

# A systematically reduced mathematical model for organoid expansion

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## 2 ABSTRACT

1

Organoids are three-dimensional multicellular tissue constructs. When cultured *in vitro*, they 3 4 recapitulate the structure, heterogeneity, and function of their in vivo counterparts. As awareness of the multiple uses of organoids has grown, e.g. in drug discovery and personalised medicine, 5 demand has increased for low-cost and efficient methods of producing them in a reproducible 6 manner and at scale. Here we focus on a bioreactor technology for organoid production, which 7 exploits fluid flow to enhance mass transport to and from the organoids. To ensure large numbers 8 of organoids can be grown within the bioreactor in a reproducible manner, nutrient delivery to, 9 and waste product removal from, the organoids must be carefully controlled. 10

We develop a continuum mathematical model to investigate how mass transport within the 11 bioreactor depends on the inlet flow rate and cell seeding density, focusing on the transport 12 13 of two key metabolites: glucose and lactate. We exploit the thin geometry of the bioreactor to systematically simplify our model. This significantly reduces the computational cost of generating 14 model solutions, and provides insight into the dominant mass transport mechanisms. We test 15 the validity of the reduced models by comparison with simulations of the full model. We then 16 exploit our reduced mathematical model to determine, for a given inlet flow rate and cell seeding 17 density, the evolution of the spatial metabolite distributions throughout the bioreactor. To assess 18 the bioreactor transport characteristics, we introduce metrics quantifying glucose conversion 19 (the ratio between the total amounts of consumed and supplied glucose), the maximum lactate 20 concentration, the proportion of the bioreactor with intolerable lactate concentrations, and the time 21 when intolerable lactate concentrations are first experienced within the bioreactor. We determine 22 the dependence of these metrics on organoid-line characteristics such as proliferation rate and 23 24 rate of glucose consumption per cell. Finally, for a given organoid line, we determine how the distribution of metabolites and the associated metrics depend on the inlet flow rate. Insights from 25 26 this study can be used to inform bioreactor operating conditions, ultimately improving the quality and number of bioreactor-expanded organoids. 27

28 Keywords: organoid culture, bioreactor, asymptotic, multiscale, transport, reduced-order model

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## **1 INTRODUCTION**

Organoid technology is becoming increasingly prominent as a biomedical tool, with applications in drug discovery and personalised medicine. In biomedical research, brain, kidney, and liver organoids are used to understand the underlying biological mechanisms in tissue development and tissue–drug interactions (Bock et al., 2020; Eisenstein, 2018; Kondo and Inoue, 2019; Tuveson and Clevers, 2019).

Organoids are three-dimensional, multicellular structures which, when grown in vitro, recapitulate 36 the structure, function, and heterogeneous cellular composition of in vivo tissues (Drost and Clevers, 37 2018). Their three-dimensional geometry means they are more representative of in vivo tissues than 2D 38 cell cultures (Young and Reed, 2016). "Organoid expansion" refers to the growth of multiple organoids 39 from pluripotent stem cells, which are typically derived from patient biopsies or from other organoids 40 (de Souza, 2018). The stem cells are embedded in a supporting extra-cellular matrix (ECM) and cultured in 41 carefully-controlled conditions designed to promote organoid growth. The surrounding ECM provides the 42 biochemical and biomechanical cues needed for the cells to proliferate and differentiate into specialised 43 cells, as happens in vivo (Eisenstein, 2018; Huang et al., 2012). 44

Current methods for organoid expansion are labour intensive, with organoids typically being produced 45 in small numbers at specialist research laboratories. New technologies are required to manufacture large 46 47 numbers of organoids with uniform and reproducible characteristics, to meet the demands of applications 48 such as high-throughput screening in drug development. One such technology exploits bioreactors, 49 which aim to deliver sufficient nutrients and growth factors to the cells to promote cell proliferation and differentiation, and to prevent the accumulation of toxins, which can lead to cell death. For a more 50 detailed overview of bioreactor technologies used for 3D cell culture see, for example, Martin et al. (2004), 51 Pörtner and Giese (2006) and Wendt et al. (2009). 52

