitary) dynamics we have $\Phi_t(W) = U(t)WU^+(t)$ with the unitary group $U(t)$ on H generated by the Hamiltonian. But more general evolutions like semigroups are admitted in the sequel. The mappings Φ_t have unique extensions to continuous $\mathbb{C}$-linear mappings $\overline{\Phi}_t$ of $\mathcal{L}_1(H)$ into $\mathcal{L}_1(H)$. The duality (1) then allows to determine the Heisenberg evolution, a family Ψ_t of continuous linear operators on $\mathcal{L}(H)$. Any Schrödinger evolution Φ_t on $\mathcal{D}(H_S \otimes H_E)$ induces a unique time evolution $\rho_t = tr_{H_E}\Phi_t(W)$ of the system H_S. In order to obtain a linear dependence on the initial state $\rho = \rho_{t=0}$ we need an affine linear mapping F of $\mathcal{D}(H_S)$ into $\mathcal{D}(H_S \otimes H_E)$. Then the mapping $\rho \mapsto W = F(\rho) \mapsto \rho_t = tr_{H_E}\Phi_t(W)$ is a linear time evolution on $\mathcal{D}(H_S)$. This time evolution has the correct initial condition $\rho_{t=0} = \rho$ if F satisfies the constraint $\mathrm{tr}_{H_E} F(\rho) = \rho$. The Heisenberg dynamics of the system then follows from the duality (1) applied to $\mathcal{L}(H_S)$ and $\mathcal{L}_1(H_S)$.*

The paper is organized a follows. In Sec. 2 we prove the main result of the paper – Theorem 1 – describing all linear liftings. In Sec. 3 we consider a theorem – Theorem 2 – that is in a sense dual to Theorem 1 and describes a reduction of observables of the system H to observables of the system H_S.

In the final Sec. 4 we consider the case of a classical state space, i. e. a space of probability measures, and the representation of the quantum mechanical state space $\mathcal{D}(H)$ by probability measures either on the set of pure states – the Choquet representation – or on the Hilbert space – a representation used in the theory of stochastic Schrödinger equations. The space $\mathcal{D}(H)$ is a convex set with the closed set $\mathcal{P}(H)$ of pure states as extremal points. Any $W \in \mathcal{D}(H)$ can be represented by an integral over the pure states $W = \int_{\mathcal{P}(H)} \mu(dP)\, P$, where $\mu(dP)$ is a probability measure on $\mathcal{P}(H)$. Since this representation has been derived by Choquet for general convex sets, see e.g. [7], we denote the (non-unique) measure $\mu(dP)$ as Choquet measure of W. In Theorem 3 we prove that there does not exist a linear mapping γ from the space $\mathcal{D}(H)$ into the set of probability measures on the set $\mathcal{P}(H)$ such that the measure $\gamma(W)$ is the Choquet measure of the state $W \in \mathcal{D}(H)$. This theorem is in fact a consequence of Theorem 1. In Sec. 4 we deduce Theorem 3 from the structural difference between the classical and the quantum mechanical state spaces. Both these spaces are convex sets. But the classical state space is a simplex whereas $\mathcal{D}(H)$ not, see e.g. [8]. Finally we investigate the representation of the state space by probability measures on the Hilbert space. Also in this case the structural difference between the quantum mechanical state space and the space of probability measures does not allow an affine linear mapping from $\mathcal{D}(H)$ into the measure space.

2 LINEAR LIFTINGS

The main result of the paper is the following theorem.

Theorem 1 *Let $F : \mathcal{D}(H_S) \longrightarrow \mathcal{D}(H_S \otimes H_E)$ be an affine linear mapping such that $\mathrm{tr}_{H_E} F(\rho) = \rho$ for all $\rho \in \mathcal{D}(H_S)$. Then there exists an element $\rho_E \in \mathcal{D}(H_E)$ such that $F(\rho) = \rho \otimes \rho_E$.*

Proof The mapping F can be extended (uniquely) to the $\mathbb{C}$-linear mapping of $\mathcal{L}_1(H_S)$ into $\mathcal{L}_1(H_S \otimes H_E)$ that we shall denote by the same symbol. This extension has the following properties:

$$F(\mathcal{L}_1^+(H_S)) \subset \mathcal{L}_1^+(H_S \otimes H_E), \tag{2}$$

$$F(\mathcal{L}_1^a(H_S)) \subset \mathcal{L}_1^a(H_S \otimes H_E); \tag{3}$$

we shall use these properties later.

Let $\{e_i, i \in \mathbb{N}\}$ (respectively, $\{f_j, j \in \mathbb{N}\}$) be an orthonormal basis in H_S (respectively, in H_E). Without loss of generality we assume H_S and H_E to be infinite-dimensional. Then $H = \mathrm{span}\,\{e_i \otimes f_j,\ i \in \mathbb{N}, j \in \mathbb{N}\}$. We realize $\mathcal{L}_1(H_S)$ as a vector space of complex valued functions on $\mathbb{N}^2$: $g_{ij} \in \mathbb{C} : i \in \mathbb{N}, j \in \mathbb{N}$. Analogously, we realize $\mathcal{L}_1(H) = \mathcal{L}_1(H_S \otimes H_E)$ as a vector space of complex valued functions on $\mathbb{N}^4$: $F(g)_{ij}^{kl}, i, j, k, l \in \mathbb{N}$. We say that $b_{ij}, i, j \in \mathbb{N}$ is a (k, l)-component of $F \in \mathcal{L}_1(H)$ and denote it by $(F)^{kl}$ if $F_{ij}^{kl} = b_{ij}$ for all $i, j \in \mathbb{N}$. We say that F has only the components of some type if all components of other type are equal to zero. Let us note that

$$\mathrm{tr}_{H_E} F(g) = g \Leftrightarrow \sum_{i=1}^{\infty} F(g)_{ii}^{kl} = g_{kl},\ \forall k, l \in N. \tag{4}$$

Consider the following basis $\{g^{kl}, k \leq l, g^{kl*}, k < l\}$ in $\mathcal{L}_1(H_S)$:

$$g_{ij}^{kl} = \begin{cases} 1 & \text{if } (i,j) \in \{(k,k),(k,l),(l,k),(l,l)\}, \\ 0 & \text{otherwise} \end{cases}, \quad g_{ij}^{kl*} = \begin{cases} 1 & \text{if } (i,j) = (k,k), \\ i & \text{if } (i,j) = (k,l), \\ -i & \text{if } (i,j) = (l,k), \\ 1 & \text{if } (i,j) = (l,l), \\ 0 & \text{otherwise.} \end{cases}$$

