4(b) - Atomic lifetimes Atom-Photon Physics

Let us consider N(t) atoms in an excited state b at a particular time t. The rate of change of N(t) is



$$N(t) = N(t=0) \exp\left(-\frac{t}{2b}\right)$$

where $\overline{z}_{b}^{-1} = \underbrace{\leq}_{k} \underset{k \in b}{ } W_{kb}^{\leq}$ is the lifetime of level b

4(b) - Atomic lifetimes Atom-Photon Physics

Please note:

• In the absence of external fields, the lifetime of an atom cannot depend on the magnetic quantum number m of the level b (atoms = spherically symmetric systems)

• The lifetimes of hydrogenic ions are shorter than those in Hydrogen. They are given by the scaling law

$$Z(Z) = Z^{-4}Z(Z=1)$$

In the dipole approximation the lifetime of 2s in H is infinite
In practice,

$$(2S) \xrightarrow{\sqrt{1}}{7} S$$

(two-photon emission)



5- Measurements (radiative lifetimes)

- Importance of oscillator strengths (f values) of spectral lines: computations of gas discharges, plasmas, stellar atmosphere, etc
- In general, theoretical predictions are not accurate: experimental measurements are necessary

Direct measurements: f can be determined through the intensity of the spontaneous emission of a given spectral line.



Problem: in many cases, N cannot be determined accurately Solution: f is determined by measuring radiative lifetimes

5(a) – The beam-foil method



Key idea: "The time dependent exponential decay of excited atoms may be converted into a spatial variation of intensity by exciting a beam of fast-moving atoms at a given position".



Problem: repopulation of excited states by radiative decay from higher levels: too long lifetimes (overcome by Koenig+Ellett(1932) using optical excitation) and low velocity of the beam (t > 10⁻⁶ s)

5(a) – The beam-foil method



Modern technique (Kay, 1963; Bashkin, 1964):

- Based on the fact that ions in the beam of a van der Graaf accelerator can be strongly excited by passing through a thin carbon foil
- Advantage: high velocity (10⁸cm/s): radiative decay extended over several cms

1. Acceleration of ionsinitial v spread determined by analyzing magnet

2. Ions collide with the foil: excitation

3. Ions decay: intensity/wavelength of emitted radiation is detected



4. Beam velocity is callibrated: one must take into account the energy loss of the ions passing through the foil

5(a) – The beam-foil method



Calculation of lifetimes

$$\kappa(l) = N_{\kappa}(0) \exp(-t/2k) t = e/v$$

Density of ions at the downstream face of the foil at position 1

$$ln Nk(l)$$
 as a function of *l*: slope gives T_k

Problem: radiative cascade: the plot shows a pronounced curvature due to the repopulation of the level being studied by the decay from higher-lying levels

- Remember: the ions are excited by collisions and this is an unselective process
- The radiative cascade can be avoided by replacing the foil by an intense laser: strict selection rules

5.(b) – The delayed coincidence method



5.(c) – The time of flight method

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van Dyck et al, Phys. Rev. Lett. 25, 1403 (1970)

- Importance: this method is widely used for determining the lifetimes of metastable states (typically >10⁻³s)
- Problems with beam-foil/delayed coincidence method:
 - Beam foil method: no decay would be observed for an apparatus of realistic size
 - Delayed coincidence method: the time necessary for obtaining sufficient data is unrealistic

Example: metastable states in He, He⁺

5.(c) – The time of flight method





Lamp: if there are 2 close metastable levels, it excites the atom to a close, non metastable level

Why "time of flight"?

One may determine the number of atoms which decay between both detectors by comparing the number of metastable atoms arriving at the detectors within specific velocity intervals



If the initial velocity distribution is uniform across the beam

$$N_i(v) = \left(\begin{array}{c} \varepsilon & \varepsilon \\ \varepsilon &$$

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Ratio (numbers of atoms at both detectors)

$$R = \frac{N_{D}(N)}{N_{A}(N)} = \frac{C_{D}}{C_{A}} \exp(-t/2)$$