This study is motivated by proprietary organoid expansion bioreactor technology developed by Cellesce 53 (Ellis et al., 2019). The 'Cellesce Expansion 1 (CXP1)' bioreactor is currently used to expand colorectal 54 cancer organoids, see Figure 1. Flow of media through the system enhances the delivery of nutrients to, 55 and the removal of waste products from, organoids seeded in a hydrogel layer. In this application, oxygen 56 is present at high concentrations, and is not a limiting factor for organoid growth. The key metabolites of 57 interest here are glucose, essential for colorectal cancer organoid growth, and lactate. Lactate can have a 58 detrimental effect on cell behaviour, such as metabolism (Romero-Garcia et al., 2016), and sufficiently 59 60 high levels can lead to cell death. Lactate can be produced via anaerobic respiration and aerobic glycolysis (Liberti and Locasale, 2016). We do not focus on the precise mechanisms of lactate production here, but 61 instead determine how the media flow promotes lactate removal. We note that while colorectal cancer 62 organoids tolerate high lactate concentrations, the intention is to use CXP1 to expand a range of normal 63 (healthy) and pathological organoids. Since different organoid types have distinct requirements (e.g. nutrient 64 levels required for cell proliferation and lactate tolerances), understanding the mass transport of glucose 65 and lactate within the bioreactor is important. While we acknowledge the biological complexity of organoid 66 culture, spatiotemporal knowledge of these two metabolites provide useful and practical information on the 67 operation of CXP1, and provides the framework for more complex models in the future. 68

Key priorities in the CXP1 bioreactor design and operation are uniformity of organoid size and system
 reproducibility, to ensure there is minimal variation in organoid characteristics between and within batches

71 grown under the same operating conditions. The main control parameters for the CXP1 bioreactor are the 72 inlet flow rate (controlled via a peristaltic pump) and the initial cell seeding density (the organoids are 73 grown from single cells). Optimisation of these control parameters requires spatiotemporal information 74 about the flow and metabolite (here glucose and lactate) concentrations throughout the bioreactor (Galban 75 and Locke, 1999a). Such data are impractical, inefficient, and expensive to collect through experimental 76 means alone.

77 To complement experimental studies, mathematical models of bioreactor systems can be used to predict 78 media flow profiles and the associated metabolite concentrations that cannot easily be measured in vitro, 79 thus providing useful insights to ensure CXP1 operation is maintained within tolerable operating regions 80 of these metabolites. Here, we adopt a continuum modelling approach, in which the dependent variables 81 (cell density, fluid velocity, metabolite concentrations) are assumed to vary continuously in space and time. Our resulting model comprises a system of partial differential equations (PDEs). A key advantage of 82 83 such a mathematical modelling approach is the ability to quickly, efficiently and accurately analyse the 84 system as control parameters are varied. A continuum, rather than discrete, cell-based approach is often used to model bioreactor systems, which is justified due to the typical cell numbers ( $\mathcal{O}(10^6)$  cells) and 85 metabolite concentrations (CXP1: 16mM in 15mL of culture media) present. We model the organoids (cell 86 87 aggregates) as effective (bulk) reaction terms over the hydrogel, which can be formally obtained through an asymptotic homogenisation procedure (see, for example, Dalwadi et al. (2018); Dalwadi and King (2020)). 88

89 Here we review existing mathematical models for metabolite transport in bioreactor systems. A variety of 90 different mathematical modelling approaches have been applied to related problems in tissue engineering, 91 including: ordinary differential equation (ODE) models (Sachs et al., 2001); PDE models (Galban and Locke, 1999b,a; Shipley et al., 2009, 2011; Shipley and Waters, 2012; Chapman et al., 2014, 2017; Pearson 92 et al., 2014); computational approaches (Mehrian et al., 2020b; Nguyen et al., 2018); and agent-based 93 models (Drasdo and Höhme, 2005; Byrne et al., 2007; Byrne and Drasdo, 2009). For a more comprehensive 94 review of continuum modelling approaches for tissue engineering, see O'Dea et al. (2012). As noted above, 95 in this work we use a continuum modelling approach to develop a PDE model for metabolite transport 96 within a specific bioreactor set-up. We focus on a systematic model reduction of this model, taking an 97 approach similar to that used in Shipley et al. (2011); Shipley and Waters (2012); Chapman et al. (2017). In 98 so doing, we highlight two key advantages of model reduction. First, we identify the physical mechanisms 99 that dominate the system behaviour on the timescale of interest. Secondly, reduced models are more 100 tractable than their full model counterpart and, as such, can be solved more rapidly numerically or, in some 101 102 cases, analytically. This facilitates more detailed exploration of parameter space, which is important for 103 subsequent optimisation of bioreactor operating conditions, and allows more detailed biological models to 104 be incorporated.

