

(Non-)Contextuality of Physical Theories as an Axiom

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QIP 2011

Plan

1. Introduction:

non-contextuality

2. Results:

a general framework to study non-contextuality;
non-contextuality and non-locality

3. Open problems:

theoretical;
applied;
a complexity perspective.

1. Introduction:

non-contextuality for

1. classical theories
2. non-signaling theories
3. quantum theory

“measuring”
means

“to ask questions to a system”

“measuring”
means

“to ask questions to a system”

what is your velocity?

**“measuring”
means**

“to ask questions to a system”

what is your velocity?

**what is your angular
momentum?**

**“measuring”
means**

“to ask questions to a system”

what is your velocity?

**what is the
colour of your
eyes?**

**what is your angular
momentum?**

**“measuring”
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“to ask questions to a system”

what is your velocity?

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**what is the
colour of your
eyes?**

is your entropy 5 bits?

**“measuring”
means**

“to ask questions to a system”

what is your velocity?

**what is your angular
momentum?**

**what is the
colour of your
eyes?**

is your entropy 5 bits?

Question 1

Question 5

Question 2

Question 4

Question 3

context:

a set of

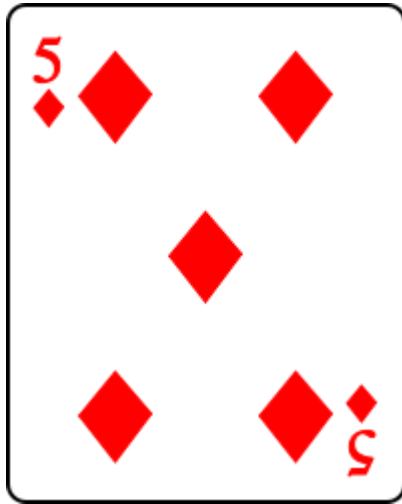
mutually compatible
questions

Question 1

Question 2

Question 3

compatibility:



Question 1:

Is the colour red?

Answer:

Yes

Question 2:

What is the suit?

Answer:

Diamonds

non-contextuality:

answers

do not depend
on contexts

Context 1

Question 1:

Is X true?

Answer:

Yes

Question 2:

Is Y true?

Answer:

Yes

non-contextuality:

answers

do not depend
on contexts

Question 2:

Is Y true?

Answer:

Yes

non-contextuality:

answers

do not depend
on contexts

Question 3:

Is Z true?

Answer:

No

Context 1

Question 1:

Is X true?

Answer:

Yes

Context 2

Question 2:

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answers

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answers

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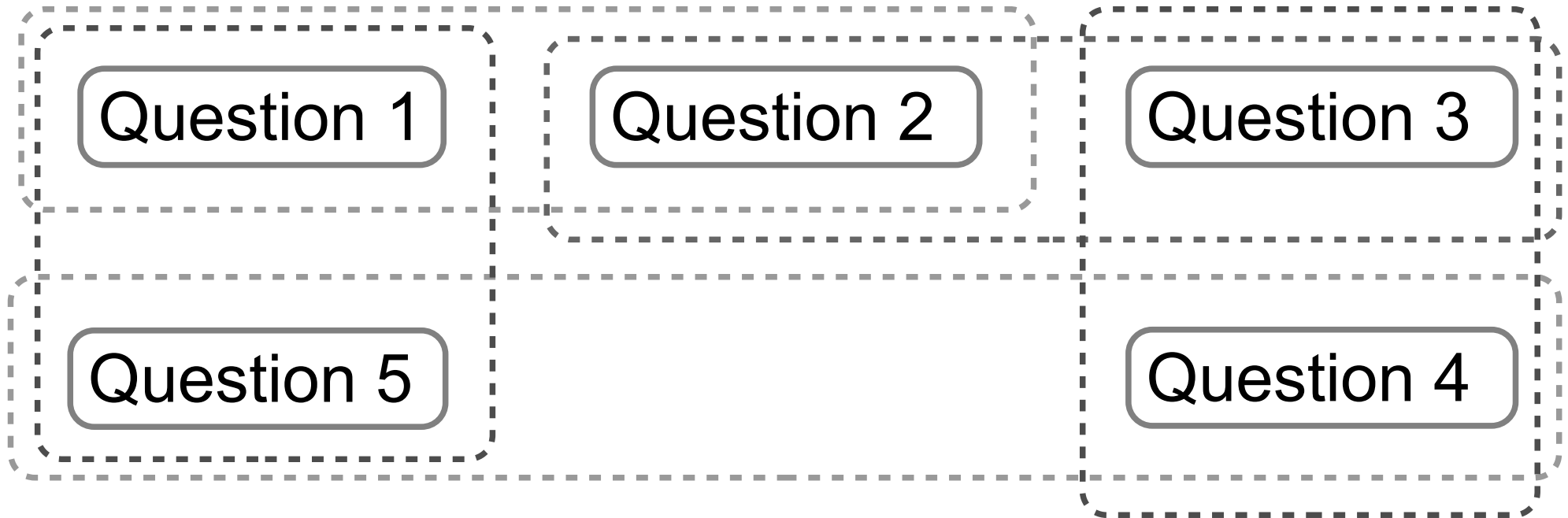
Question 3:

Is Z true?

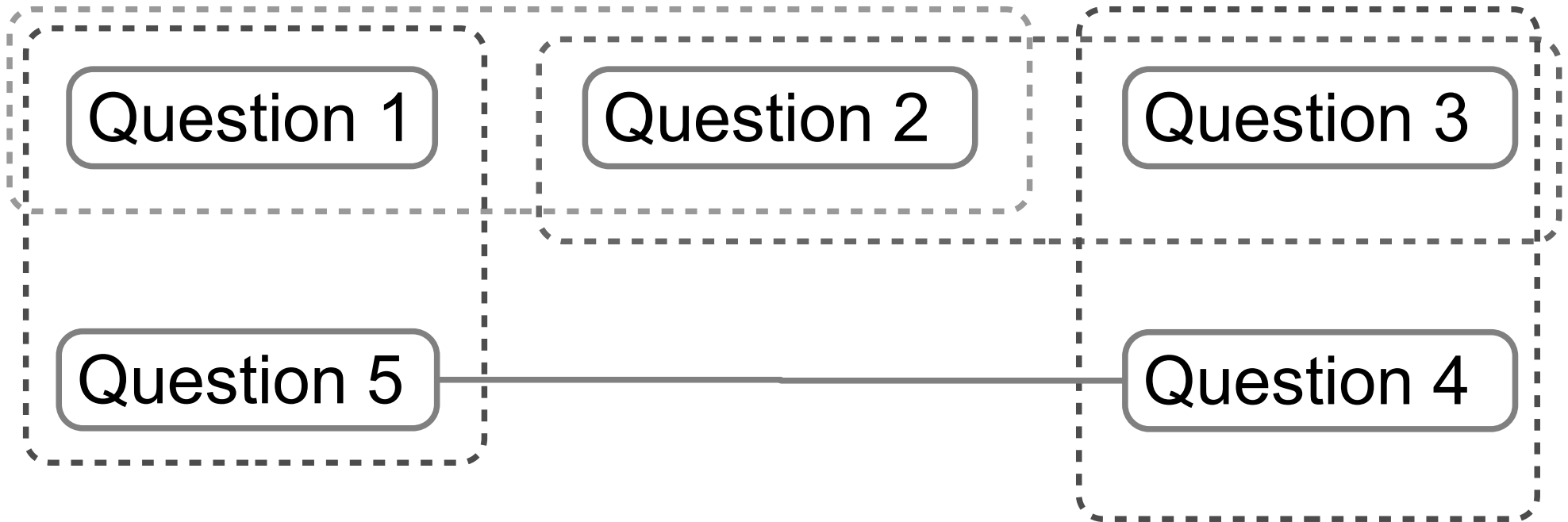
Answer:

No

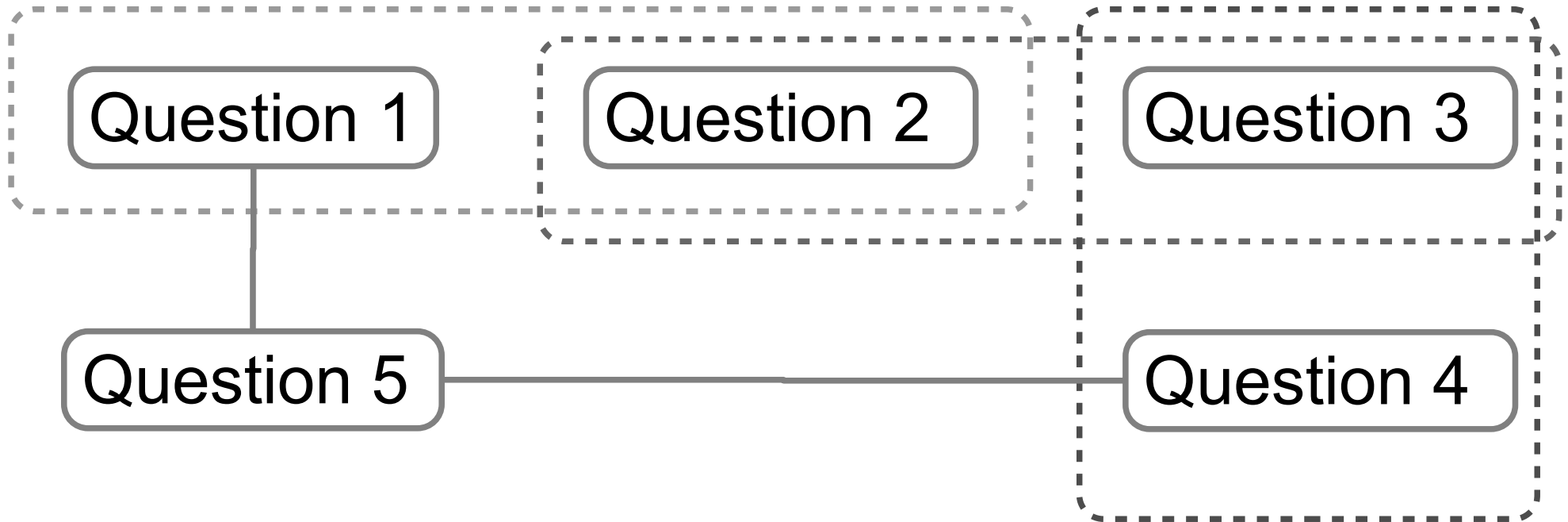
compatibility structure



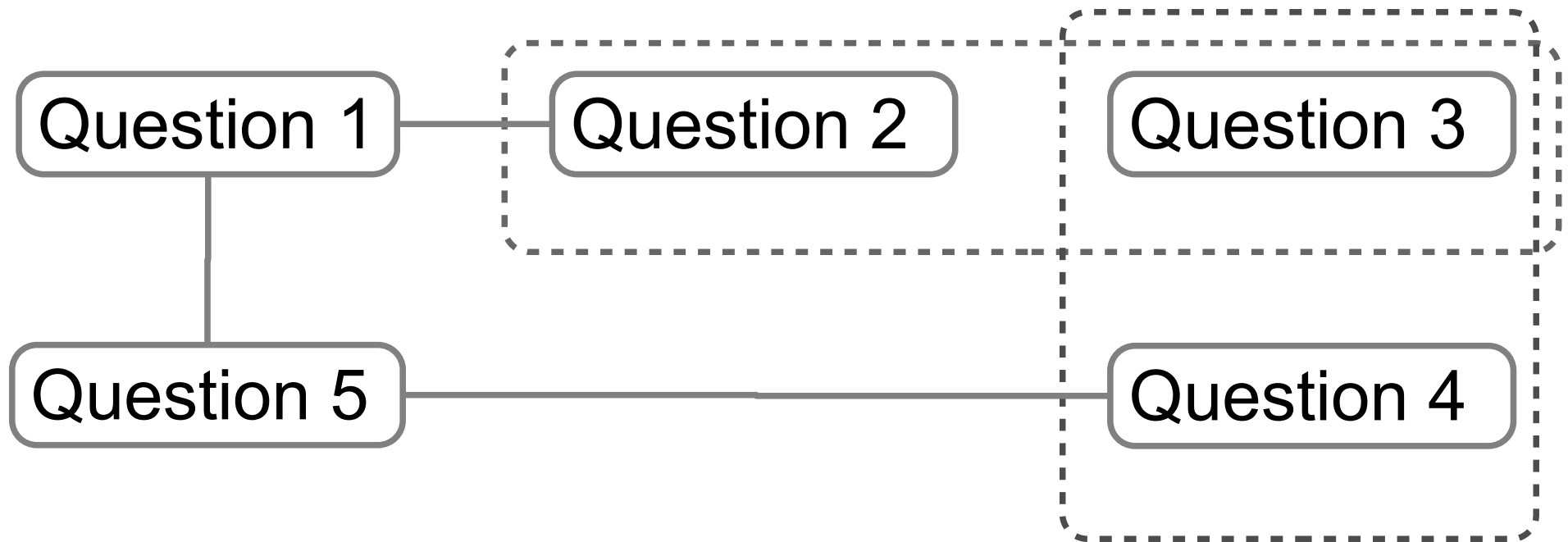
compatibility structure



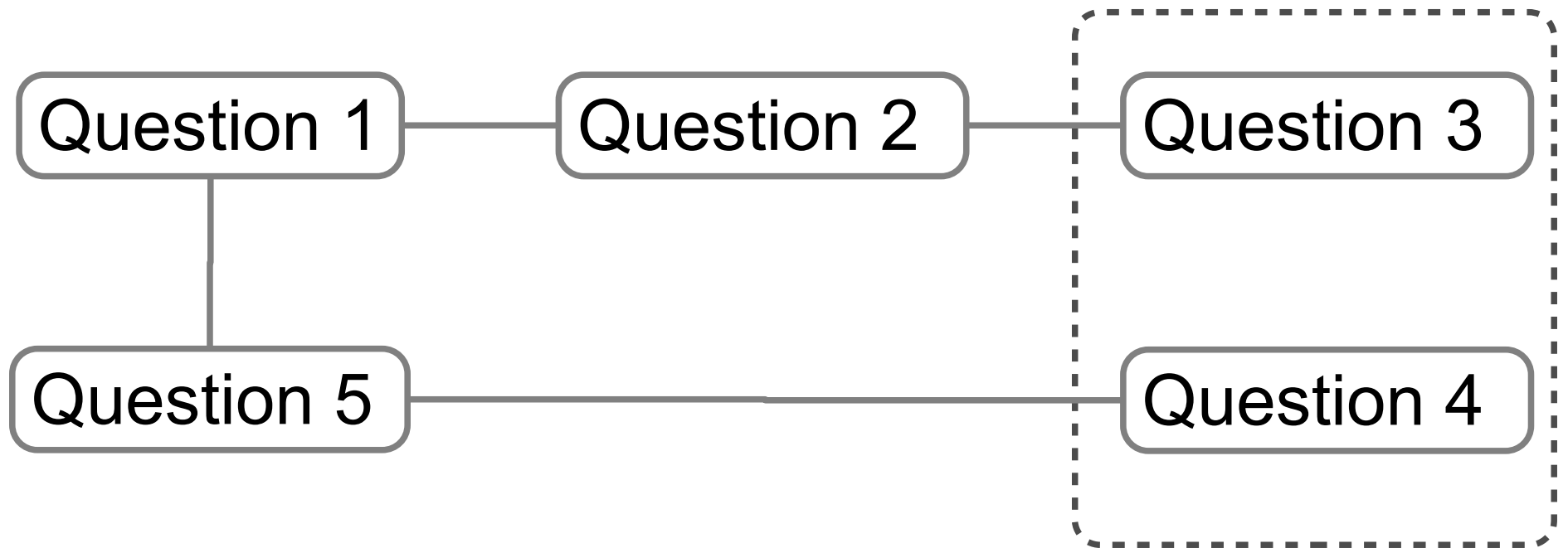
compatibility structure



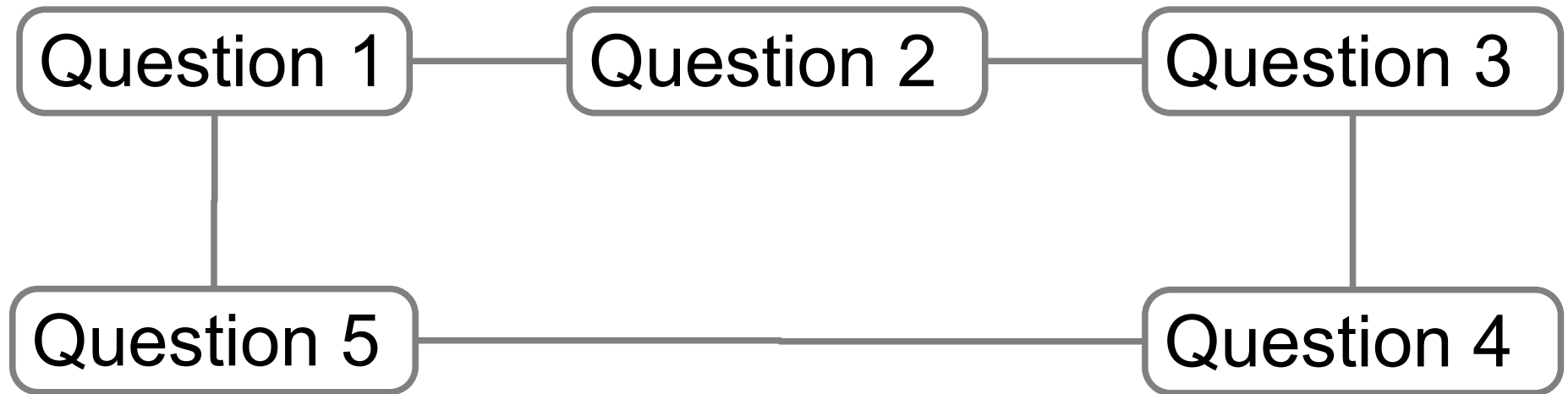
compatibility structure



compatibility structure

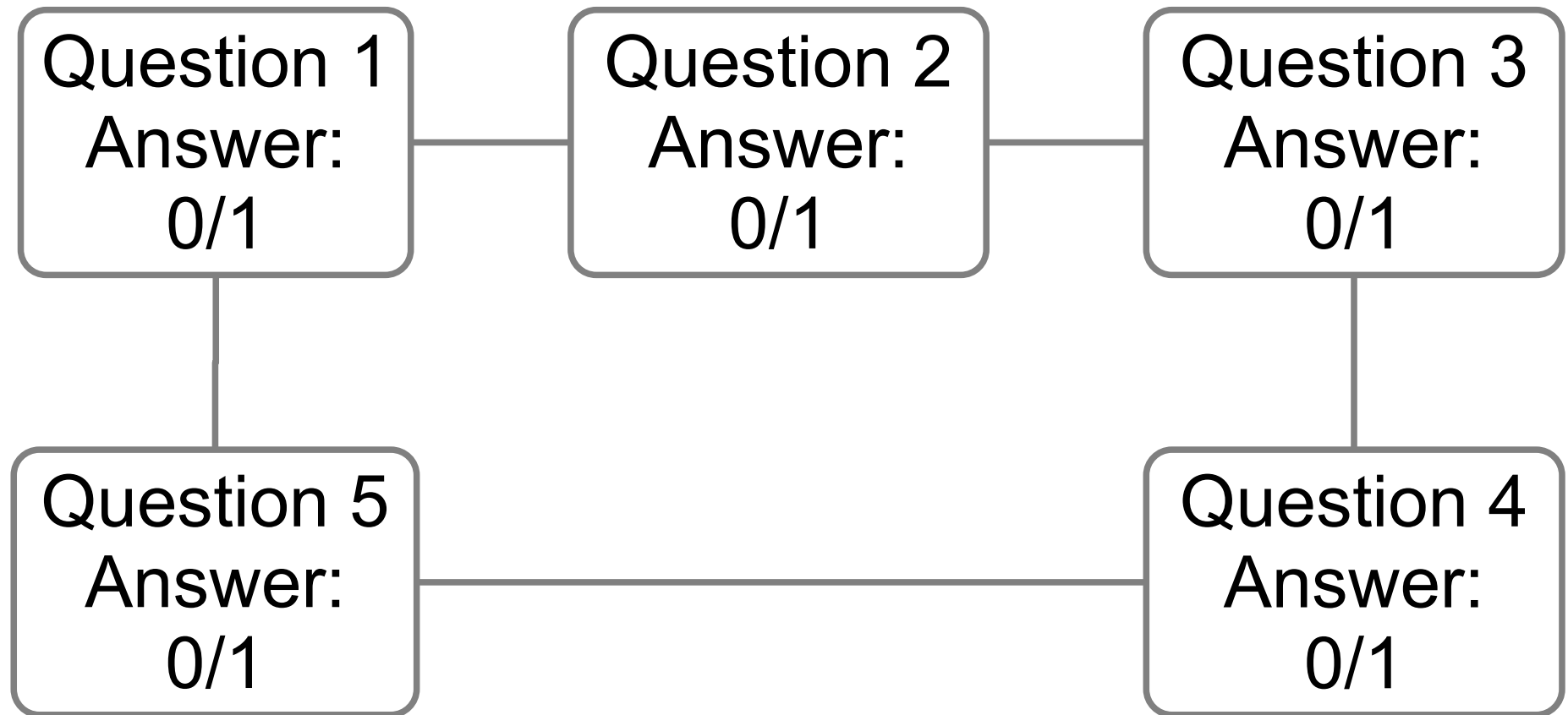


compatibility structure

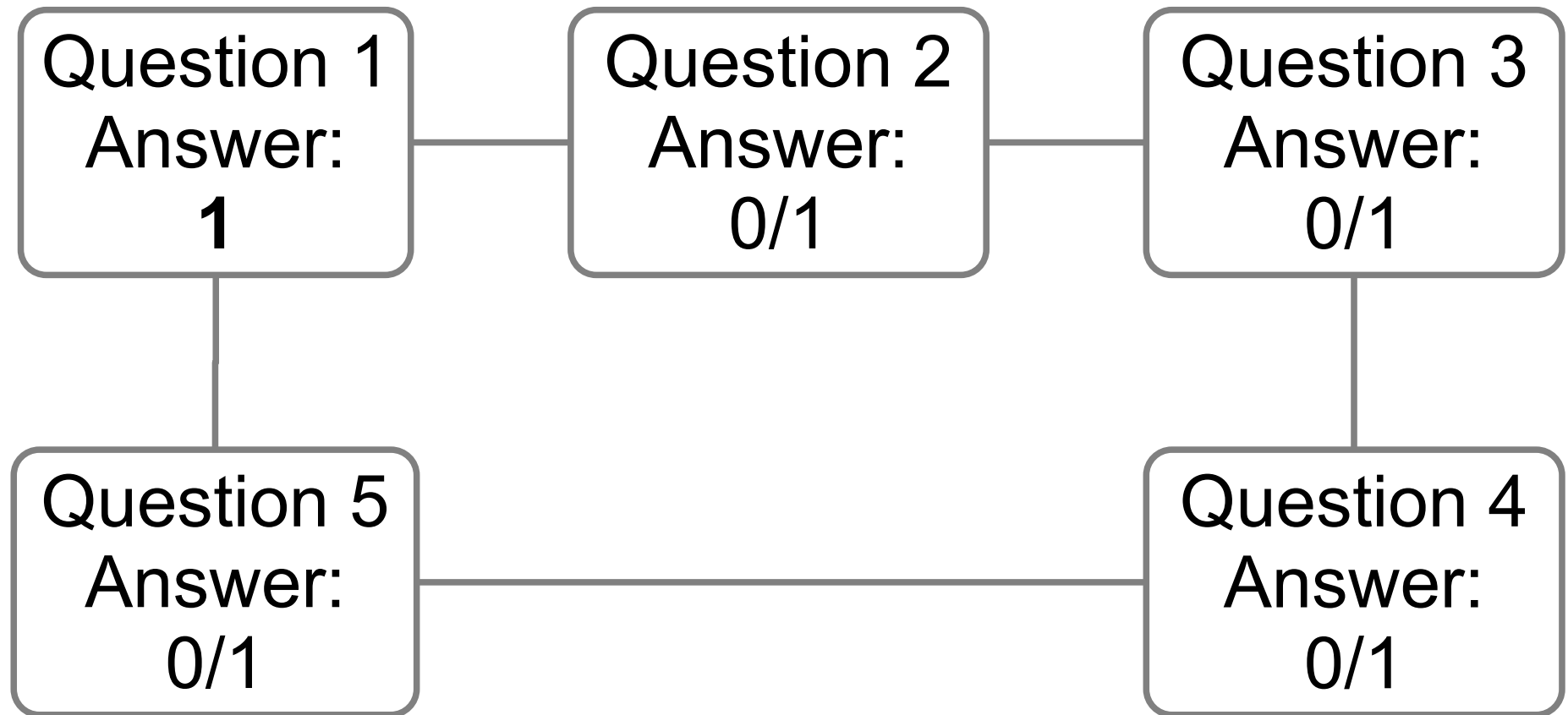


exclusiveness:

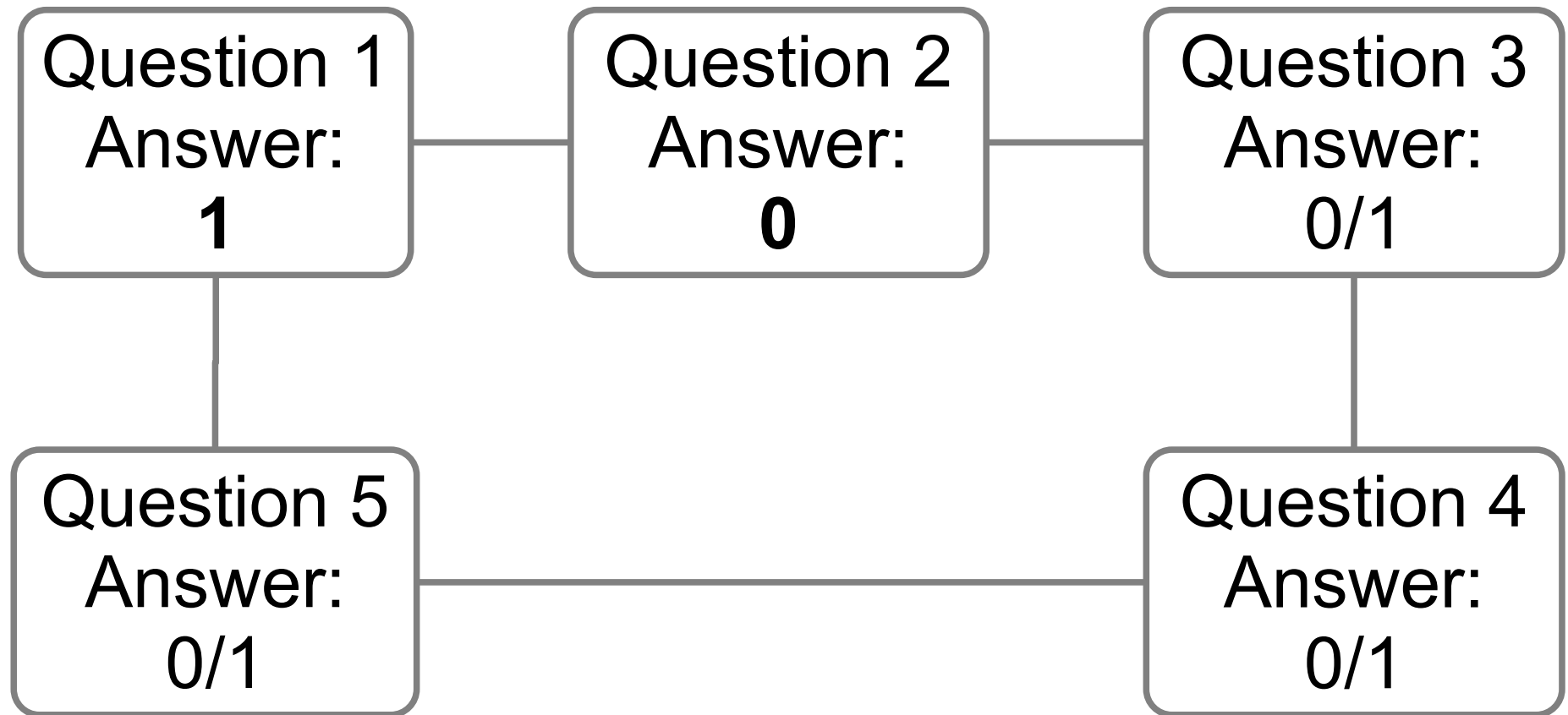
answers to
adjacent questions
are **not both 1**



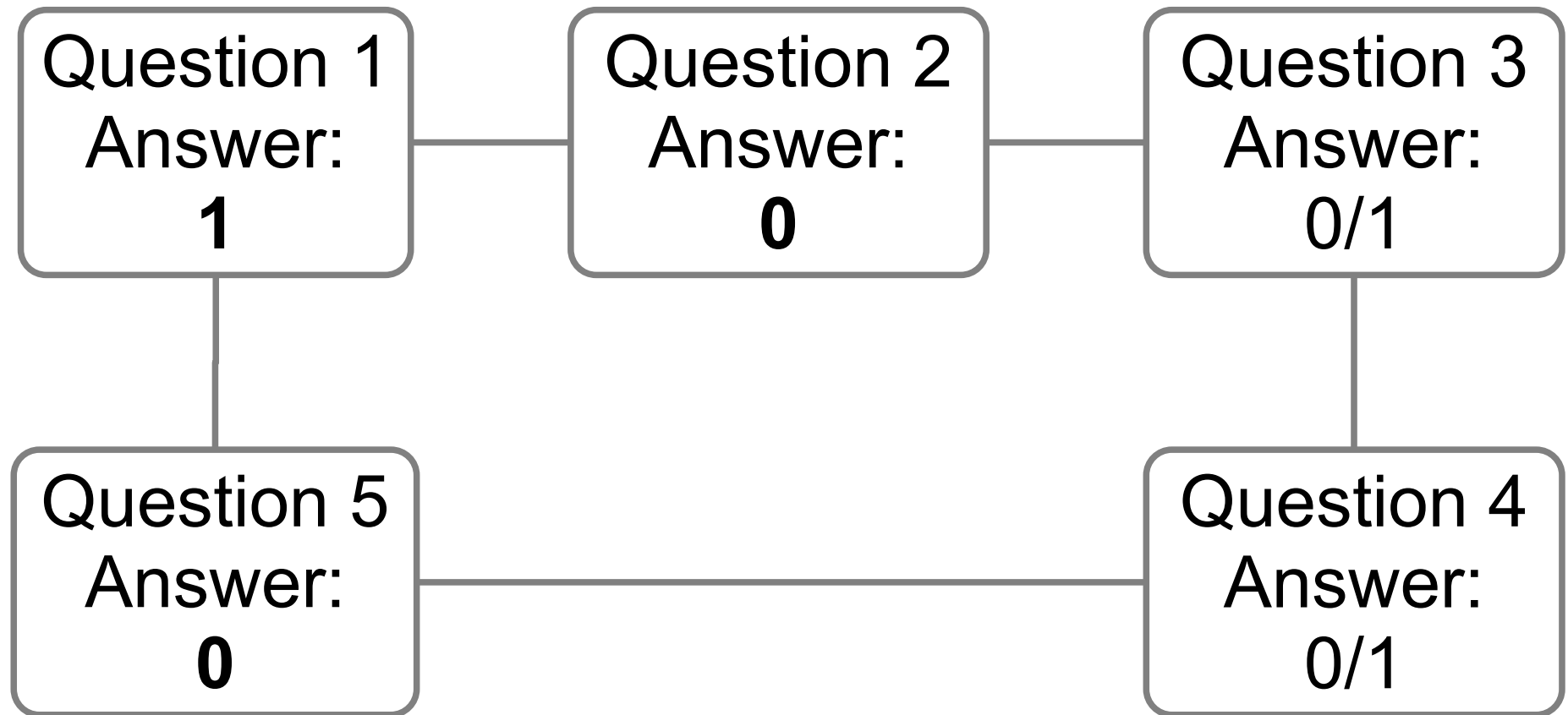
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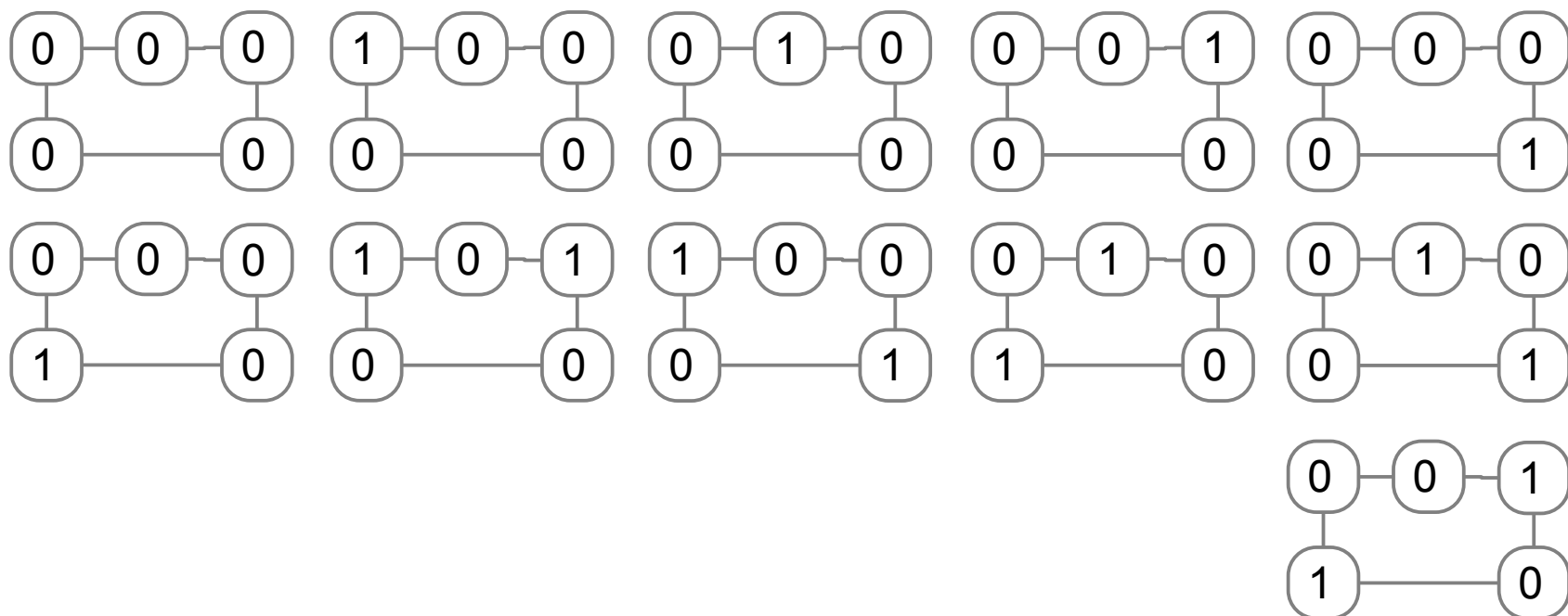
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exclusiveness:
answers to
adjacent questions
are **not both 1**