Firstly, all g^{kl} and g^{kl*} are positive operators, therefore $F(g^{kl})$ and $F(g^{kl*})$ are also positive and hence $F(g^{kl})_{ii}^{mm} \geq 0$ and $F(g^{kl*})_{ii}^{mm} \geq 0$ for all $i, m \in \mathbb{N}$. Due to (4) $\sum_{i=1}^{\infty} F(g^{kl})_{ii}^{mm} = g_{mm}^{kl} = 0$ for $m \neq k, m \neq l$ and $\sum_{i=1}^{\infty} F(g^{kl*})_{ii}^{mm} = g_{mm}^{kl*} = 0$ for $m \neq k, m \neq l$. From this follows that

$$F(g^{kl})_{ii}^{mm} = 0, m \neq k, m \neq l, i \in \mathbb{N}, \quad F(g^{kl*})_{ii}^{mm} = 0, m \neq k, m \neq l, i \in \mathbb{N}. \tag{5}$$

Secondly, all g^{kl} and g^{kl*} are self-adjoint, therefore $F(g^{kl})$ and $F(g^{kl*})$ are also self-adjoint and hence

$$F(g^{kl})_{ji}^{nm} = \overline{F(g^{kl})_{ij}^{mn}} \text{ and } F(g^{kl*})_{ji}^{nm} = \overline{F(g^{kl*})_{ij}^{mn}} \text{ for all } i, j, k, l, m, n \in \mathbb{N}. \tag{6}$$

$\square$

The further proof is organized as follows. First, we show that $F(g^{kk})$ has only (k,k)-component (Step 1), $F(g^{kl}), k<l$ (resp., $F(g^{kl*}), k<l$) has only (k,k), (k,l), (l,k), (l,l)-components (Step 2). Furthermore, we prove that non-zero components of $F(g^{kl})$ are equal (Step 3) and that non-zero components of $F(g^{kl*})$ satisfy $(F)^{kk} = -i(F)^{kl} = i(F)^{lk} = (F)^{ll}$ (Step 4). Finally, we denote elements of the only non-zero component of $F(g^{11})$ by a_{ij} and show that any non-zero component of $F(g^{kl})$ is equal to a_{ij} (Step 5) and that the non-zero components of $F(g^{kl*})$ satisfy $(F)^{kk} = -i(F)^{kl} = i(F)^{lk} = (F)^{ll} = a_{ij}$ (Step 6), which completes the proof.

In the proof we shall also use the following (obvious) lemma.

Lemma 1 *Let $a \geq 0, b \geq 0, c \geq 0$ be real numbers; then*

$$\left\{(t,p) \in \mathbb{R}^2 : (1+t)(1+p) \geq 1, t+1 \geq 0\right\} \subset$$
$$\left\{(t,p) \in \mathbb{R}^2 : (b+at)(b+cp) \geq b^2, b+at \geq 0\right\}$$
$$\Leftrightarrow a = c \leq b.$$

Proof of Theorem 1 (continued)
Step 1. Consider $g^{kk} \in \mathcal{L}_1(H_S), k \in \mathbb{N}$ and restrict $F(g^{kk})$ to the space $\langle e_m \otimes f_i, e_n \otimes f_j \rangle$, where $(m,n) \neq (k,k)$. In this basis $F(g^{kk})$ has the form

$$\begin{pmatrix} F(g^{kk})^{mm}_{ii} & F(g^{kk})^{nm}_{ji} \\ F(g^{kk})^{mn}_{ij} & F(g^{kk})^{nn}_{jj} \end{pmatrix}.$$

Because either $m \neq k$ or $n \neq k$ we have due to (5) that either $F(g^{kk})^{mm}_{ii} = 0$ or $F(g^{kk})^{nn}_{jj} = 0$. $F(g^{kk})$ is positive, hence $F(g^{kk})^{mm}_{ii} F(g^{kk})^{nn}_{jj} - F(g^{kk})^{nm}_{ji} F(g^{kk})^{mn}_{ij} \geq 0$. Combining this conditions together with (6) we get

$$F(g^{kk})^{mn}_{ij} = 0 \; \forall i,j,m,n \in \mathbb{N}, (m,n) \neq (k,k),$$

i.e. $F(g^{kk})$ has only (k,k)-component.

Step 2. Consider $g^{kl} \in \mathcal{L}_1(H_S), k,l \in \mathbb{N}, k<l$ and restrict $F(g^{kl})$ to the subspace $\langle e_m \otimes f_i, e_n \otimes f_j \rangle$, where (m,n) does not take values (k,k), (k,l), (l,k), (l,l). In this basis $F(g^{kl})$ has the form

$$\begin{pmatrix} F(g^{kl})^{mm}_{ii} & F(g^{kl})^{nm}_{ji} \\ F(g^{kl})^{mn}_{ij} & F(g^{kl})^{nn}_{jj} \end{pmatrix}.$$

Due to the conditions on (m,n) it follows from (5) that either $F(g^{kl})^{mm}_{ii} = 0$ or $F(g^{kl})^{nn}_{jj} = 0$. $F(g^{kl})$ is positive, hence $F(g^{kl})^{mm}_{ii} F(g^{kl})^{nn}_{jj} - F(g^{kl})^{nm}_{ji} F(g^{kl})^{mn}_{ij} \geq 0$. Combining this conditions together with (6) we get

$$F(g^{kl})^{mn}_{ij} = 0 \; \forall i,j,m,n \in \mathbb{N} \text{ with } (m,n) \notin \{(k,k),(k,l),(l,k),(l,l)\},$$

i.e. $F(g^{kl})$ has only (k,k), (k,l), (l,k), (l,l)-components.

Analogously (substituting g^{kl*} for g^{kl}) we prove that $F(g^{kl*})$ has only (k,k), (k,l), (l,k), (l,l)-components.

Step 3. First, let us show that the main diagonals of the non-zero components of $F(g^{kl})$ are equal, i.e. $F(g^{kl})^{kk}_{ii} = F(g^{kl})^{kl}_{ii} = F(g^{kl})^{lk}_{ii} = F(g^{kl})^{ll}_{ii}$. Restrict $F(g^{kl})$ to the subspace $\langle e_k \otimes f_i, e_l \otimes f_i\rangle$. In this basis $F(g^{kl})$ has the form

$$\begin{pmatrix} F(g^{kl})^{kk}_{ii} & F(g^{kl})^{lk}_{ii} \\ F(g^{kl})^{kl}_{ii} & F(g^{kl})^{ll}_{ii} \end{pmatrix}.$$