105 We develop a mathematical model of the CXP1 system, with the goal of determining how glucose and lactate levels within the CXP1 bioreactor change as the operating conditions (e.g. media inlet flow rate 106 107 and cell seeding density), and organoid growth characteristics, vary. We introduce a reaction-advection-108 diffusion system for glucose and lactate transport in the CXP1 bioreactor. The hydrogel and media are viewed as two distinct regions, coupled by interfacial conditions. We restrict attention to a two-dimensional 109 slice through the bioreactor, and obtain numerical solutions to the governing equations. Motivated by 110 typical parameter values of the bioreactor, we perform an asymptotic analysis to systematically reduce 111 the model from a two-dimensional geometry to a one-dimensional model, in which vertically-averaged 112 concentration profiles vary with horizontal position along the length of the bioreactor. We validate this 113 reduced model through successful comparisons with numerical solutions of the full system. We exploit the 114

reduced models to explore the parameter space of cell characteristics and bioreactor operating regimes. 115 116 To assess glucose and lactate levels, we introduce the following quantitative, time-dependent metrics: glucose conversion (the ratio between the total amounts of consumed and supplied glucose); maximum 117 lactate concentration within the bioreactor; proportion of domain with intolerable lactate levels (i.e. lactate 118 levels above a tolerated concentration); and time when intolerable lactate levels are first experienced. 119 For a given organoid type, we determine how these metrics change as the inlet flow rate varies. In this 120 way, we aim to show how quantitative insights gained from this modelling approach can inform the 121 selection of experimental bioreactor operating conditions, and ultimately improve the quality and quantity 122 of bioreactor-expanded organoids. 123

The structure of the paper is as follows. In the Methods section, we introduce the full mathematical 124 model, and then systematically derive two reduced models (referred to as the longwave approximation and 125 the *sublimit approximation*) for glucose and lactate transport within the bioreactor. In the Results section, 126 127 we verify that simulations of the reduced models are in good agreement with solutions of the full model for physiologically relevant parameter regimes. We demonstrate the advantages of the model reductions, 128 highlighting, in particular, the physical insights obtained from systematic derivation of the reduced models 129 from the full system. We then use the longwave approximation model to investigate how the glucose and 130 lactate concentrations within the bioreactor change for different organoid lines. We examine the evolution 131 of the concentration profiles and demonstrate how our quantitative metrics to assess metabolite behaviour 132 are heavily dependent on organoid line characteristics, such as proliferation and nutrient consumption 133 rates. We then investigate, for a specific organoid line, how the media inlet flow rate affects the metabolite 134 135 concentrations, and explain how this information can be used to optimise the bioreactor control parameters. The paper concludes with a Discussion where we summarise our results and outline future directions for 136 our modelling approach. 137

# 2 METHODS

We derive an unsteady two-dimensional model for glucose and lactate transport within the CXP1 bioreactor.
Schematics of the CXP1 bioreactor and our model geometry are presented in Figure 1. We use COMSOL
Multiphysics® to solve the full mathematical model numerically and use the insights provided by the
numerical simulations to motivate systematic reductions of the full model. The resulting reduced models
are solved using a combination of analytical (method of characteristics) and numerical (Chebfun toolbox
and ode45 in MATLAB) techniques.

## 144 2.1 Bioreactor set-up

We consider organoids grown from single cells seeded in a homogeneous thin layer of hydrogel in the 145 bioreactor (lower yellow layer in Figure 1). A typical initial seeding density for the CXP1 bioreactor is 146  $4 \times 10^5$  cell mL<sup>-1</sup> –  $6 \times 10^5$  cell mL<sup>-1</sup>. We assume that all cells seeded within the hydrogel are viable and 147 become organoids, and that there is negligible settling (which is a fair assumption given the relative time 148 of the gelation of the well-mixed solution, compared to the settling time of the cells). The hydrogel acts 149 as a porous scaffold for the seeded cells, providing the anchorage for cells and the biomechanical and 150 biochemical cues required for cell growth (Huang et al., 2012). The bioreactor is placed within an incubator 151 which maintains constant temperature, O<sub>2</sub> (atmospheric levels) and CO<sub>2</sub> concentration. Nutrient-rich 152 culture media, with typical glucose concentration of 16mM, is stored in an upstream reservoir and is fed 153 into the system through an inlet pipe, and slowly flows across the bioreactor (upper blue layer in Figure 1), 154 with typical flow velocity of  $10^{-6}$ m s<sup>-1</sup>. The media is then removed from the bioreactor through an outlet 155 pipe. The top of the culture media layer is a free surface. We assume there is no flow within the hydrogel. 156 We consider colorectal cancer organoids, which are expanded in the bioreactor for 7 days. The organoids 157