the answers 1 have expectation ≤ 2
if non-contextual and exclusive

classical theories
give expectation ≤ 2

non-contextuality:
probabilities of answers
do not depend
on contexts

Context 1

Question 1:

Is X true?

$\text{Pr}[\text{answer Yes}] = a$

Question 2:

Is Y true?

$\text{Pr}[\text{answer Yes}] = b$

non-contextuality:

probabilities of answers
do not depend
on contexts

non-contextuality:
probabilities of answers
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Question 2:

Is Y true?

$$\Pr[\text{answer Yes}] = b$$

Question 3:

Is Z true?

$$\Pr[\text{answer Yes}] = c$$

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Question 2:

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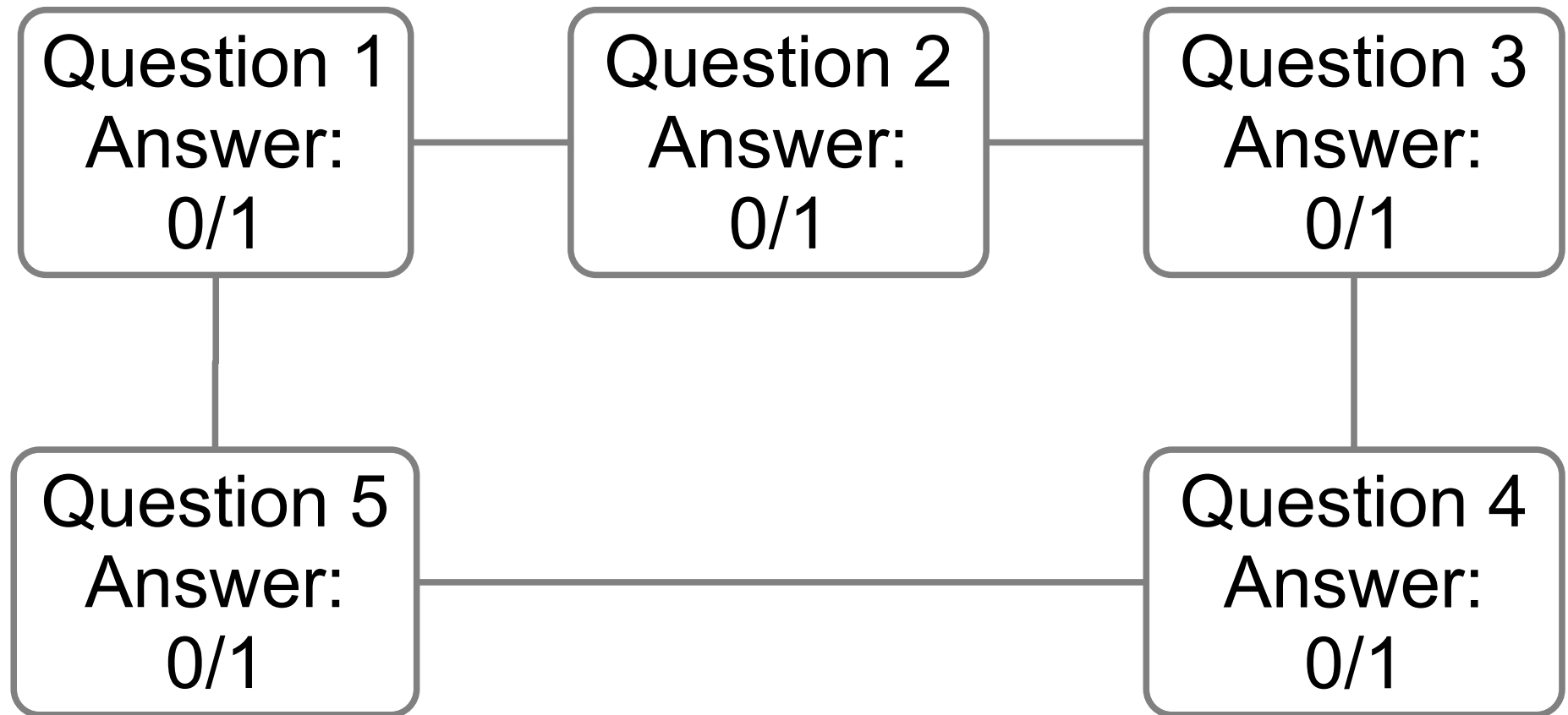
Question 3:

Is Z true?

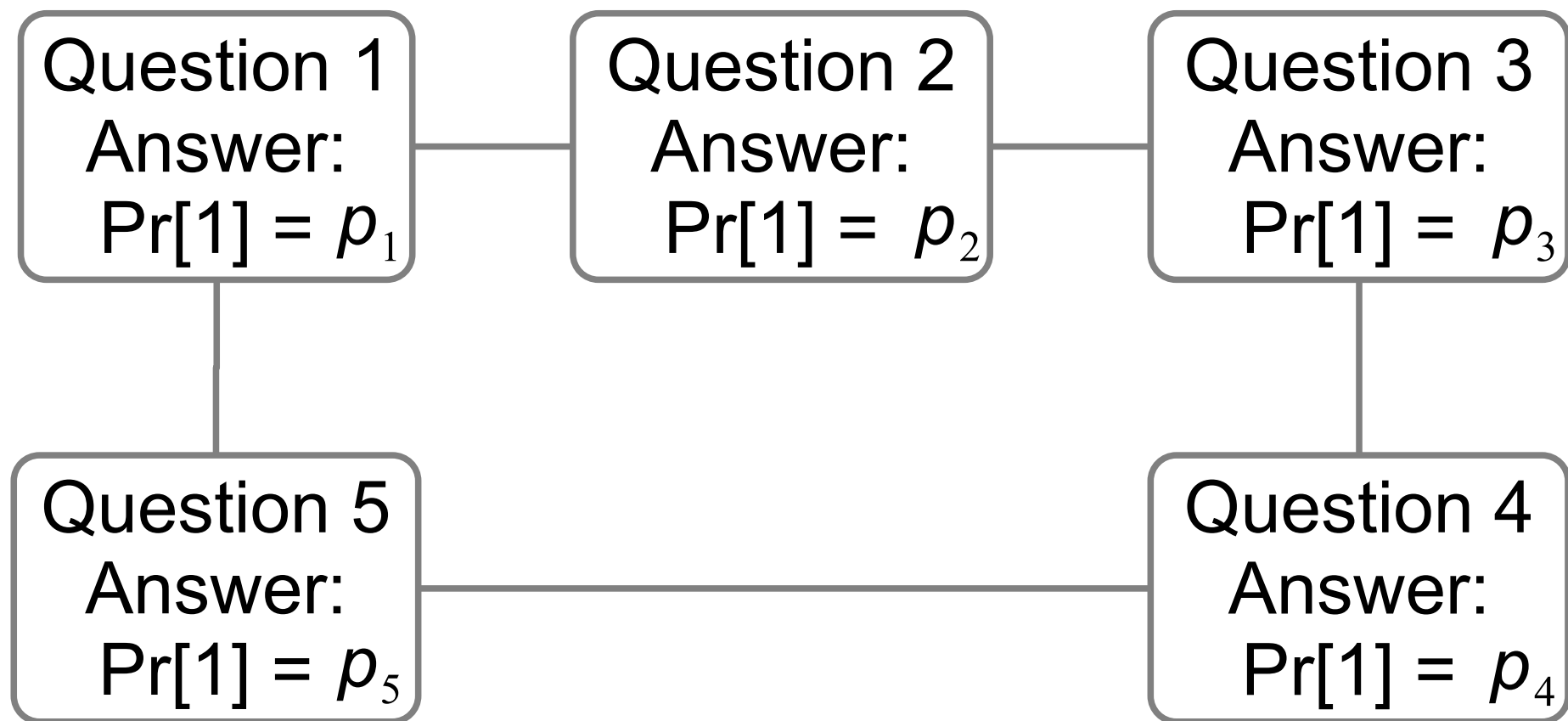
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exclusiveness:

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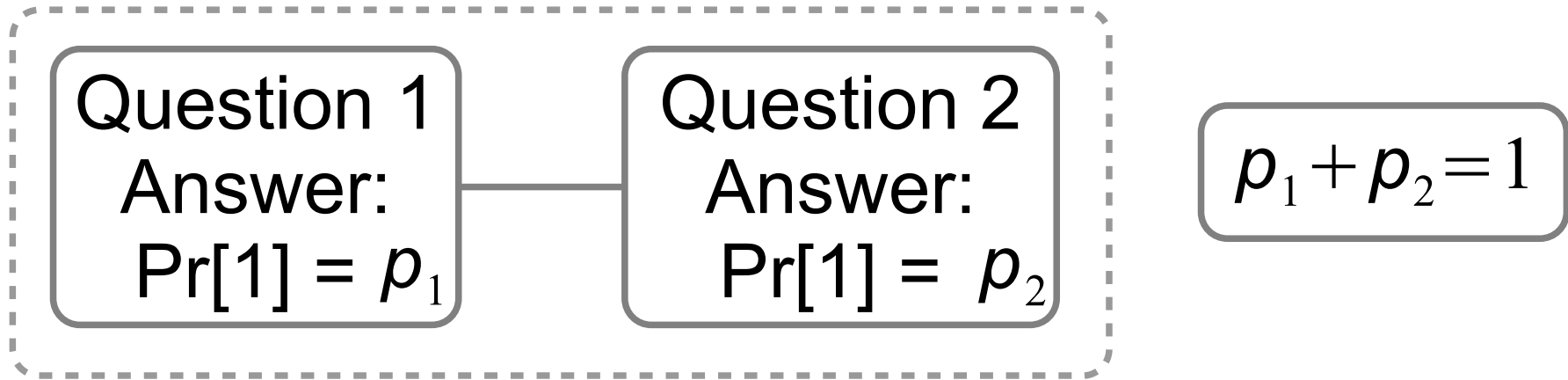


exclusiveness:
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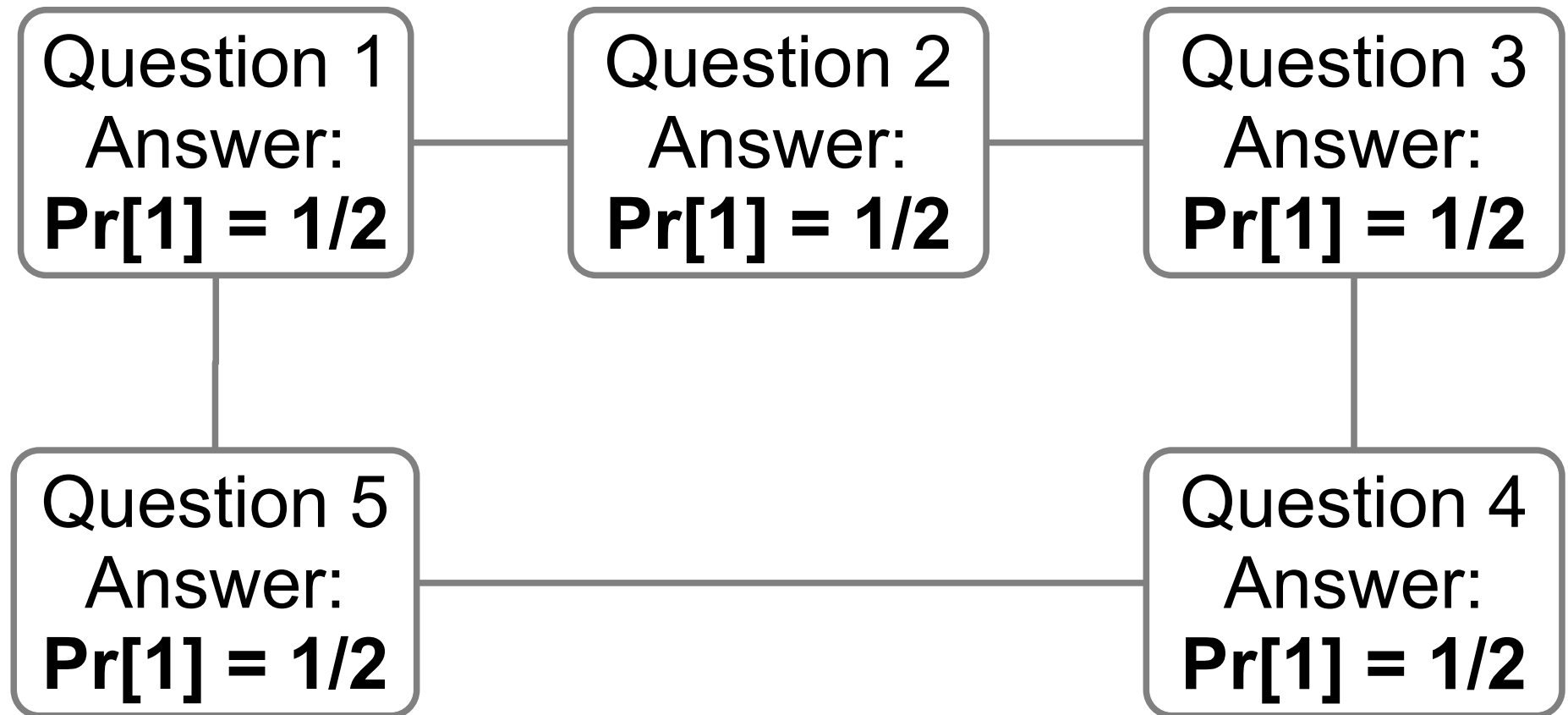