This matrix is positive, hence $F(g^{kl})^{kk}_{ii}F(g^{kl})^{ll}_{ii} - F(g^{kl})^{lk}_{ii}F(g^{kl})^{kl}_{ii} \geq 0$, i.e. $|F(g^{kl})^{kl}_{ii}| \leq \sqrt{F(g^{kl})^{kk}_{ii}F(g^{kl})^{ll}_{ii}}$ (note that $F(g^{kl})^{kk}_{ii}$ and $F(g^{kl})^{ll}_{ii}$ are real and non-negative). Due to (4) $\sum_{i=1}^{\infty} F(g^{kl})^{kk}_{ii} = \sum_{i=1}^{\infty} F(g^{kl})^{kl}_{ii} = \sum_{i=1}^{\infty} F(g^{kl})^{lk}_{ii} = \sum_{i=1}^{\infty} F(g^{kl})^{ll}_{ii} = 1$ and hence

$$\begin{aligned} 1 &= \sum_{i=1}^{\infty} \mathrm{Re}F(g^{kl})^{kl}_{ii} \leq \sum_{i=1}^{\infty} |F(g^{kl})^{kl}_{ii}| \leq \sum_{i=1}^{\infty} \sqrt{F(g^{kl})^{kk}_{ii}F(g^{kl})^{ll}_{ii}} \\ &\leq \sum_{i=1}^{\infty} \frac{F(g^{kl})^{kk}_{ii} + F(g^{kl})^{ll}_{ii}}{2} = 1, \end{aligned}$$

and therefore all parts of the inequality must be equal. We have

$$\sqrt{F(g^{kl})^{kk}_{ii}F(g^{kl})^{ll}_{ii}} = \frac{F(g^{kl})^{kk}_{ii} + F(g^{kl})^{ll}_{ii}}{2} \Longrightarrow F(g^{kl})^{kk}_{ii} = F(g^{kl})^{ll}_{ii}$$

and

$$\mathrm{Re}F(g^{kl})^{kl}_{ii} = |F(g^{kl})^{kl}_{ii}| = F(g^{kl})^{kk}_{ii} \Longrightarrow F(g^{kl})^{kl}_{ii} = F(g^{kl})^{lk}_{ii} = F(g^{kl})^{kk}_{ii}.$$

Hence the diagonal elements $F(g^{kl})^{kk}_{ii} = F(g^{kl})^{kl}_{ii} = F(g^{kl})^{lk}_{ii} = F(g^{kl})^{ll}_{ii}$ are equal.

Secondly, let us show that the corresponding non-diagonal elements of the non-zero components of $F(g^{kl})$ are equal, i.e. $F(g^{kl})^{kk}_{ij} = F(g^{kl})^{kl}_{ij} = F(g^{kl})^{lk}_{ij} = F(g^{kl})^{ll}_{ij}$, where $i \neq j$. Denote $a_i{=}F(g^{kl})^{kk}_{ii}$. Restrict $F(g^{kl})$ to the subspace $\langle e_k \otimes f_i, e_k \otimes f_j, e_l \otimes f_j\rangle$. In this basis $F(g^{kl})$ has the form

$$\begin{pmatrix} F(g^{kl})^{kk}_{ii} & F(g^{kl})^{kk}_{ji} & F(g^{kl})^{lk}_{ji} \\ F(g^{kl})^{kk}_{ij} & F(g^{kl})^{kk}_{jj} & F(g^{kl})^{lk}_{jj} \\ F(g^{kl})^{kl}_{ij} & F(g^{kl})^{kl}_{jj} & F(g^{kl})^{ll}_{jj} \end{pmatrix} = \begin{pmatrix} a_i & \bar{y} & \bar{x} \\ y & a_j & a_j \\ x & a_j & a_j \end{pmatrix} = A.$$

If $a_j = 0$ then obviously $x = y = 0$ as A is positive. If $a_j \neq 0$ then

$$\det A = -y(\bar{y}a_j - \bar{x}a_j) + x(\bar{y}a_j - \bar{x}a_j) = -a_j|y-x|^2 \geq 0 \Longrightarrow x = y,$$

and we have derived $F(g^{kl})^{kk}_{ij} = F(g^{kl})^{kl}_{ij} = F(g^{kl})^{lk}_{ij} = F(g^{kl})^{ll}_{ij}$ for all i, j.

Step 4. Firstly, let us prove $F(g^{kl*})^{kk}_{ii} = -iF(g^{kl*})^{kl}_{ii} = iF(g^{kl*})^{lk}_{ii} = F(g^{kl*})^{ll}_{ii}$, i.e. that this condition holds on the main diagonals of non-zero components of $F(g^{kl*})$. Analogously to the previous step, we get $|F(g^{kl*})^{kl}_{ii}| \leq \sqrt{F(g^{kl*})^{kk}_{ii}F(g^{kl*})^{ll}_{ii}}$ (note that $F(g^{kl*})^{kk}_{ii}$ and $F(g^{kl*})^{ll}_{ii}$ are real and non-negative). Due to (4) $\sum_{i=1}^{\infty} F(g^{kl*})^{kk}_{ii} = \sum_{i=1}^{\infty} F(g^{kl*})^{ll}_{ii} = 1$, $\sum_{i=1}^{\infty} F(g^{kl*})^{kl}_{ii} = i$, $\sum_{i=1}^{\infty} F(g^{kl*})^{lk}_{ii} = -i$ and hence

$$\begin{aligned} 1 &= \sum_{i=1}^{\infty} \mathrm{Im}F(g^{kl*})^{kl}_{ii} \leq \sum_{i=1}^{\infty} |F(g^{kl*})^{kl}_{ii}| \leq \sum_{i=1}^{\infty} \sqrt{F(g^{kl*})^{kk}_{ii}F(g^{kl})^{ll}_{ii}} \\ &\leq \sum_{i=1}^{\infty} \frac{F(g^{kl*})^{kk}_{ii} + F(g^{kl*})^{ll}_{ii}}{2} = 1, \end{aligned}$$

and therefore all parts of inequality must be equal. Analogously to the previous step we have

$$F(g^{kl*})_{ii}^{kk} = F(g^{kl*})_{ii}^{ll}$$

and

$$\mathrm{Im}F(g^{kl*})_{ii}^{kl} = |F(g^{kl*})_{ii}^{kl}| = F(g^{kl*})_{ii}^{kk} \Longrightarrow F(g^{kl*})_{ii}^{kl} = F(g^{kl*})_{ii}^{lk} = F(g^{kl*})_{ii}^{kk}$$

and hence $F(g^{kl*})_{ii}^{kk} = -iF(g^{kl*})_{ii}^{kl} = iF(g^{kl*})_{ii}^{lk} = F(g^{kl*})_{ii}^{ll}$.

Secondly, let us show that this property holds also for corresponding non-diagonal elements of the non-zero components of $F(g^{kl*})$, i.e. $F(g^{kl*})_{ij}^{kk} = -iF(g^{kl*})_{ij}^{kl} = iF(g^{kl*})_{ij}^{lk} = F(g^{kl*})_{ij}^{ll}$ if $i \neq j$.