158 are grown from single stem cells (roughly  $10\mu$ m in diameter) until they are approximately  $40 - 80\mu$ m in 159 diameter and comprise approximately 50 cells. The organoids are then extracted from the hydrogel and 160 tested for size, viability, and number of cells per organoid. The total number of organoids per bioreactor 161 is also recorded. Finally, the extracted organoids are frozen and stored for future use (for example, drug 162 assays).

We consider the bioreactor design, *e.g.* the hydrogel and media depths, to be fixed (though modelling can provide insights into the role of system geometry on the resulting metabolite concentrations). The glucose concentration in the upstream reservoir is also fixed. The bioreactor operating parameters that can be varied are the media inlet flow rate and the cell seeding density in the hydrogel. The key biological question we seek to answer using mathematical modelling is "how do the bioreactor operating conditions and cell characteristics influence the glucose and lactate concentrations within the CXP1 bioreactor".

## 169 2.1.1 Parameter values

The CXP1 geometry and relevant parameter values (*e.g.* bioreactor length, hydrogel and culture media layer depths, maximum culture media flow velocity, and initial cell seeding density) are outlined in Ellis et al. (2019) and stated in Table 1. The hydrogel used in the CXP1 protocol is Corning Matrigel Matrix and the culture media is a modified form of Dulbecco's modified Eagle medium (DMEM), both of which are described in Ellis et al. (2019).

175 The diffusivities of glucose and lactate in hydrogel and media used in our model are taken from the 176 literature (see Table 1). Our model can be specialised for different cell lines, via characterisation of their rates of proliferation and glucose consumption. In Table 1, we state typical values for rates of cell 177 178 proliferation and glucose consumption, estimated from CXP1 experimental data of several colorectal 179 cancer organoid cell lines. We were also able to obtain averaged values for lactate concentration in the 180 culture media layer at the end of the experiment empirically, which are similar to the values predicted by the model. Estimating model parameter values from experimental data can be challenging, although there 181 have been advances in predicting cellular proliferation rates, e.g. via machine learning methods (Mehrian 182 183 et al., 2020a).

While the current CXP1 operating conditions have been empirically chosen to be specialised for colorectal cancer organoids, a key advantage of mathematical modelling is that it facilitates consideration of metabolite transport within CXP1 for other cell lines (which is the intent of Cellesce). This knowledge will streamline the adaptation of the CXP1 bioreactor to expanding organoids with significantly different behaviour, *e.g.* non-cancerous organoids.

# 189 2.2 Mathematical model

# 190 2.2.1 Governing equations

191 Motivated by the specific bioreactor set-up, parameter values, cell densities, and metabolite concentrations, discussed in Section 2.1, we neglect stochastic effects and adopt a continuum modelling 192 approach. We consider a two-dimensional slice of the bioreactor, and adopt a Cartesian coordinate system 193  $\boldsymbol{x} = (x, z)$  with origin at the bottom–left corner of the domain (see Figure 1). We denote time by t. The 194 hydrogel region of the bioreactor is  $(x, z) \in [0, L] \times [0, h_H]$  (yellow region in Figure 1) and the media 195 region is  $(x, z) \in [0, L] \times [h_H, h_M]$  (blue region in Figure 1). We denote the glucose concentration 196 by c = c(x, z, t) and the lactate concentration by w = w(x, z, t), with subscripts M and H to denote 197 198 concentrations in the media and hydrogel, respectively. We define the model parameters introduced below, together with their typical values, in Table 1. 199