exclusiveness:
answers to
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Context 1



exclusiveness:
answers to
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are **not both 1**



exclusiveness:
answers to
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are **not both 1**

non-signaling theories*
give expectation ≤ 2.5

axiomatically:

classical theories

give expectation ≤ 2

non-signaling theories

give expectation ≤ 2.5

axiomatically:

classical theories

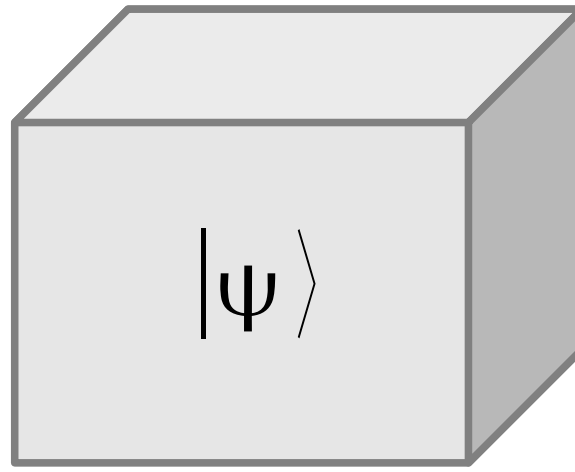
give expectation ≤ 2

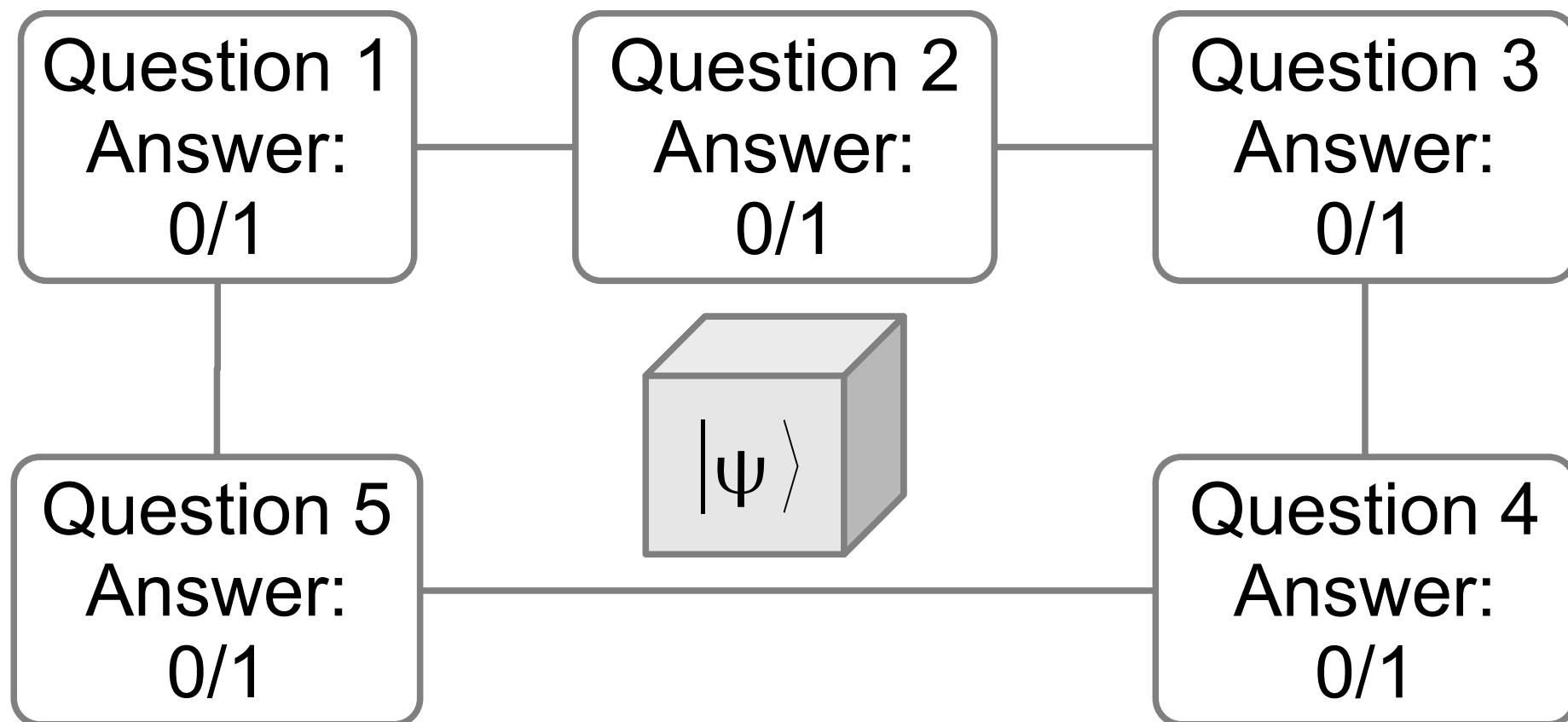
quantum theory

give expectation $\leq ?$

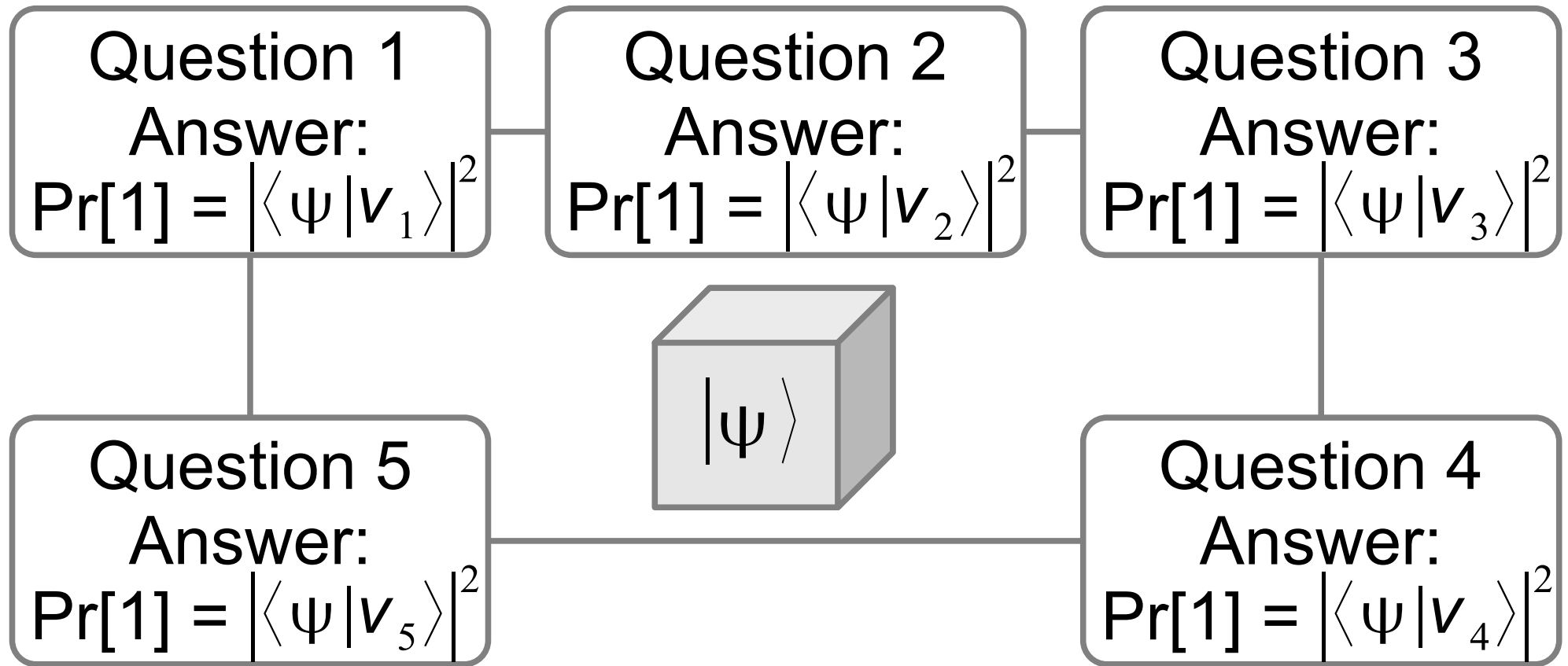
non-signaling theories

give expectation ≤ 2.5

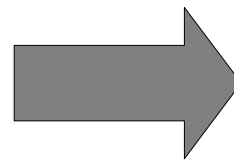




exclusiveness:
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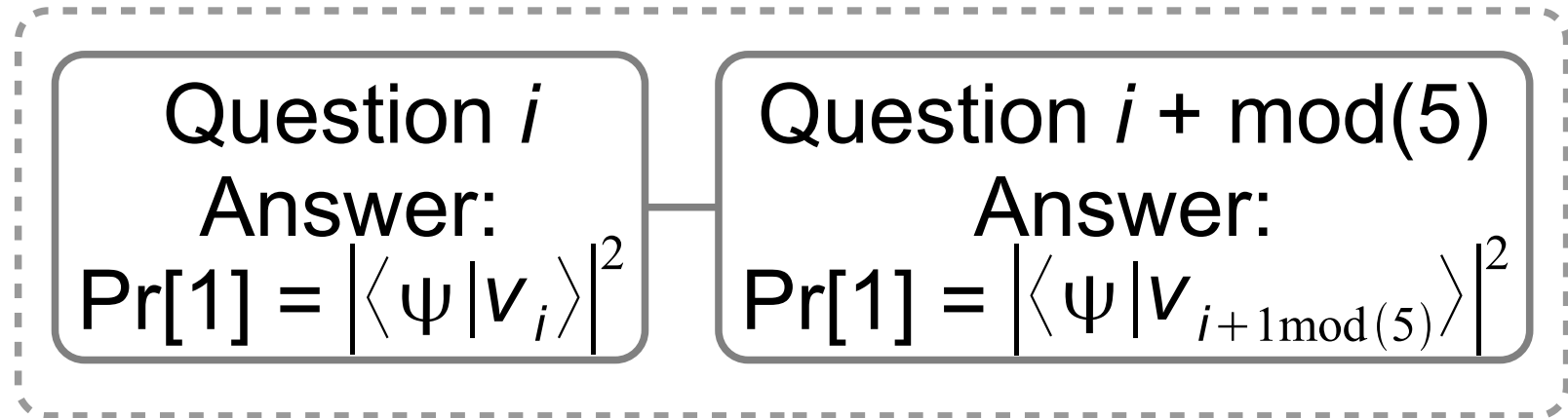
extra axiom:
Born rule



expectation:

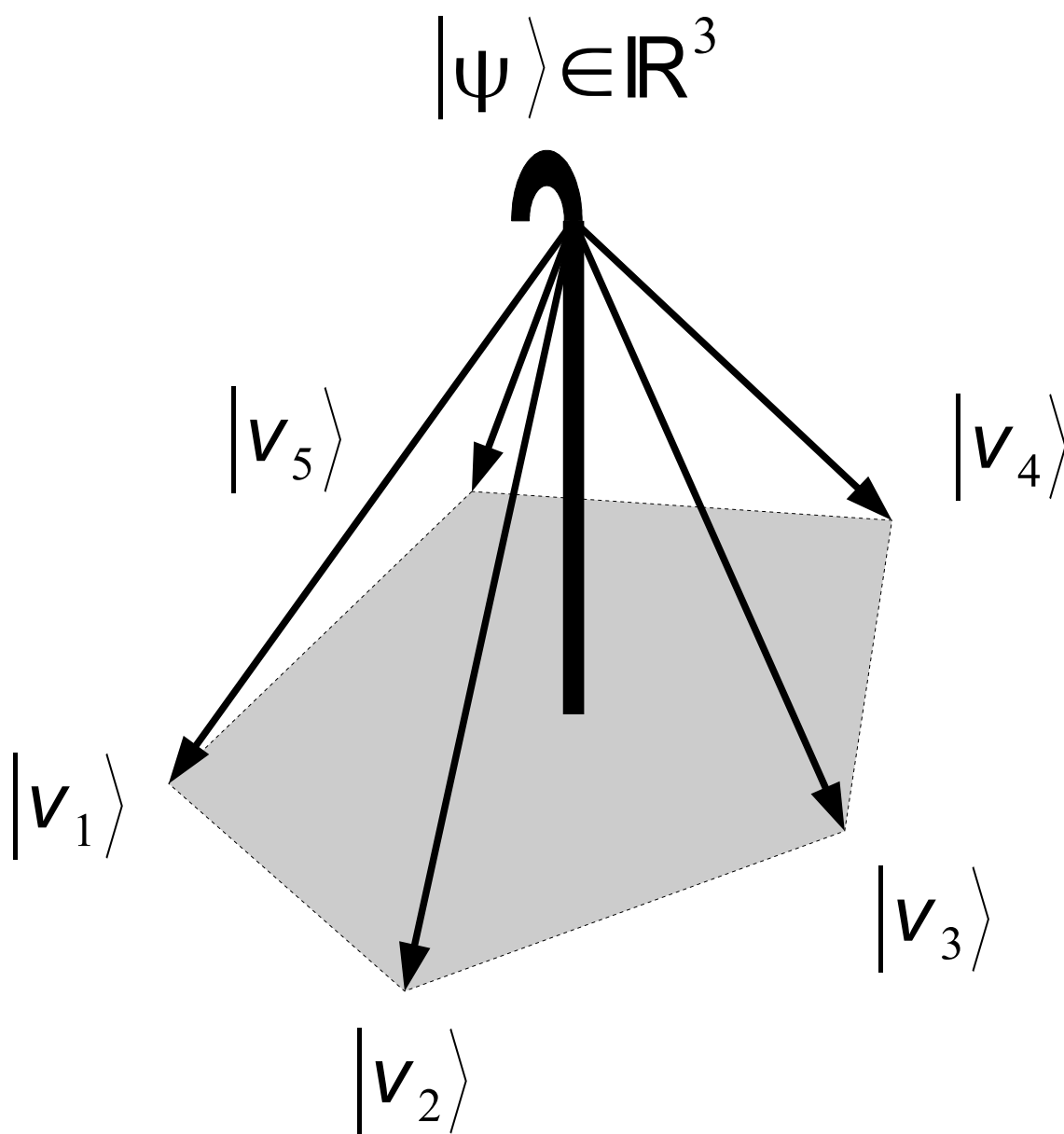
$$\sum_{i=1}^5 |\langle \psi | v_i \rangle|^2$$