Denote $a_i = F(g^{kl*})_{ii}^{kk}$. Restrict $F(g^{kl*})$ to the subspace $\langle e_k \otimes f_i, e_k \otimes f_j, e_l \otimes f_j \rangle$. In this basis $F(g^{kl*})$ has the form

$$\begin{pmatrix} F(g^{kl*})_{ii}^{kk} & F(g^{kl*})_{ji}^{kk} & F(g^{kl*})_{ji}^{lk} \\ F(g^{kl*})_{ij}^{kk} & F(g^{kl*})_{jj}^{kk} & F(g^{kl*})_{jj}^{lk} \\ F(g^{kl*})_{ij}^{kl} & F(g^{kl*})_{jj}^{kl} & F(g^{kl*})_{jj}^{ll} \end{pmatrix} = \begin{pmatrix} a_i & \bar{y} & \bar{x} \\ y & a_j & ia_j \\ x & -ia_j & a_j \end{pmatrix} = A.$$

If $a_j = 0$ then obviously $x = iy = 0$ as A is positive. If $a_j \neq 0$ then

$$\det A = -y(\bar{y}a_j - i\bar{x}a_j) + x(-i\bar{y}a_j - \bar{x}a_j) = -a_j|iy - x|^2 \geq 0 \Longrightarrow x = iy,$$

and hence $F(g^{kl*})_{ij}^{kk} = -iF(g^{kl*})_{ij}^{kl} = iF(g^{kl*})_{ij}^{lk} = F(g^{kl*})_{ij}^{ll}$ holds for all i, j.

Step 5. Firstly, let us show that the main diagonals of non-zero components of all $F(g^{kl}), k \leq l$, are equal, i.e. we have to prove

$$F(g^{kk})_{ii}^{kk} = F(g^{ll})_{ii}^{ll} = F(g^{kl})_{ii}^{kl} \text{ for all } k < l.$$

Consider $g(t) = g^{kl} + tg^{kk} + pg^{ll}$, where $t + p + tp \geq 0, p + 1 \geq 0$ and $k < l$. The operator $g(t)$ is positive hence $F(g(t))$ is also positive. Restrict $F(g(t))$ to the subspace $\langle e_k \otimes f_i, e_l \otimes f_i \rangle$. In this basis $F(g(t))$ has the form

$$\begin{pmatrix} F(g(t))_{ii}^{kk} & F(g(t))_{ii}^{lk} \\ F(g(t))_{ii}^{kl} & F(g(t))_{ii}^{ll} \end{pmatrix} = \begin{pmatrix} F(g^{kl})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk} & F(g^{kl})_{ii}^{lk} \\ F(g^{kl})_{ii}^{kl} & F(g^{kl})_{ii}^{ll} + pF(g^{ll})_{ii}^{ll} \end{pmatrix}.$$

This matrix is positive, hence

$$(F(g^{kl})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk})(F(g^{kl})_{ii}^{ll} + pF(g^{ll})_{ii}^{ll}) \geq (F(g^{kl})_{ii}^{kl})^2$$

(note that $F(g^{kl})_{ii}^{kl}$, $F(g^{kl})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk}$, and $F(g^{kl})_{ii}^{ll} + pF(g^{ll})_{ii}^{ll}$ are real and non-negative). We apply Lemma 1 with $a = F(g^{kk})_{ii}^{kk}$, $b = F(g^{kl})_{ii}^{kl}$, $c = F(g^{ll})_{ii}^{ll}$, which gives us

$$F(g^{kk})_{ii}^{kk} = F(g^{ll})_{ii}^{ll} \leq F(g^{kl})_{ii}^{kl}.$$

Taking into account the fact that $\sum_{i=1}^{\infty} F(g^{kk})_{ii}^{kk} = \sum_{i=1}^{\infty} F(g^{ll})_{ii}^{ll} = \sum_{i=1}^{\infty} F(g^{kl})_{ii}^{kl} = 1$ we get

$$F(g^{kk})_{ii}^{kk} = F(g^{ll})_{ii}^{ll} = F(g^{kl})_{ii}^{kl}.$$

Secondly, let us show that the remaining elements of the non-zero components of $F(g^{kl}), k \leq l$ are equal. For that purpose we prove

$$F(g^{kk})_{ij}^{kk} = F(g^{ll})_{ij}^{ll} = F(g^{kl})_{ij}^{kl} \text{ for all } i \neq j \text{ and all } k < l$$

using again the operator $g(t)$. Denote $a_i = F(g^{kl})_{ii}^{kk}$. Restrict $F(g(t))$ to the subspace $\langle e_k \otimes f_i, e_k \otimes f_j, e_l \otimes f_j \rangle$. In this basis $F(g(t))$ has the form

$$\begin{aligned}
&\begin{pmatrix} F(g(t))_{ii}^{kk} & F(g(t))_{ji}^{kk} & F(g(t))_{ji}^{lk} \\ F(g(t))_{ij}^{kk} & F(g(t))_{jj}^{kk} & F(g(t))_{jj}^{lk} \\ F(g(t))_{ij}^{kl} & F(g(t))_{jj}^{kl} & F(g(t))_{jj}^{ll} \end{pmatrix} \\
= &\begin{pmatrix} F(g^{kl})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk} & F(g^{kl})_{ji}^{kk} + tF(g^{kk})_{ji}^{kk} & F(g^{kl})_{ji}^{lk} \\ F(g^{kl})_{ij}^{kk} + tF(g^{kk})_{ij}^{kk} & F(g^{kl})_{jj}^{kk} + tF(g^{kk})_{jj}^{kk} & F(g^{kl})_{jj}^{lk} \\ F(g^{kl})_{ij}^{kl} & F(g^{kl})_{jj}^{kl} & F(g^{kl})_{jj}^{ll} + pF(g^{ll})_{jj}^{ll} \end{pmatrix} \\
= &\begin{pmatrix} a_i + ta_i & x + ty & x \\ \bar{x} + t\bar{y} & a_j + ta_j & a_j \\ \bar{x} & a_j & a_j + pa_j \end{pmatrix} = A.
\end{aligned}$$

If $a_j = 0$ then obviously $x = y = 0$ as A is positive. If $a_j \neq 0$ then

$$\begin{aligned}
\det A &= a_i a_j^2 (1+t)^2 (1+p) + (x+ty) a_j \bar{x} + x(\bar{x} + t\bar{y}) a_j - x a_j (1+t) \bar{x} \\
&\quad -(x+ty)(\bar{x}+t\bar{y}) a_j (1+p) - a_i a_j^2 (1+t) \\
&= -\frac{a_j}{1+t} |x(1+t) - (x+ty)|^2 \geq 0 \Longrightarrow x = y.
\end{aligned}$$