In the hydrogel, the glucose and lactate are transported via diffusion and glucose is consumed by organoids, which subsequently produce lactate. For the organoids (cell aggregates), we model the reaction terms through effective (bulk) sink/source terms over the hydrogel. Such an approach can be mathematically justified through a formal averaging procedure, such as the asymptotic homogenisation carried out for related systems in Dalwadi et al. (2018); Dalwadi and King (2020). The equations governing metabolite transport within the hydrogel,  $(x, z) \in [0, L] \times [0, h_H]$ , are then:

$$\frac{\partial c_H}{\partial t} = D_{CH} \nabla^2 c_H - r(t, \boldsymbol{x}, c_H, w_H) n(t), \qquad (2.1)$$

$$\frac{\partial w_H}{\partial t} = D_{WH} \nabla^2 w_H + s(t, \boldsymbol{x}, c_H, w_H) n(t), \qquad (2.2)$$

where r and s denote the rates of glucose consumption and lactate production per cell, respectively (units mol cell<sup>-1</sup> s<sup>-1</sup>) and n(t) is the cell density at time t (units cell m<sup>-2</sup>). We assume the cells proliferate at rate p, so that the cell density is

$$n(t) = N_0 e^{pt},\tag{2.3}$$

where  $N_0$  is the spatially uniform initial cell-seeding density. While cell growth is likely to have some dependence on the glucose consumption and local lactate concentration, we assume, as a first approximation, that glucose and lactate concentrations are not growth-rate limiting. Thus, due to the spatially uniform initial cell density, the cell density does not vary in space.

During glycolysis, one glucose molecule produces energy and two lactate molecules (Liberti and Locasale,
2016). Motivated by this, we impose

$$s = 2r. (2.4)$$

In general, we expect the glucose consumption to be a monotonically increasing function of glucoseconcentration. For simplicity, we assume that

$$r = \nu_C c_H, \tag{2.5}$$

217 where  $\nu_C$  is a constant (units m<sup>2</sup> cell<sup>-1</sup> s<sup>-1</sup>) representing the rate of glucose consumption per unit cell 218 density.

In the media,  $(x, z) \in [0, L] \times [h_H, h_M]$ , the advection–diffusion equations for metabolite transport are:

$$\frac{\partial c_M}{\partial t} + u(z)\frac{\partial c_M}{\partial x} = D_{CM}\nabla^2 c_M,$$
(2.6)

$$\frac{\partial w_M}{\partial t} + u(z)\frac{\partial w_M}{\partial x} = D_{WM}\nabla^2 w_M,$$
(2.7)

219 where u(z) is the horizontal media flow. Given the slow nature of the flow and geometry of the flow domain, 220 the flow is well-approximated by pressure-driven lubrication flow with a free surface, so that u(z) is the 221 half-Poiseuille flow:

$$u(z) = [u] \frac{(z - h_H)^2}{(h_M - h_H)^2},$$
(2.8)

222 where [u] is the maximum flow velocity.

Governing equations Eqs. (2.1)–(2.8) require appropriate boundary, initial, and interfacial conditions. The boundaries in the hydrogel are solid walls and we impose zero flux of glucose and lactate at x = 0, L:

$$-D_{CH}\frac{\partial c_H}{\partial x} = -D_{WH}\frac{\partial w_H}{\partial x} = 0.$$
(2.9)

We assume the concentrations of glucose and lactate in the inlet pipe are maintained at the constant values  $c_{-\infty}$  and 0, respectively. We assume pointwise continuity of metabolite flux at the inlet, x = 0:

$$u(z)c_M - D_{CM}\frac{\partial c_M}{\partial x} = u(z)c_{-\infty}, \ u(z)w_M - D_{WM}\frac{\partial w_M}{\partial x} = 0;$$
(2.10)

and we impose no diffusive flux of metabolites at the outlet, x = L:

$$-D_{CM}\frac{\partial c_M}{\partial x} = -D_{WM}\frac{\partial w_M}{\partial x} = 0,$$
(2.11)

noting that the metabolites can leave the bioreactor via advection. We impose no-flux conditions for the metabolites at the base of the hydrogel, z = 0, and at the top of the media layer,  $z = h_M$ :

$$-D_{CH}\frac{\partial c_H}{\partial z} = -D_{WH}\frac{\partial w_H}{\partial z} = 0 \text{ at } z = 0 \text{ and } -D_{CM}\frac{\partial c_M}{\partial z} = -D_{WM}\frac{\partial w_M}{\partial z} = 0 \text{ at } z = h_M.$$
(2.12)

223 At the media-hydrogel interface,  $z = h_H$ , we impose continuity of metabolite concentration and flux:

$$c_M = c_H, w_M = w_H, \text{ and } D_{CM} \frac{\partial c_M}{\partial z} = D_{CH} \frac{\partial c_H}{\partial z}, D_{WM} \frac{\partial w_M}{\partial z} = D_{WH} \frac{\partial w_H}{\partial z}.$$
 (2.13)

A schematic of these boundary conditions on the domain geometry is given in Figure 2.