Context i

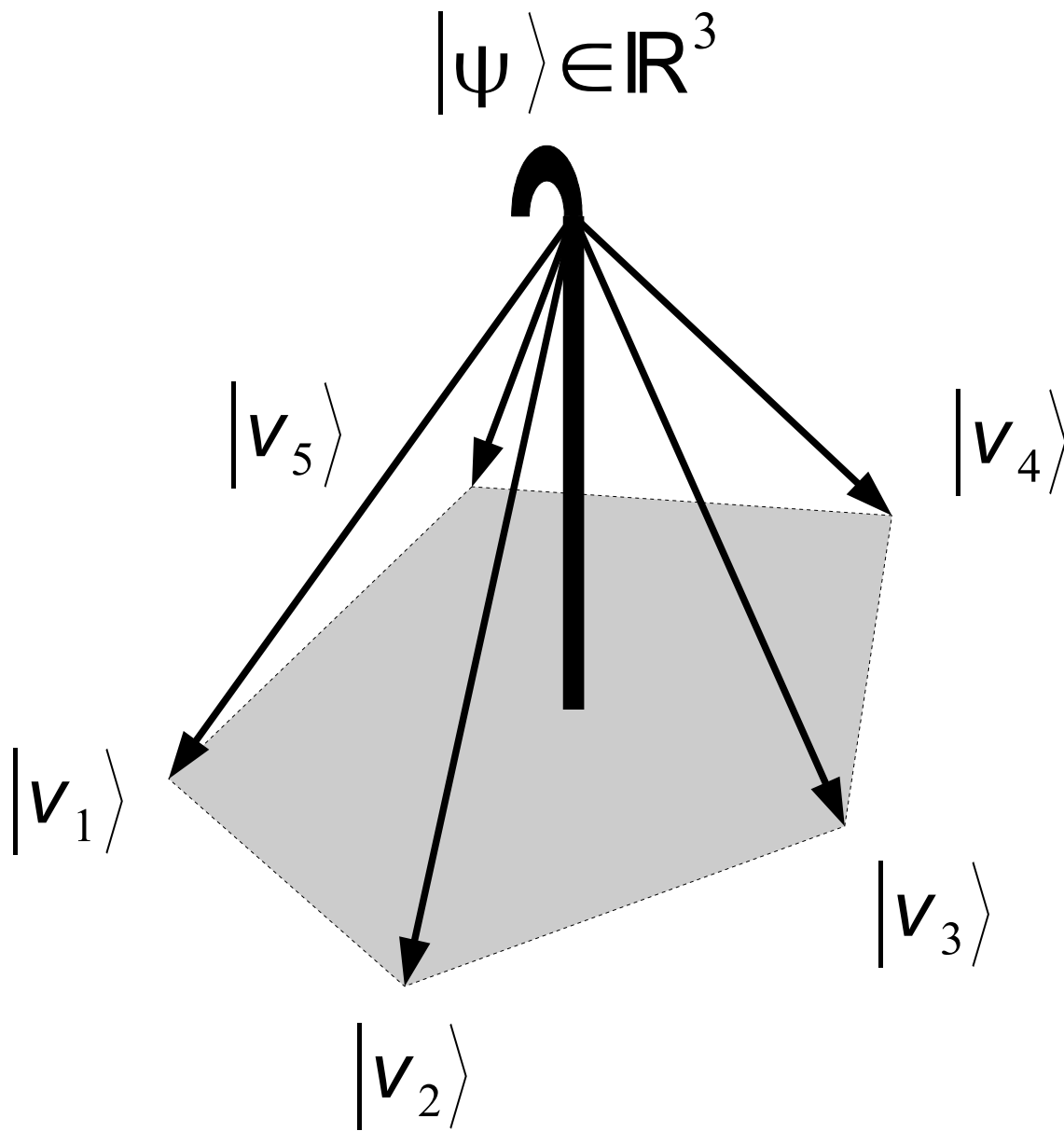


$$\langle v_i | v_{i+1 \bmod(5)} \rangle = 0$$

compatibility



$$|\langle \psi | v_i \rangle|^2 = \frac{1}{\sqrt{5}}$$

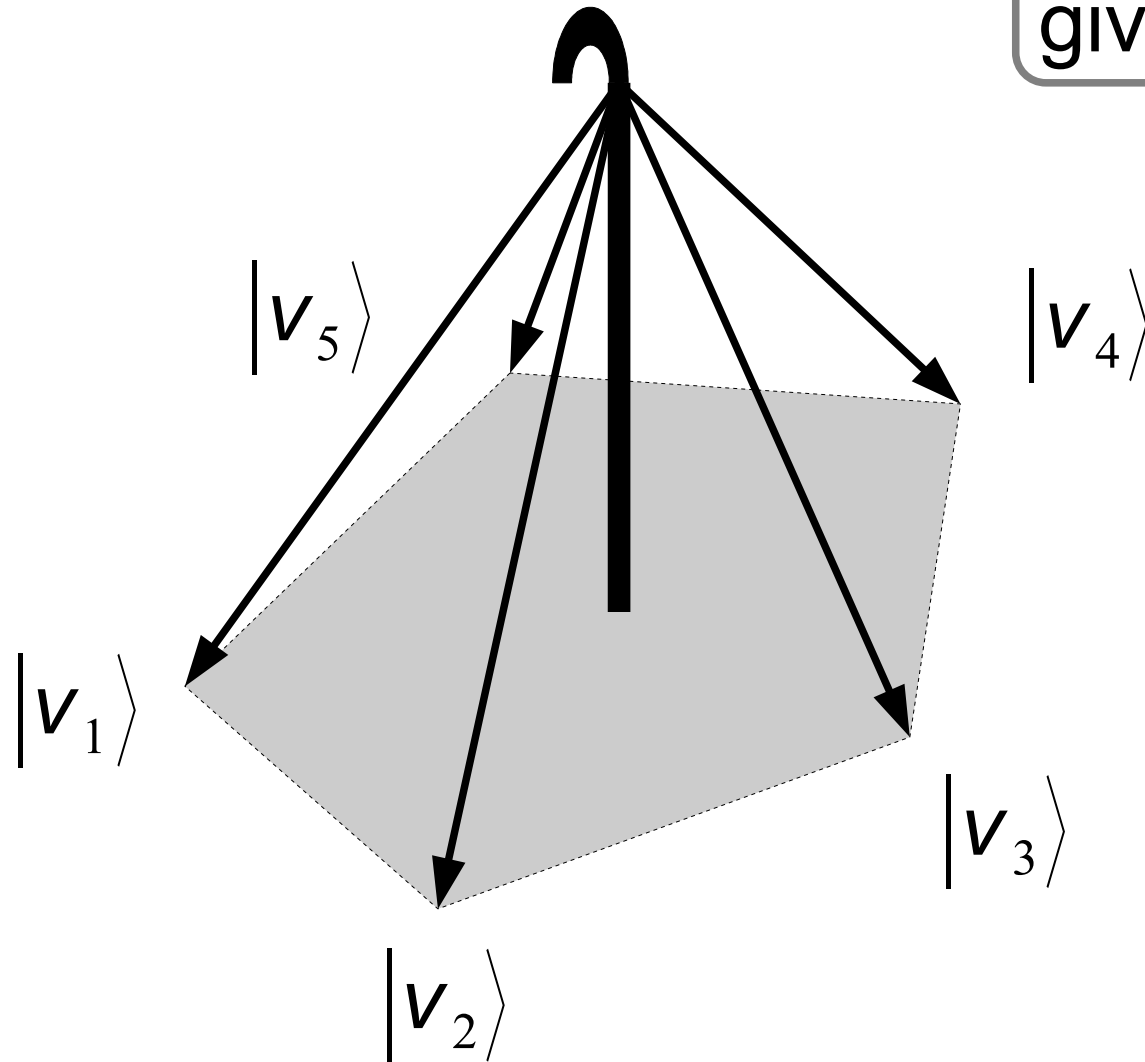


$$|\langle \psi | v_i \rangle|^2 = \frac{1}{\sqrt{5}}$$

$$\sum_{i=1}^5 |\langle \psi | v_i \rangle|^2 = \sqrt{5}$$

$$|\psi\rangle \in \mathbb{R}^3$$

quantum theory
gives expectation $\leq \sqrt{5}$



$$|\langle \psi | v_i \rangle|^2 = \frac{1}{\sqrt{5}}$$

$$\sum_{i=1}^5 |\langle \psi | v_i \rangle|^2 = \sqrt{5}$$

classical theories*
give expectation ≤ 2

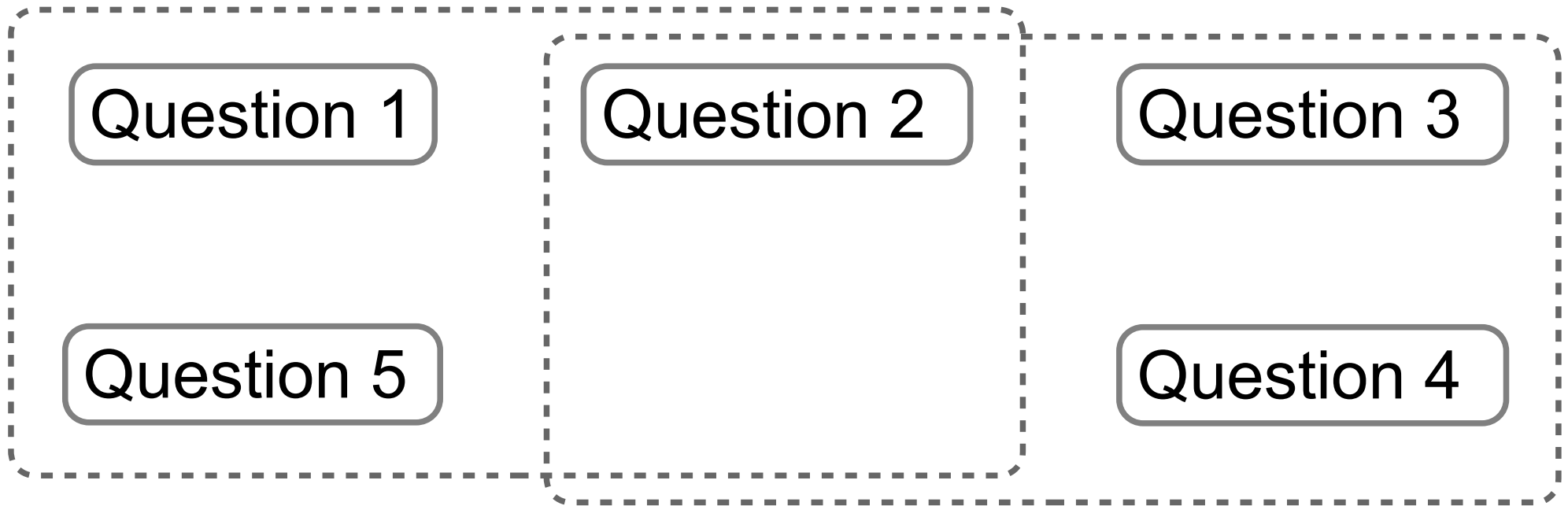
quantum theory**
gives expectation $\leq \sqrt{5} \approx 2.23$

non-signaling theories*
give expectation ≤ 2.5

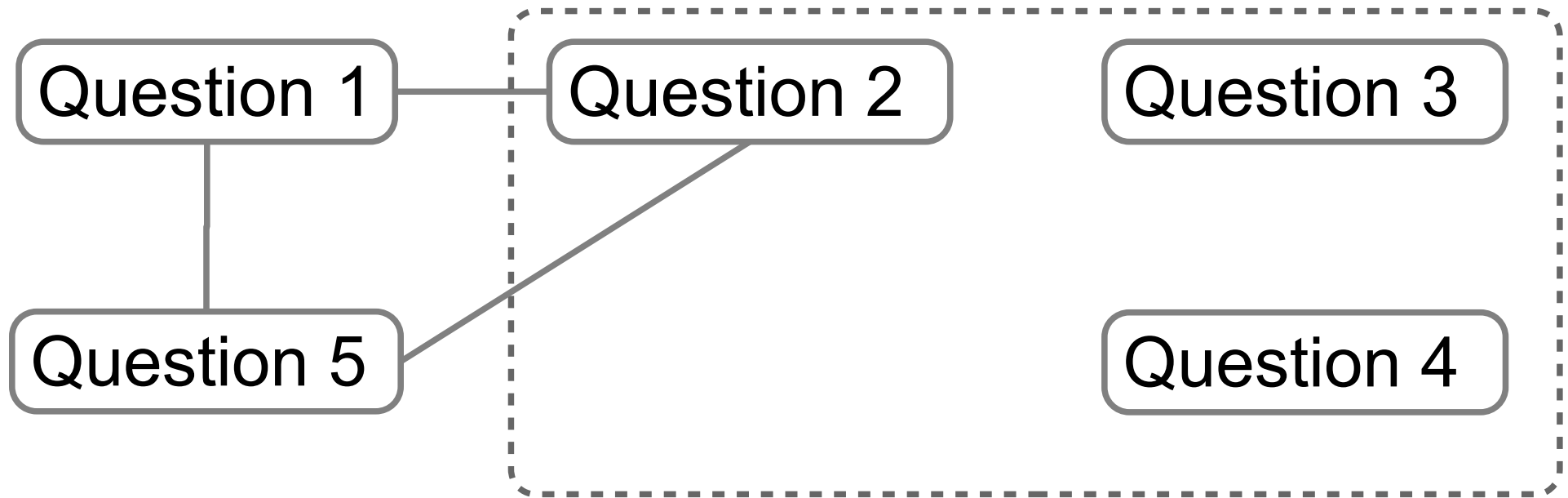
2. Results:

- a general framework
to study non-contextuality:
- 1. general compatibility structures
- 2. perfectness
- 3. non-locality

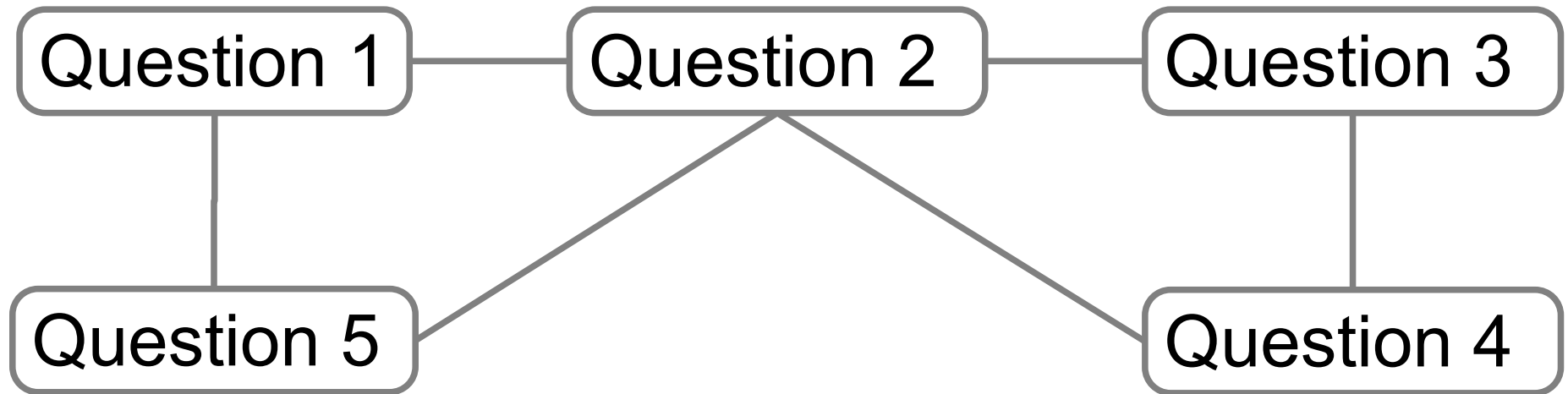
**every graph/hypergraph is
a compatibility structure**



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**every graph/hypergraph is
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Classification theorem

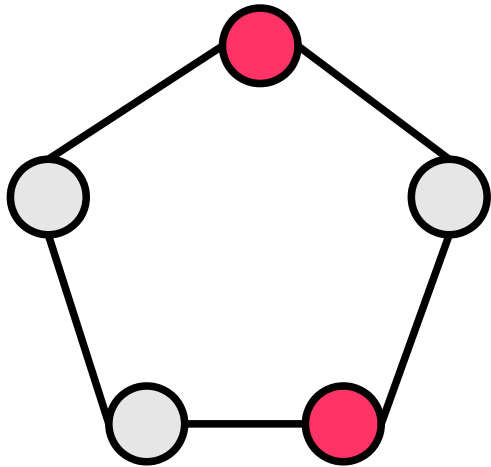
Let Γ be a compatibility structure, seen as a hypergraph. Let G be the graph obtained by connecting contextual questions. The maximum expectation values for exclusive answers are

$$\alpha(G) \leq \vartheta(G) \leq \alpha^{FP}(\Gamma)$$

for classical, quantum, and non-signaling theories, respectively; where $\alpha(G)$ is the independence number, $\vartheta(G)$ is the Lovász ϑ -function, and $\alpha^{FP}(\Gamma)$ is the fractional packing number.

Independence number $\alpha(G)$

Is the maximum number of mutually non-adjacent vertices in a graph G .