This means that

$$F(g^{kl})_{ij}^{kk} = F(g^{kk})_{ij}^{kk}.$$

Step 6. Firstly, let us show that the main diagonals of non-zero components of all $F(g^{kl*}), k < l$, satisfy the equality

$$F(g^{kk})_{ii}^{kk} = -iF(g^{kl*})_{ii}^{kl} = iF(g^{kl*})_{ii}^{lk} = F(g^{ll})_{ii}^{ll}.$$

Thereby we use the same arguments as in the previous step considering the operator $g^*(t) = g^{kl*} + tg^{kk} + pg^{ll}$, where $t + p + tp \geq 0, p + 1 \geq 0$ and $k < l$. The operator $g^*(t)$ is positive, hence $F(g^*(t))$ is also positive. Restrict $F(g^*(t))$ to the subspace $\langle e_k \otimes f_i, e_l \otimes f_i \rangle$. In this basis $F(g^*(t))$ has the form

$$\begin{pmatrix} F(g^*(t))_{ii}^{kk} & F(g^*(t))_{ii}^{lk} \\ F(g^*(t))_{ii}^{kl} & F(g^*(t))_{ii}^{ll} \end{pmatrix} = \begin{pmatrix} F(g^{kl*})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk} & -F(g^{kl*})_{ii}^{kl} \\ F(g^{kl*})_{ii}^{kl} & F(g^{kl*})_{ii}^{ll} + pF(g^{ll})_{ii}^{ll} \end{pmatrix}. \quad (7)$$

Note that $-iF(g^{kl*})_{ii}^{kl}$, $F(g^{kl*})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk}$, and $F(g^{kl*})_{ii}^{ll} + pF(g^{ll})_{ii}^{ll}$ are real and non-negative. The matrix (7) is positive, hence

$$(F(g^{kl*})_{ii}^{kk} + tF(g^{kk})_{ii}^{kk})(F(g^{kl*})_{ii}^{kk} + pF(g^{ll})_{ii}^{ll}) \geq -|F(g^{kl*})_{ii}^{kl}|^2 = |-iF(g^{kl*})_{ii}^{kl}|^2.$$

We apply Lemma 1 with $a = F(g^{kk})^{kk}_{ii}$, $b = -iF(g^{kl*})^{kl}_{ii}$, $c = F(g^{ll})^{ll}_{ii}$, which gives us

$$F(g^{kk})^{kk}_{ii} = F(g^{ll})^{ll}_{ii} \leq -iF(g^{kl*})^{kl}_{ii}.$$

Taking into account the fact that $\sum_{i=1}^{\infty} F(g^{kk})^{kk}_{ii}) = \sum_{i=1}^{\infty} F(g^{ll})^{ll}_{ii} = -\sum_{i=1}^{\infty} iF(g^{kl*})^{kl}_{ii} = 1$ we get

$$F(g^{kk})^{kk}_{ii} = -iF(g^{kl*})^{kl}_{ii} = iF(g^{kl*})^{lk}_{ii} = F(g^{ll})^{ll}_{ii} \text{ for all } k < l.$$

Secondly, let us show that the remaining elements of the non-zero components of $F(g^{kl*})$, $k < l$, satisfy

$$F(g^{kk})^{kk}_{ij} = F(g^{kl*})^{kk}_{ij} \quad \text{if } i \neq j$$

using again the operator $g^*(t)$. Denote $a_i = F(g^{kl*})^{kk}_{ii}$. Restrict $F(g^*(t))$ to the subspace $\langle e_k \otimes f_i, e_k \otimes f_j, e_l \otimes f_j \rangle$. In this basis $F(g^*(t))$ has the form

$$\begin{pmatrix} F(g^*(t))^{kk}_{ii} & F(g^*(t))^{kk}_{ji} & F(g^*(t))^{lk}_{ji} \\ F(g^*(t))^{kk}_{ij} & F(g^*(t))^{kk}_{jj} & F(g^*(t))^{lk}_{jj} \\ F(g^*(t))^{kl}_{ij} & F(g^*(t))^{kl}_{jj} & F(g^*(t))^{ll}_{jj} \end{pmatrix}$$

$$= \begin{pmatrix} F(g^{kl*})^{kk}_{ii} + tF(g^{kk})^{kk}_{ii} & F(g^{kl*})^{kk}_{ji} + tF(g^{kk})^{kk}_{ji} & F(g^{kl*})^{lk}_{ji} \\ F(g^{kl*})^{kk}_{ij} + tF(g^{kk})^{kk}_{ij} & F(g^{kl*})^{kk}_{jj} + tF(g^{kk})^{kk}_{jj} & F(g^{kl*})^{lk}_{jj} \\ F(g^{kl*})^{kl}_{ij} & F(g^{kl*})^{kl}_{jj} & F(g^{kl*})^{ll}_{jj} + pF(g^{ll})^{ll}_{jj} \end{pmatrix}$$

$$= \begin{pmatrix} a_i + ta_i & x + ty & -ix \\ \bar{x} + t\bar{y} & a_j + ta_j & -ia_j \\ i\bar{x} & ia_j & a_j + pa_j \end{pmatrix} = A.$$

If $a_j = 0$ then obviously $x = y = 0$ as A is positive. If $a_j \neq 0$ then, analogously to the previous step,

$$\det A = -\frac{a_j}{1+t}|x(1+t) - (x+ty)|^2 \geq 0 \Longrightarrow x = y.$$

This means that

$$F(g^{kl*})^{kk}_{ij} = F(g^{kk})^{kk}_{ij}.$$

Denote $a_{ij} = F(g^{11})^{11}_{ij}$ and consider $\rho_E \in \mathcal{L}_1^+(H_S)$ that has the form a_{ij} in the basis $\{e_i, i \in \mathbb{N}\}$. It is easy to see now that $F(\rho) = \rho \otimes \rho_E$ for each $\rho \in \mathcal{L}_1(H)$. The theorem is proved. $\square$

Remark 3 *The theorem implies that the linear lifting F is continuous.*

Remark 4 *If we skip the constraint $\mathrm{tr}_{H_E} F(\rho) = \rho$, more general liftings are possible. Let $\rho_E \in \mathcal{D}(H_E)$ be a reference state, and K_n a family of bounded operators in H which satisfy $\sum_n K_n^+ K_n = Id$, then*

$$\rho \longmapsto F(\rho) = \sum_n K_n \left(\rho \otimes \rho_E\right) K_n^+ \tag{8}$$

is a linear and continuous mapping $\mathcal{D}(H_S) \to \mathcal{D}(H)$. Such liftings are used in general investigations of the process of measurement [9] and in information theory, see e. g. [10].

Remark 5 *It is well known that any mixed state ρ of a system $\mathcal{S}$ can be obtained as the reduced state of a pure state in an extended system $\mathcal{S} \times \mathcal{E}$, if only $\dim H_E \geq \dim H_S$, see e.g. [11]. But due to Theorem 1 the pure state cannot depend linearly on the state ρ. The representation by a pure state is actually a generalization of the classical Gram's theorem from linear algebra. To see this let H_S be realized as $\mathcal{L}_2(\Omega, \mathcal{B}_\Omega, \mu_\Omega)$ where Ω is a set , $\mathcal{B}_\Omega$ is a σ-algebra of its subsets, μ_Ω a non-negative σ-additive measure on $\mathcal{B}_\Omega$. Then the space $H = H_S \otimes H_E$ is isomorphic to the space $\mathcal{L}_2(\Omega, \mathcal{B}_\Omega, \mu_\Omega, H_E)$ of H_E-valued Bochner square μ_Ω-integrable functions on Ω. The corresponding isomorphic map $H_S \otimes H_E \to \mathcal{L}_2(\Omega, \mathcal{B}_\Omega, \mu_\Omega, H_E)$ is denoted by φ. On the other hand, the space $H_S \otimes H_S$ can be realized as $\mathcal{L}_2(\Omega \times \Omega, \mathcal{B}_\Omega \otimes \mathcal{B}_\Omega, \mu_\Omega \otimes \mu_\Omega)$, and hence the space $\mathcal{L}_1^+(H_S)$ can be considered as a vector subspace of the latter space which includes all Hilbert-Schmidt operators in H_S. Any normalized vector $a \in H_S \otimes H_E$, $\|a\| = 1$, spans a one-dimensional subspace of $H_S \otimes H_E$ and defines a unique projection operator $P_a \in \mathcal{D}(H_S \otimes H_E)$. If $f_a \in \mathcal{L}_2(\Omega, \mathcal{B}_\Omega, \mu_\Omega, H_E)$ is defined by $f_a = \varphi(a)$ then the reduced state of pure state P_a is given by*

$$S(\omega_1, \omega_2) = \langle f_a(\omega_1), f_a(\omega_2) \rangle_{H_E} . \tag{9}$$

Now the generalization of Gram's theorem can be formulated as follows: For any $S \in \mathcal{L}_1^+(H_S)$ there exists a vector $a \in H_S \otimes H_E$, $\|a\| = 1$, for which (9) holds. If Ω is a finite set and μ_Ω is the counting measure, we obtain the classical Gram's theorem.

3 REDUCING OBSERVABLES

The problem of linear liftings of states is closely related to the problem of reducing observables of the total system H to observables of the subsystem H_S.

Lemma 2 *Let $F : \mathcal{L}_1(H_S) \longrightarrow \mathcal{L}_1(H_S \otimes H_E)$ be a continuous mapping and let $F^* : \mathcal{L}(H_S \otimes H_E) \longrightarrow \mathcal{L}(H_S)$ be its adjoint mapping; then $F^*(B \otimes Id_E) = B$ for all $B \in \mathcal{L}(H_S)$ iff $\mathrm{tr}_{H_E} F(\rho) = \rho$ for all $\rho \in \mathcal{L}_1(H_S)$.*

Proof If $B \in \mathcal{L}(H_S)$ then, according to the definition of the duality between $\mathcal{L}(H)$ and $\mathcal{L}_1(H)$, $\langle B \otimes Id_E, F(\rho) \rangle = tr_H (B \otimes Id_E) F(\rho) = \mathrm{tr}_{H_S} B\rho = \langle B, \rho \rangle$. This identity together with the definition of the duality between $\mathcal{L}(H_S)$ and $\mathcal{L}_1(H_S)$ implies that

$$F^*(B \otimes Id) = B. \tag{10}$$

On the other hand, if F^* satisfies (10) then, for $B \in \mathcal{L}(H_S)$ and $\rho \in \mathcal{L}_1(H_S)$,

$$\langle B \otimes Id_E, F(\rho) \rangle = \langle F^*(B \otimes Id_E), \rho \rangle = \langle B, \rho \rangle = \mathrm{tr}_{H_S} B\rho. \tag{11}$$

But $\langle B \otimes Id_E, F(\rho) \rangle = \mathrm{tr}_{H_S} B(\mathrm{tr}_{H_E} F(\rho))$. Hence $\langle B, \rho \rangle = \mathrm{tr}_{H_S} B\rho = \langle B, \mathrm{tr}_{H_E} F(\rho) \rangle$, and as the latter identity holds for any B, we finally obtain $\rho = \mathrm{tr}_{H_E} F(\rho)$. The lemma is proved. □

Theorem 1 and Lemma 2 imply the following theorem.

Theorem 2 *If $R : \mathcal{L}(H_S \otimes H_E) \to \mathcal{L}(H_S)$ is a linear mapping, continuous in the ultraweak or $(\sigma(\mathcal{L}(H), \mathcal{L}_1(H)), \sigma(\mathcal{L}(H_S), \mathcal{L}_1(H_S)))$ topology, see e.g. Sec. VI.6 of [4], and if $R(B \otimes Id_E) = B$ is true for all $B \in \mathcal{L}(H_S)$ then there exists an element $\rho_E \in \mathcal{D}(H_S)$ such that $R(A) = \mathrm{tr}_{H_E} A(Id_S \otimes \rho_E)$ for all $A \in \mathcal{L}(H)$.*

4 PROBABILITY MEASURES

The classical analog of the case considered in Theorem 1 is much simpler and admits non-factorizing answers. Let T be a topological space, then $\mathcal{C}_b(T)$ is the vector space of all bounded continuous functions on T, $\mathcal{M}(T)$ is the vector space of all Borel (signed) measures on T equipped with the topology $\sigma(\mathcal{M}(T), \mathcal{C}_b(T))$, and $\mathcal{M}_p(T)$ is the closed convex set of probability measures on T. The Dirac measure at point $t \in T$ will be denoted by δ_t. Let Q and P be topological spaces, $E = Q \times P$ the product space, and $\mathcal{G} : \mathcal{M}_p(E) \to \mathcal{M}_p(Q)$ be the mapping induced by the projection $\mathrm{pr}_Q : E \to Q$. The mapping $\mathcal{G}$ can be (uniquely) extended by linearity to an $\mathbb{R}$-linear mapping $\mathcal{M}(E) \to \mathcal{M}(Q)$. For any measure $\mu \in \mathcal{M}(E)$ the measure $\mathcal{G}\mu \in \mathcal{M}(Q)$ is called the marginal of μ. The right inverse of $\mathcal{G}$ will be called a lifting.