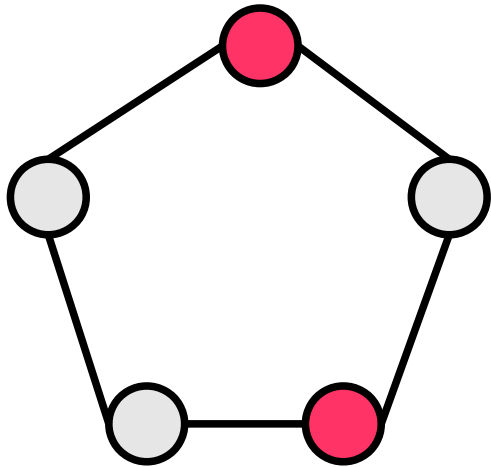


$$\alpha(C_5) = 2$$

NP-complete;
hard to approximate

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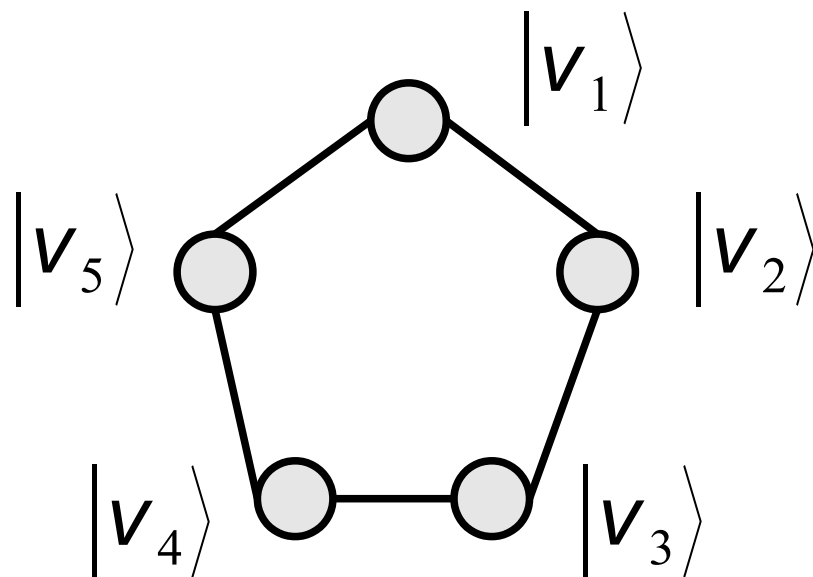
NP-complete;
hard to approximate

classical theories
give expectation ≤ 2

Lovász ϑ -function $\vartheta(G)$

An **orthogonal representation** of G is a set of unit vectors associated to the vertices such that two vectors are orthogonal if the corresponding vertices are adjacent:

$$\vartheta(G) := \max_{\text{orth. repr.}} \sum_{i=1}^n |\langle \psi | v_i \rangle|^2$$



$$\vartheta(C_5) = \sqrt{5}$$

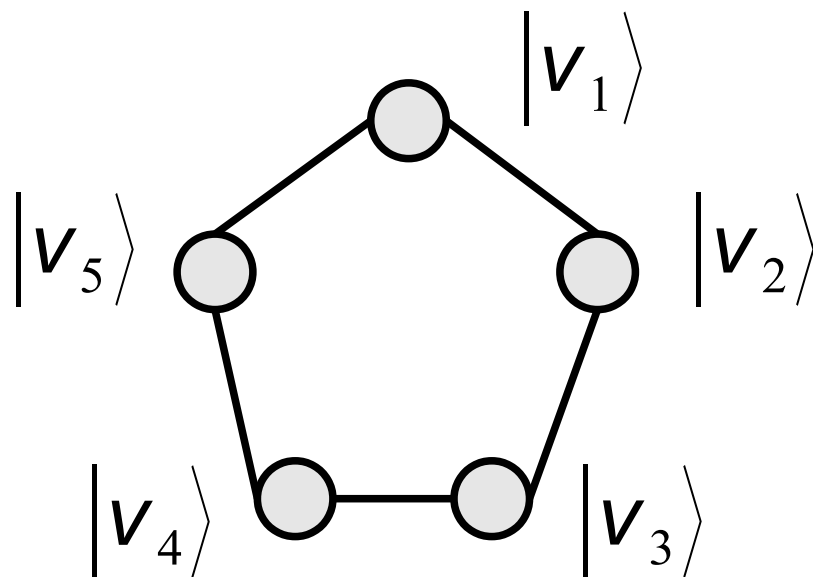
semidefinite program*

*Lovász (1978)

Lovász ϑ -function $\vartheta(G)$

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semidefinite program*

quantum theory
gives expectation $\leq \sqrt{5}$

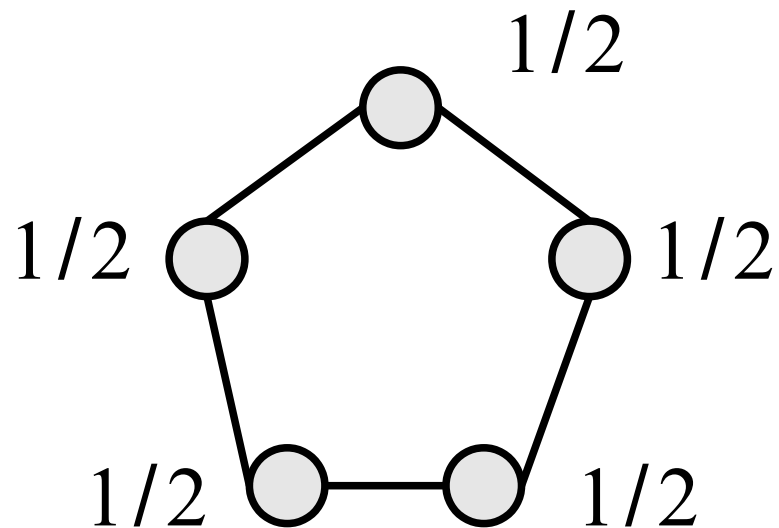
*Lovász (1978)

Fractional packing number $\alpha^{FP}(\Gamma)$

Let Γ be a compatibility structure, seen as a hypergraph:

$$\alpha^{FP}(\Gamma) = \max \sum_i w_i$$

$$\text{s.t. } \forall i \quad 0 \leq w_i \leq 1 \quad \text{and} \quad \forall \text{ context } C \in \Gamma, \quad \sum_{i \in C} w_i \leq 1$$



$$\alpha^{FP}(C_5) = 5/2$$

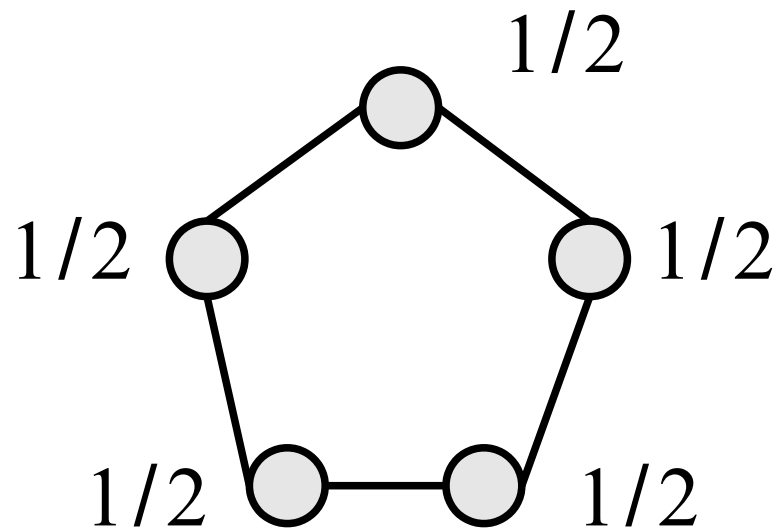
linear program

Fractional packing number $\alpha^{FP}(\Gamma)$

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$$\alpha^{FP}(C_5) = 5/2$$

linear program

non-signaling theories
give expectation ≤ 2.5

Remark.

$\wp(C_5)$ is the max. violation of the **Klyachko-Can-Binicioglu-Shumovsky (KCBS) inequality***

The inequality can be used to detect genuine quantum effects and it is the simplest non-contextual inequality violated by a qutrit (because the orthogonal representation has dimension 3).

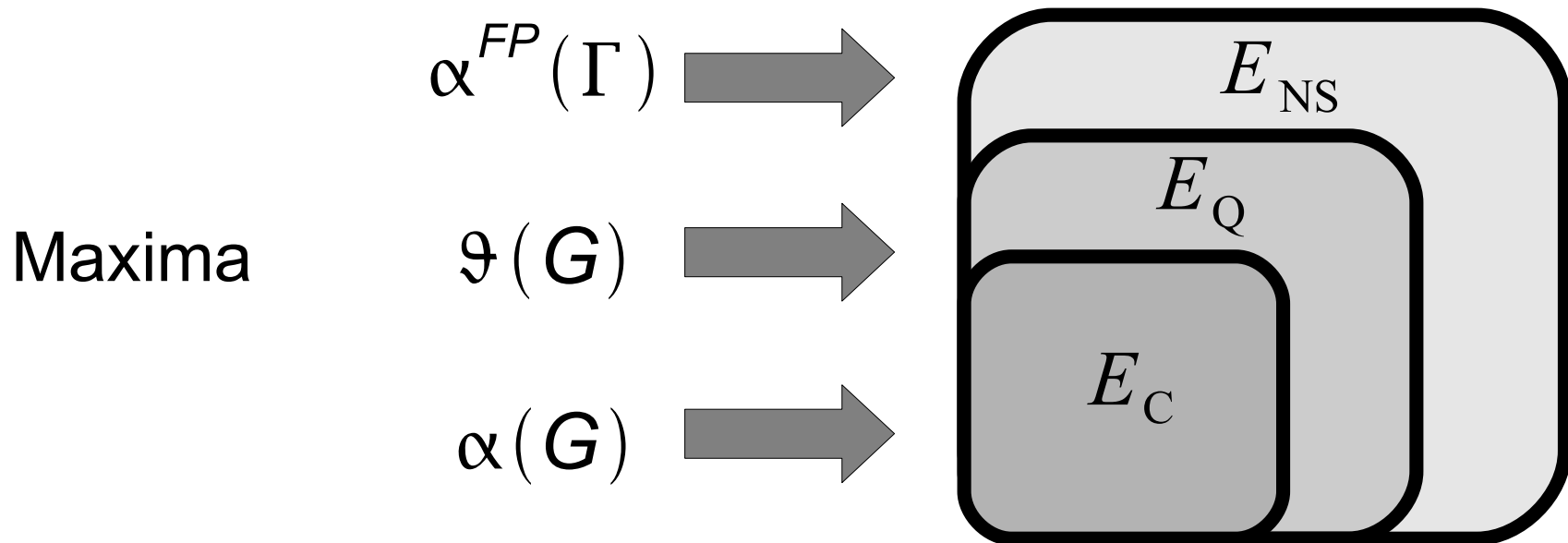
Quantum mechanics as a “sandwich theory”*

Let $E_C \subset E_Q \subset E_{NS}$ be the **convex sets** of the vectors realizing the expectations for classical, quantum, and non-signaling theories, respectively.

*cf. Knuth (1994)

Quantum mechanics as a “sandwich theory”*

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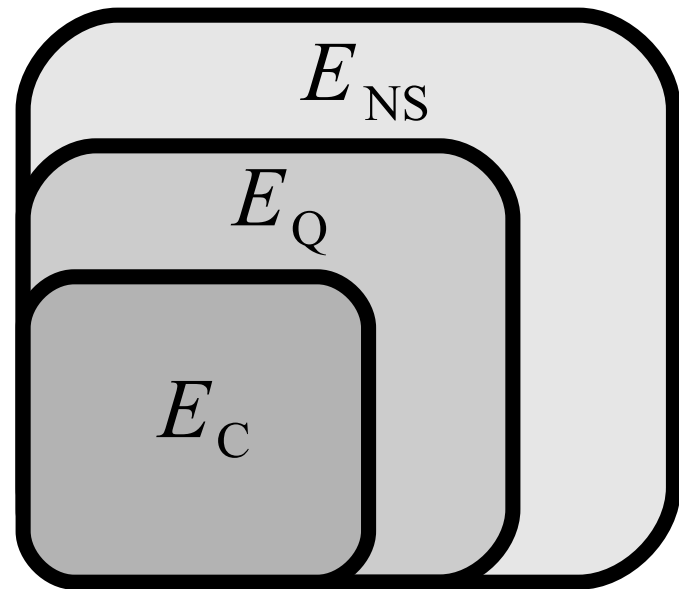


*cf. Knuth (1994)

Quantum mechanics as a “sandwich theory”*

Let $E_C \subset E_Q \subset E_{NS}$ be the **convex sets** of the vectors realizing the expectations for classical, quantum, and non-signaling theories, respectively.

membership in E_Q of a vector can be tested with a semidefinite program!



*cf. Knuth (1994)

Remark.

Standard result about the Lovász function can be then used to give the max. violation of known inequalities. For example, the max. violation for the n -cycle generalization of the KCBS inequality, recently computed in * is

$$\vartheta(C_n) = \frac{n \cos(\pi/n)}{(1 + \cos(\pi/n))}$$

Classicality and perfectness

A graph G is **perfect*** if $\alpha(H) = \vartheta(H) = \chi(\bar{H})$ for every induced subgraph H . So, perfect graphs are “the most classical ones”. For a perfect graph

$$E_C = E_Q = E_{NS}$$

Whenever $\alpha(G) < \vartheta(G)$ we have a difference between classical theories and quantum mechanics and a “state dependent” proof of the Kochen-Specker theorem**.

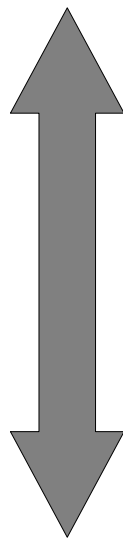
The KCBS inequality is based on C_5 which is the smallest non-perfect graph.

*Berge (1961); **where effects sum to unity

Two remarks

1. Many intractable problems are tractable for perfect graphs (*i.e.*, when classical and quantum theories coincide).
2. There are graphs s.t. $\alpha(G)=2$ and $\vartheta(G)=\Omega(n^{1/3})$ (*i.e.*, classical and quantum theories can have arbitrarily large separation)*.

non-contextuality



non-locality

Observation

Non-local experiments give compatibility structures:
compatible questions are the local measurement.

Observation

Non-local experiments give compatibility structures:
compatible questions are the local measurement.

Alice

settings $x \in X$

outcomes $a \in A$

Bob

settings $y \in Y$

outcomes $b \in B$

Compatibility graph for a non-local experiment

$$G = (V, E) \quad V = A \times B \times X \times Y$$

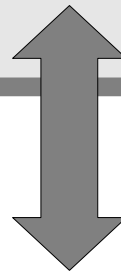
$$\{abxy, a'b'x'y\} \in E \text{ iff} \\ (x = x' \wedge a \neq a') \vee (y = y' \wedge b \neq b')$$

[Γ is the hypergraph of all cliques* in G]

Compatibility graph for a non-local experiment

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exclusiveness;

compatibility:

the observables of
Alice and Bob **commute**.