Lemma 3 *Let $f : Q \to \mathcal{M}_p(E)$ be a continuous function such that $\mathcal{G}f(q) = \delta_q$, then the mapping $\mathcal{F} : \mathcal{M}(Q) \to \mathcal{M}(E)$ defined by*

$$\mathcal{F}\upsilon := \int_Q f(q)\upsilon(dq) \tag{12}$$

is a linear lifting. Any linear lifting has this representation.

Proof Take the Dirac measure δ_q then the integral is $\mathcal{F}\delta_q = f(q) \in \mathcal{M}_p(E)$ and we have $\mathcal{G}\mathcal{F}\delta_q = \mathcal{G}f(q) = \delta_q$. The general case follows by linearity and continuity. On the other hand, if G is a linear lifting, then (12) follows with the function $f(q) = \mathcal{F}\delta_q$. □

If $f(q)$ factorizes into $f(q) = \delta_q \times \chi$ with $\chi \in \mathcal{M}(P)$, the lifting (12) factorizes into $\mathcal{F}(\upsilon) = \upsilon \times \chi$. But one can obviously choose non-factorizing functions $f(q)$ such that $\mathcal{F}(\upsilon)$ is not a product measure. To give an explicit example we split Q into two disjoint measurable sets $Q = Q_1 \cup Q_2$ and denote by $\chi_1(q)$ and $\chi_2(q)$ the characteristic functions of the sets Q_1 and Q_2. Then

$$f(q) = \chi_1(q)\delta_q \times \delta_{p_1} + \chi_2(q)\delta_q \times \delta_{p_2} \tag{13}$$

with two points $p_j \in P$, $j = 1, 2$, $p_1 \neq p_2$, yields an example of a non-factorizing lifting.

The state space $\mathcal{D}(H)$ of a quantum mechanical system is a closed convex set with the pure states $\mathcal{P}(H)$ as extremal points. Any $W \in \mathcal{D}(H)$ can be represented by the Choquet integral [7]

$$W = \int_{\mathcal{P}(H)} P\,\mu(dP) \tag{14}$$

where $\mu(dP)$ is a – in general non-unique – measure in the convex set $\mathcal{M}_p(\mathcal{P}(H))$ of probability measures on $\mathcal{P}(H)$, see e. g. [12]. This representation relates the quantum mechanical state space with the space of probability measures, and one might ask whether it is possible to find an affine linear mapping $\gamma : \mathcal{D}(H) \to \mathcal{M}(\mathcal{P}(H))$ such that (12) is valid for all $W \in \mathcal{D}(H)$ with the measure $\mu(dP) = \gamma_W(dP)$.

Theorem 3 *There does not exist an affine linear mapping $\gamma : \mathcal{D}(H) \to \mathcal{M}_p(\mathcal{P}(H))$ such that the representation (12) holds for all $W \in \mathcal{D}(H)$ with $\mu(dP) = (\gamma W)(dP)$.*

Proof If such a mapping γ exists then any pure state has to be represented by an atomic measure on the one-point set containing just this pure state. Moreover this mapping can be extended to an $\mathbb{R}$-linear mapping $\gamma : \mathcal{L}_1^a(H) \to \mathcal{M}(\mathcal{P}(H))$. Since there are finite sets of pure states which are linearly dependent in $\mathcal{L}_1^a(H)$ – e.g. any four projection operators on the Hilbert subspace $\mathbb{C}^2$ of $\mathcal{H}$ – whereas the set of atomic measures is linear independent in $\mathcal{M}(\mathcal{P}(H))$ we obtain contradiction to the linearity of γ. □

The proof given here exploits the different structures of the convex sets $\mathcal{D}(H)$ and $\mathcal{M}_p(\mathcal{P}(H))$: the space of measures is a simplex whereas $\mathcal{D}(H)$ not. Theorem 3 is also closely related to our main Theorem 1, it is actually a consequence of it. To see that implication assume such an affine linear mapping γ exists. Then the lifting problem of Sec. 2. has the following solution in contradiction to Theorem 1.

In the first step the statistical operator $\rho \in \mathcal{D}(H_S)$ is mapped onto the measure $\gamma\rho \in \mathcal{M}_p(\mathcal{P}(H_S))$. Following Lemma 3 we can lift this measure to a measure $\sigma \in \mathcal{M}_p(\mathcal{P}(H_S) \times \mathcal{P}(H_E))$. Thereby we can choose a lifting such that σ is not a product measure, take e.g. (13). The operator

$$W = \int_{\mathcal{P}(H_S)\times\mathcal{P}(H_E)} P_S \otimes P_E \, \sigma(dP_S \times dP_E) \tag{15}$$

has the partial trace $\mathrm{tr}_{H_E} W = \int_{\mathcal{P}(H_S)} P_S \, (\gamma\rho)(dP_S) = \rho$. All steps of the mapping $\rho \to W$ are affine linear. Since the measure σ does not factorize, the statistical operator W has not the product form $\rho \otimes \rho_E$, and we have obtained a contradiction to Theorem 1.

In addition to the representation of states by a probability distribution on the set of pure states there exists a representation of any state by a random vector distributed by a probability measure on the Hilbert space. Such a representation is used in the theory of Schrödinger (-Belavkin) stochastic equations (see [13], [14] and references therein), which gives both, a phenomenological description of continuous measurements, and a Markovian approximations for the reduced dynamics.

By $\mathcal{M}(H)$ we denote the space of all σ-additive signed measures on the σ-algebra of Borel subsets of H. The space of probability measures on H is denoted by $\mathcal{M}_p(H)$, the set of all measures concentrated on $H\backslash\{0\}$ by $\mathcal{M}^0(H)$, and the set of all probability measures concentrated on $H\backslash\{0\}$ by $\mathcal{M}_p^0(H) = \mathcal{M}^0(H) \cap \mathcal{M}_p(H)$.

In the theory of stochastic Schrödinger equations a probability measure $\nu \in \mathcal{M}_p^0(H)$ represents the state $B \in \mathcal{D}(H)$ if

$$\int_H \langle z, Az\rangle \|z\|^{-2} \nu(dz) = \omega_B(A) \equiv \mathrm{tr}_H AB \tag{16}$$

is valid for all observables $A \in \mathcal{L}(H)$. Thereby any measure $\nu \in \mathcal{M}_p^0(H)$ represents a state, and any state $W \in \mathcal{D}(H)$ can be represented by such a measure.

For the proof of the first statement take $A \in \mathcal{L}(H)$. Then the function $|\langle z, Az\rangle| \, \|z\|^{-2}$ is bounded by $\|A\|$ for all $z \neq 0$, and the integral $\omega_\nu(A) := \int_H \|z\|^{-2}\langle z, Az\rangle \nu(dz)$ is defined. Moreover, it is easy to see that all demands of Gleason's theorem, see Remark 1, are fulfilled. Hence there exists a state $W \in \mathcal{D}(H)$ such that $\omega_\nu(A) = \mathrm{tr}_H AW$.

On the other hand, given a statistical operator a probability measure for the representation (16) can be constructed as follows. For any $B \in \mathcal{D}(H)$, let $\nu_B^0 \in \mathcal{M}_p^0(H)$ be a probability measure with the correlation operator B, i.e. for all $z_1, z_1 \in H$ the identity

$\langle z_1, Bz_2\rangle = \int \langle z_1, z\rangle\langle z, z_2\rangle \nu_B^0(dz)$ is true. It is worth noticing that among the measures ν_B^0 there exist precisely one Gaussian measure with zero mean value. The positive measure $\nu_B \in \mathcal{M}^0(H)$ is then defined by $\nu_B = \langle z, z\rangle \nu_B^0 = \|z\|^2 \nu_B^0$; i.e. for any Borel subset $\mathcal{A}$ of $H^{\mathbb{R}}$ we have $\nu_B(\mathcal{A}) = \int_{\mathcal{A}} \langle z, z\rangle \nu_B^0(dz)$. The identity $\mathrm{tr}_H B = 1$ implies that ν_B is a probability measure on H; in fact $\nu_B(H) = \int \langle z, z\rangle \nu_B^0(dz) = \mathrm{tr}_H B = 1$. For any observable $A \in \mathcal{L}(H)$ the function $H \longmapsto \mathbb{R}^1 : z \longmapsto \frac{1}{\|z\|^2}\langle z, Az\rangle$ is a random variable on the probability space (H, ν_B). The mean value $\bar{A}$ of this random variable

$$\bar{A} = \int_H \langle z, Az\rangle \|z\|^{-2} \nu_B(dz) = \int_H \langle z, Az\rangle \nu_B^0(dz) = \mathrm{tr}_H AB$$

is exactly the expectation of the observable A in the state $B \in \mathcal{D}(H)$. Hence the measure $\nu_B \in \mathcal{M}_p^0(H)$ represents the state B.

There exists an affine linear mapping from the measures $\upsilon \in \mathcal{M}_p^0(H)$ into the set of measures of the Choquet representation. Let $\varphi : H\backslash\{0\} \to \mathcal{P}(H)$ be the mapping $a \mapsto P_a$, where P_a is the projection operator onto the subspace $\{\lambda a \mid \lambda \in \mathbb{C}\}$, i.e. $P_a b = \langle b \mid a\rangle \|a\|^{-2} a$ for all $b \in H$. Then the measure $\nu\varphi^{-1} \in \mathcal{M}_p(\mathcal{P}(H))$ is defined by $\nu\varphi^{-1}(\mathcal{R}) = \nu(\varphi^{-1}(\mathcal{R}))$ for any measurable set $\mathcal{R} \subset \mathcal{P}(H)$ of projection operators. This mapping $\nu \mapsto \nu\varphi^{-1}$ is affine linear. If $\nu \in \mathcal{M}_p^0$ represents a state $W \in \mathcal{D}(H)$, then (16) and the definition of $\nu\varphi^{-1}$ yield

$$\langle z_1 \mid Wz_2\rangle \overset{(16)}{=} \int_H \langle z_1 \mid z\rangle \langle z \mid z_2\rangle \|z\|^{-2} \nu(dz) = \int_{\mathcal{P}(H)} \langle z_1 \mid Pz_2\rangle (\nu\varphi^{-1})(dP).$$

But that means $W = \int_{\mathcal{P}(H)} P\,(\nu\varphi^{-1})(dP)$, and $\nu\varphi^{-1}$ is the Choquet measure of the state W.

The measures in the representation (16) are highly nonunique; the arbitrariness is even larger than in the case of the Choquet representation, and one might ask again for an affine linear lifting $\mathcal{D}(H) \to \mathcal{M}_p^0(H)$. But assume such an affine linear lifting $\gamma : \mathcal{D}(H) \to \mathcal{M}_p^0(H)$ exists, then it induces an affine linear lifting $\mathcal{D}(H) \to \mathcal{M}_p(\mathcal{P}(H))$ by $W \mapsto \gamma(W) \mapsto (\gamma(W))\,\varphi^{-1}$ and we have obtained a contradiction to Theorem 3.

Corollary 1 *There does not exist an affine linear mapping $\gamma : \mathcal{D}(H) \to \mathcal{M}_p^0(H)$ such that for any $W \in \mathcal{D}(H)$ the measure $\gamma(W)$ represents the state W.*

Acknowledgment

This work was done during a stay of O. G. Smolyanov at the University of Kaiserslautern. OGS would like to thank the Deutsche Forschungsgemeinschaft (DFG) for financial support.

References

[1] C. Favre and P. A. Martin, Dynamique quantique des systémes amortis ≪non markoviens≫, *Helv. Phys. Acta*, 41, 333–361 (1968)

[2] Ph. A. Martin, *Modèles en Méchanique Statistique des Processus Irréversibles*, Chap III, Lect. Notes in Physics Vol. 103 (Springer, Berlin, 1979)

[3] E. B. Davies, *Quantum Theory of Open Systems*, Chap 10 (Academic Press, London, 1976)

[4] M. Reed and B. Simon, *Methods of Modern Mathematical Physics I: Functional Analysis*, (Academic Press, New York, 1972)

[5] W. Guz, On quantum dynamical semigroups, *Rep. Math. Phys.*, 6, 455–464 (1974)

[6] A. M. Gleason, Measures on the closed subspaces of a Hilbert space, *J. Math. and Mech.*, 6, 885–893 (1957)

[7] G. Choquet, *Lectures on Analysis. Vol II*, (Benjamin, New York, 1969)

[8] E. C. Beltrametti and G. Casinelli, *The Logic of Quantum Mechanics.* (Addison-Wesley, Reading, 1981)

[9] K. Kraus, *States, Effects and Operations*, Lecture Notes in Physics 190 (Springer, Berlin, 1983)

[10] L. Accardi and M. Ohya, Compound channels, transition expectations, and liftings, *Appl. Math. Optim.*, 39, 33–59 (1999).

[11] A. Peres, *Quantum Theory: Concepts and Methods*, (Kluwer, Dordrecht, 1993)

[12] J. Kupsch, The structure of the quantum mechanical state space and induced superselection rules, *Pramana - J. Phys.*, 51(5), 615–624 (1998), quant-ph/9612033

[13] S. Albeverio, V. N. Kolokol'tsov, and O. G. Smolyanov, Continuous quantum measurement: local and global approaches, *Rev. Math. Phys.*, 9, 827–840 (1997)

[14] V. P. Belavkin and O. G. Smolyanov, Feynman path integral corresponding to the Schrödinger stochastic equation, *Doklady Math.*, 57, 430–434 (1998)