A classification theorem for correlations

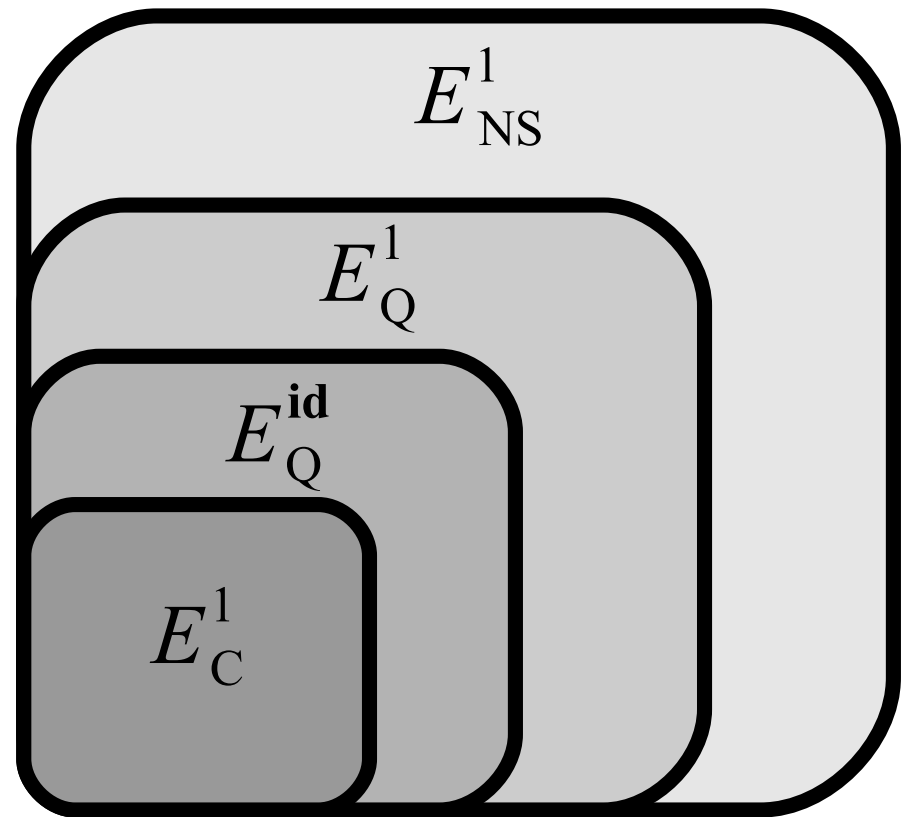
Let Γ be the compatibility hypergraph for a non-local experiment:

$$E_C^1(\Gamma) \subset E_Q^{\text{id}}(\Gamma) \subset E_Q^1(\Gamma) \subset E_{\text{NS}}^1(\Gamma)$$

are the sets of correlations obtainable by local hidden variables, local quantum measurements on a bipartite state, *idem* but without completeness relation for the measurement, and non-signaling theories, respectively:

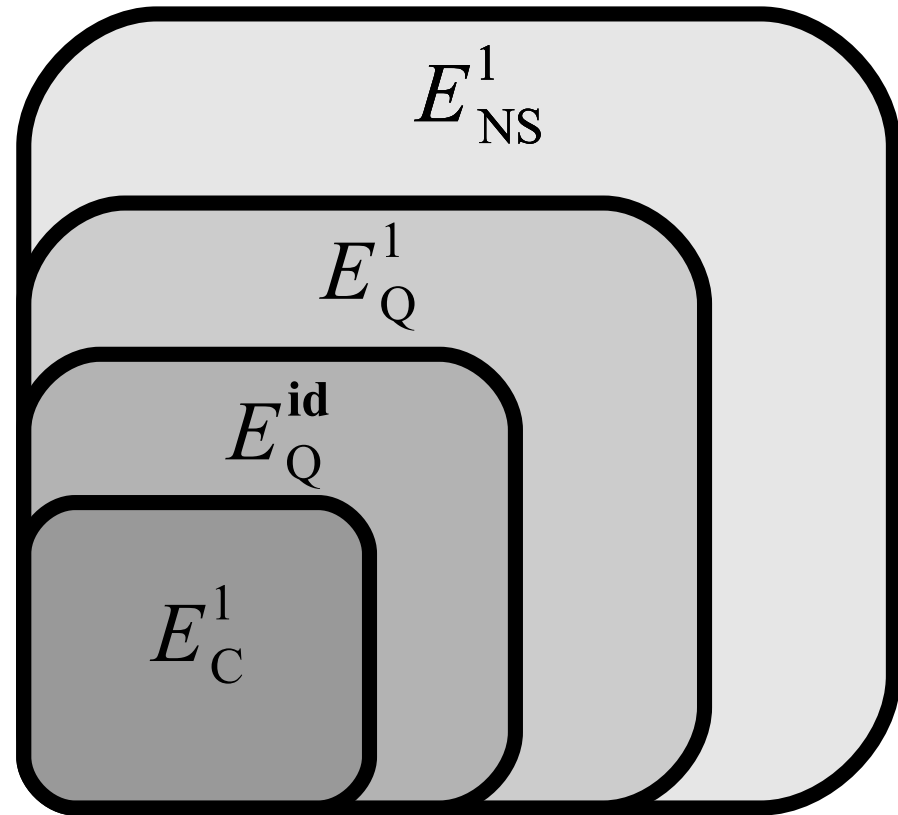
$$E_{X=C,Q,\text{NS}}^1(\Gamma) := E_X(\Gamma) \cap \{ \vec{w} : \forall xy \sum_{w_{ab|xy}} w_{ab|xy} = 1 \}$$

$$E_Q^{\text{id}}(\Gamma) := \{ (w_{ab|xy})_{abxy} : \forall xy \sum_{ab} P_{ab|xy} = \mathbf{id} \}$$



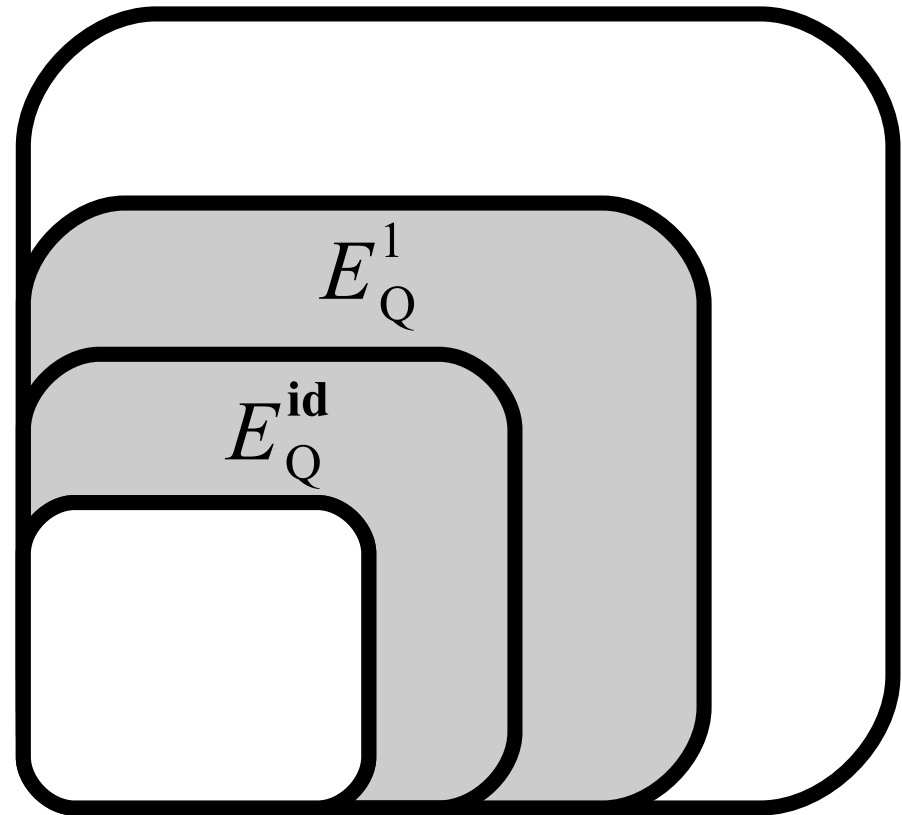
Fact

maximization over E_{NS}^1 and E_{C}^1
is equivalent to
maximization over E_{NS} and E_{C}



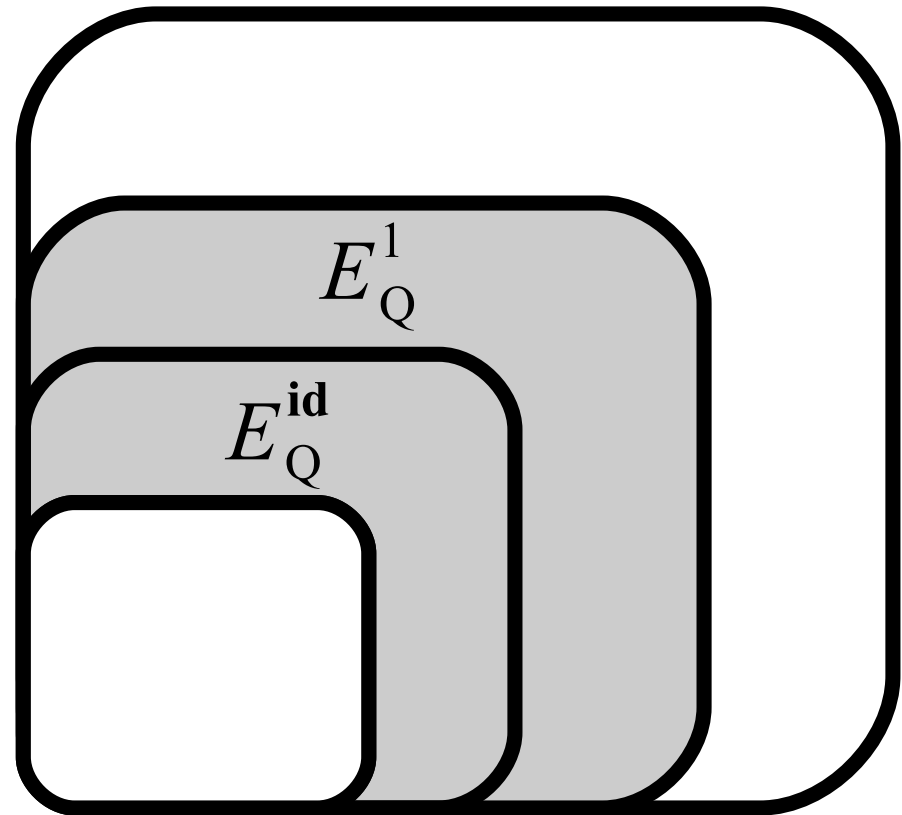
Fact

maxima via SDP
are efficient
upper bounds
to maximum
quantum violations



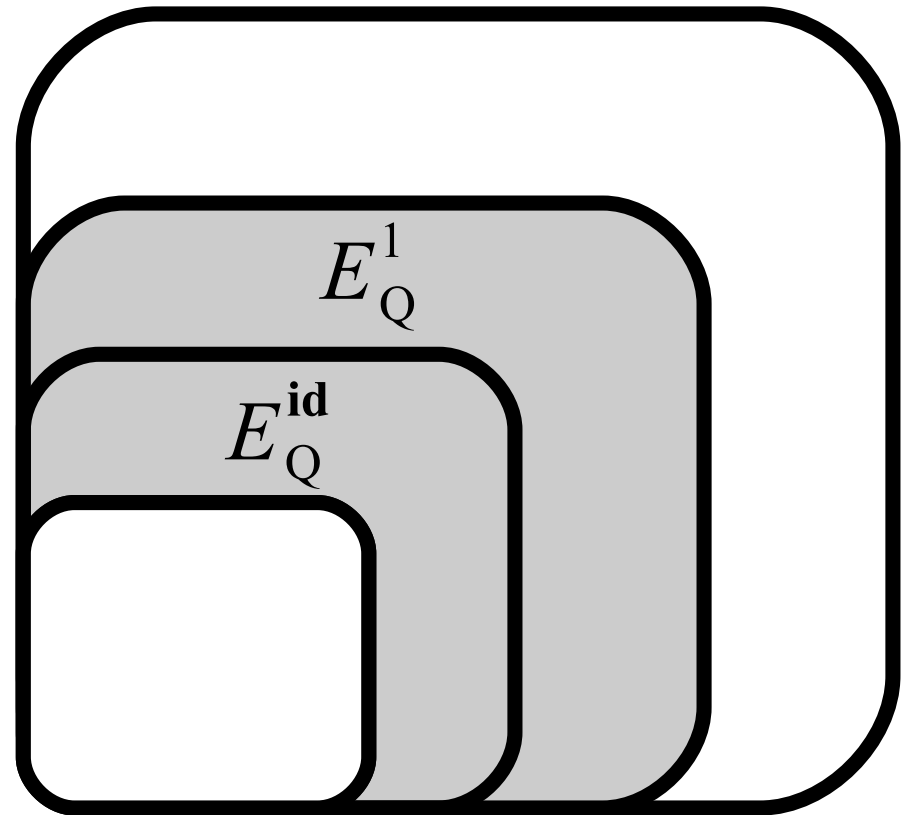
Fact

there is no efficient algorithm, unless the polynomial hierarchy collapses*



Fact

maxima via SDP
are efficient
upper bounds
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quantum violations



Problem: how well E_Q^1 approximates E_Q^{id} ?

Clause-Horne-Shimony-Holt (CHSH) inequality*

CHSH inequality

settings and outcomes: $A=B=X=Y=\{0,1\}$

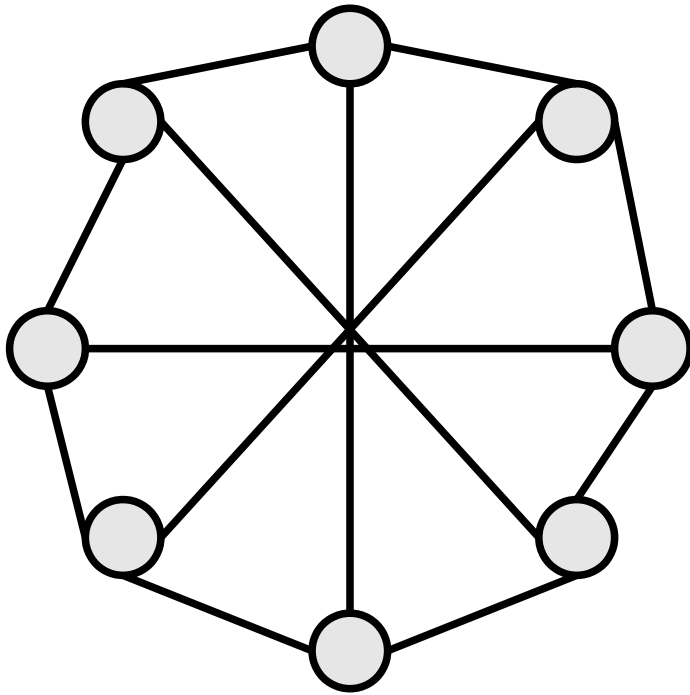
constraint:
$$\sum_{w_{ab|xy} : x \cdot y = a \text{ XOR } b} w_{ab|xy}$$

CHSH inequality

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constraint:

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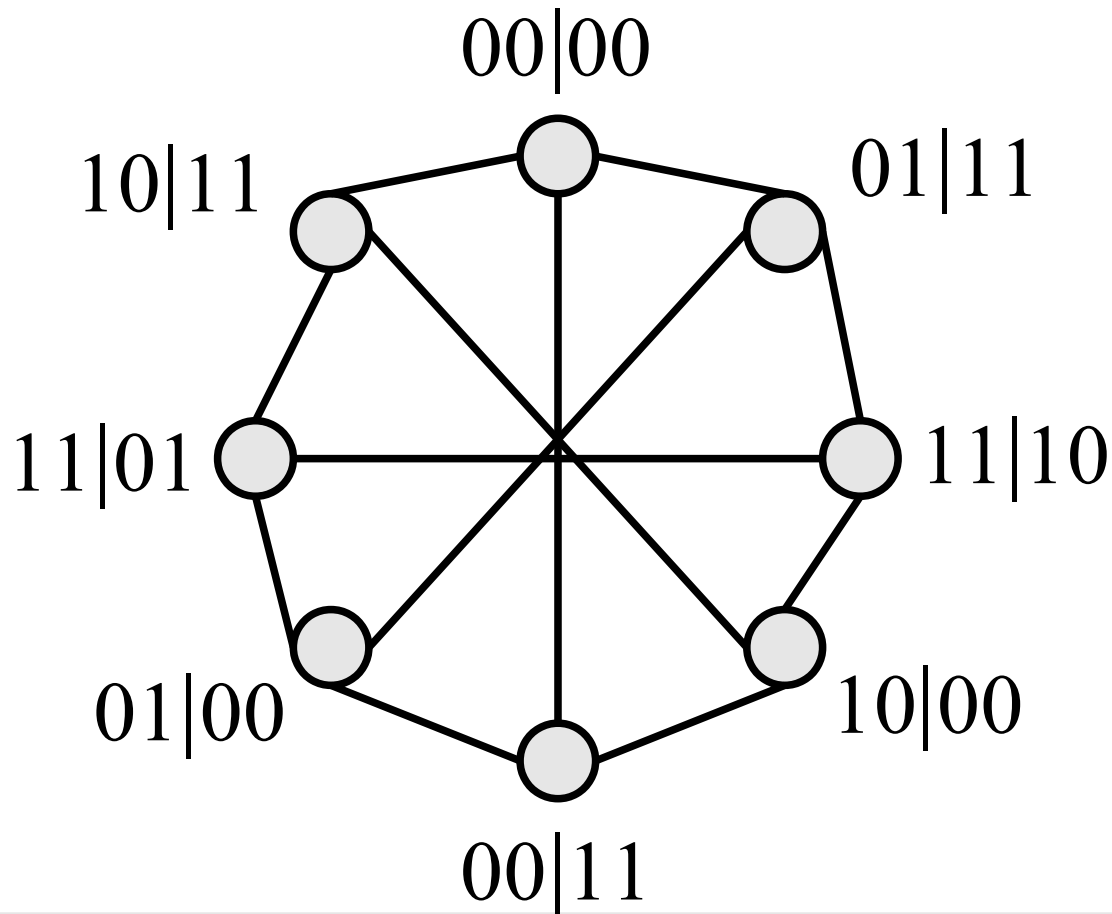


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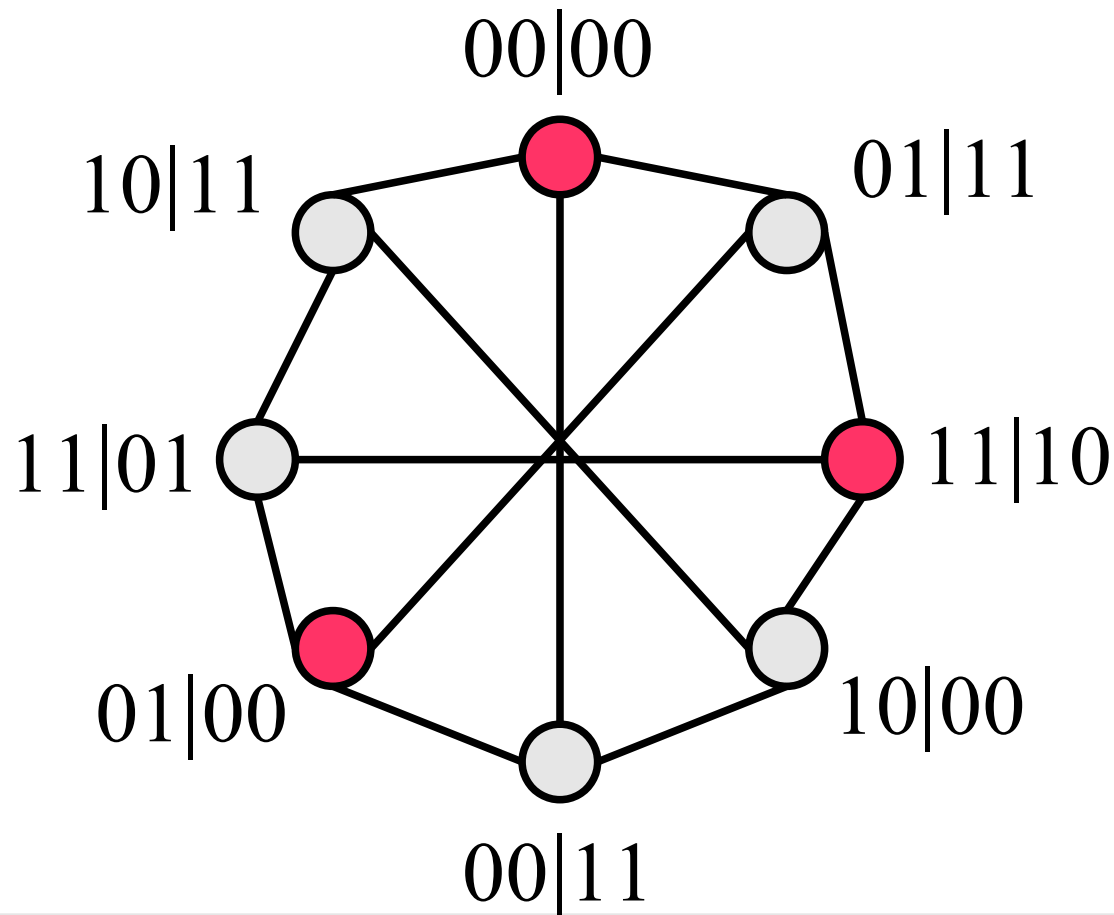


CHSH inequality

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$$\alpha(G)=3$$

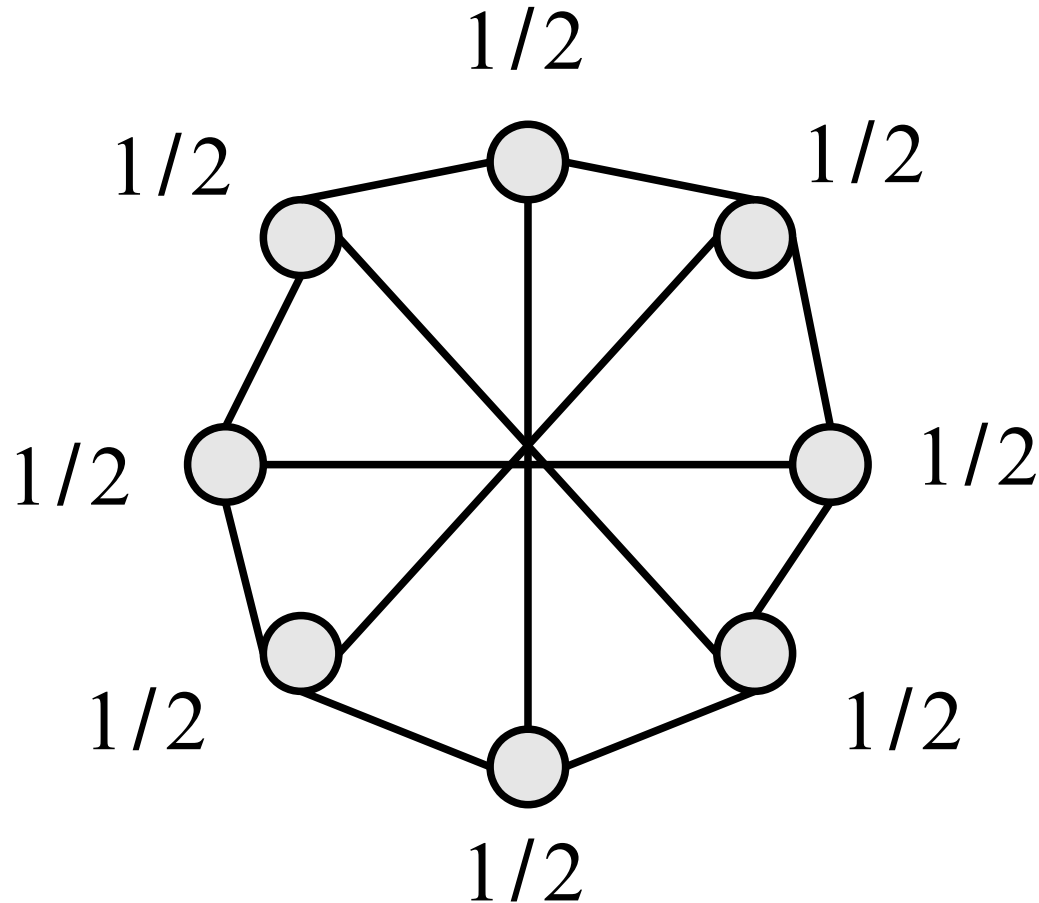
classical max.

CHSH inequality

settings and outcomes: $A=B=X=Y=\{0,1\}$

constraint:

$$\sum_{w_{ab|xy} : x \cdot y = a \text{ XOR } b} w_{ab|xy}$$



$$\alpha(G)=3$$

classical max.

$$\alpha^{\text{FP}}(\Gamma)=4$$

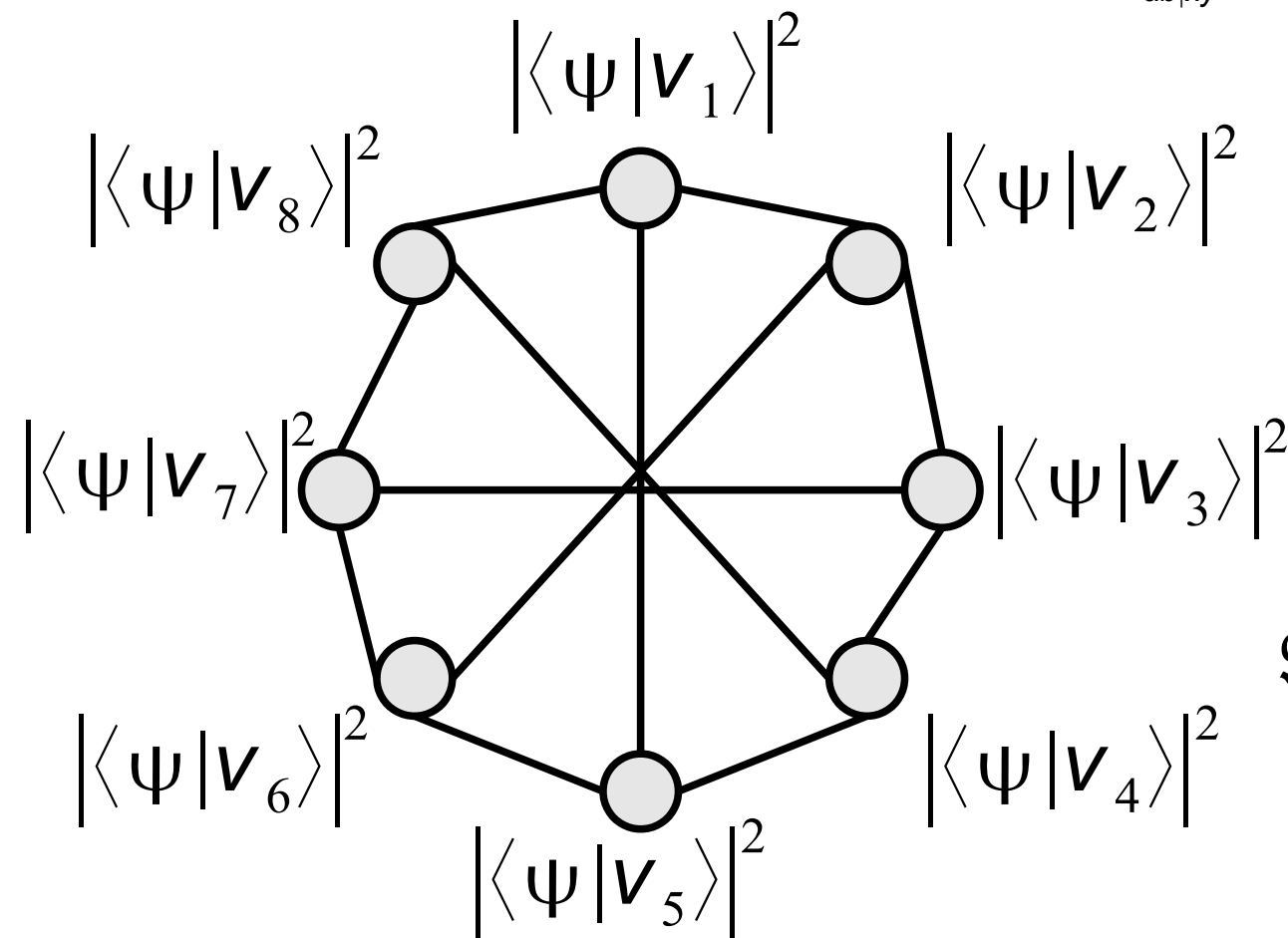
non-signaling max.

CHSH inequality

settings and outcomes: $A=B=X=Y=\{0,1\}$

constraint:

$$\sum_{w_{ab|xy} : x \cdot y = a \text{ XOR } b} w_{ab|xy}$$



$$\alpha(G)=3$$

classical max.

$$\alpha^{\text{FP}}(\Gamma)=4$$

non-signaling max.

$$\vartheta(G)=2+\sqrt{2}\approx 3.4$$

quantum max.

CHSH inequality

$$\alpha(G) = 3$$

classical max.

$$\alpha^{\text{FP}}(G) = 4$$

non-signaling max.

it attains the
Tsirelson bound



$$\vartheta(G) = 2 + \sqrt{2} \approx 3.4$$

quantum max.

Collins-Gisin inequality (I3322)*

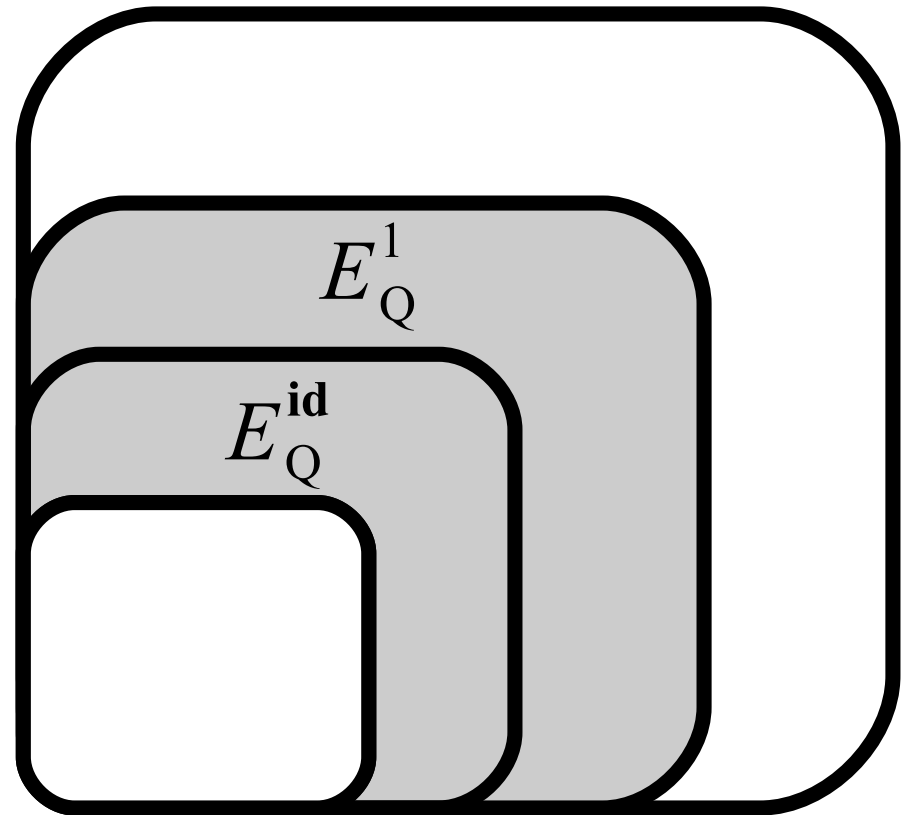
max. 6.2514



max. 6.2508**



$$E_Q^{\text{id}} \subset E_Q^1$$



3. Open problems:

1. theoretical:
relations to Bell inequalities
2. applied:
loophole-free experiments
3. a complexity perspective:
degree of perfectness

“theoretical open problem”

Can any violation of a non-contextual inequality be converted into a (comparably large) violation of a Bell inequality?

“applied open problem”

Can any violation of a non-contextual inequality be converted into a (comparably large) violation of a Bell inequality?

So far, forty years after Bell paper, all Bell experiments have loopholes: are graphs with a large separation between the independence number and the Lovász function good candidates for loophole-free experiments with inefficient detectors?

“complexity open problem”

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Perfect graphs have many efficient algorithms that in general are NP-hard. We have shown that compatibility structures from perfect graphs have coincident classical and quantum description. Can we define a notion of parametric complexity according to the classical-quantum gap?

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