As initial conditions, we assume that the glucose concentration in the media equals the glucose concentration in the upstream reservoir,  $c = c_{-\infty}$ , the glucose concentration in the hydrogel is zero, and that there is no lactate throughout the bioreactor:

$$c_H = 0, \quad c_M = c_{-\infty}, \quad w_H = w_M = 0 \quad \text{at } t = 0.$$
 (2.14)

#### 228 2.2.2 Typical timescales

The typical parameter values, given in Table 1, reveal that the physical processes included in our model 229 act over three different timescales: hours, days, and months, as shown in Table 2. Diffusion in the z-230 direction occurs over the timescale of hours; media flow, glucose consumption, lactate production, and 231 cell proliferation occur over the timescale of a day; and x-diffusion occurs over the timescale of months. 232 This scaling analysis reveals that flow markedly enhances metabolite transport in the x-direction and 233 that, within the media, advection dominates diffusive transport of metabolites in the horizontal direction. 234 The separation of timescales renders the system stiff and, as such, care is needed when implementing 235 numerical methods for its solution. At the same time, it leads naturally to the identification of large and 236 small dimensionless parameters which can be exploited for model reduction (see Section 2.3). 237

## 238 2.2.3 Non-dimensionalisation

We non-dimensionalise the problem to identify the relative importance of each transport mechanism. We introduce the following non-dimensional variables, for  $i \in \{H, M\}$ :

$$X = \frac{x}{L}, \ Z = \frac{z}{\epsilon L}, \ T = \frac{t}{[t]}, \ U(Z) = \frac{u}{[u]}, \ C_i = \frac{c_i}{c_{-\infty}}, \ W_i = \frac{w_i}{c_{-\infty}},$$
(2.15)

where X = (X, Z),  $\epsilon = h_M/L \ll 1$  is the ratio between vertical and horizontal lengthscales, [t] is the timescale, and [u] is the maximum flow velocity. The bioreactor domain is then  $(X, Z) \in [0, 1] \times [0, 1]$  and the media–hydrogel interface is at dimensionless position  $Z = H_H =: h_H/(\epsilon L)$ . Metabolite concentrations are non–dimensionalised with the upstream reservoir glucose concentration,  $c_{-\infty}$ . We fix the timescale of interest to be 1 day, so that we consider the transport on the same timescale as cell growth.

Using the scalings Eq. (2.15), the governing equations Eqs. (2.1)–(2.7) become, for  $X \in [0, 1]$ ,

$$\epsilon^2 \frac{\partial C_H}{\partial T} = d_{CH} \left( \epsilon^2 \frac{\partial^2 C_H}{\partial X^2} + \frac{\partial^2 C_H}{\partial Z^2} \right) - \epsilon^2 \rho C_H e^{PT} \quad \text{for } Z \in [0, H_H], \quad (2.16)$$

$$\epsilon^2 \frac{\partial W_H}{\partial T} = d_{WH} \left( \epsilon^2 \frac{\partial^2 W_H}{\partial X^2} + \frac{\partial^2 W_H}{\partial Z^2} \right) + 2\epsilon^2 \rho C_H e^{PT} \text{ for } Z \in [0, H_H], \quad (2.17)$$

$$\epsilon^2 \frac{\partial C_M}{\partial T} + \epsilon^2 \mu U(Z) \frac{\partial C_M}{\partial X} = d_{CM} \left( \epsilon^2 \frac{\partial^2 C_M}{\partial X^2} + \frac{\partial^2 C_M}{\partial Z^2} \right) \qquad \text{for } Z \in [H_H, 1], \quad (2.18)$$

$$\epsilon^2 \frac{\partial W_M}{\partial T} + \epsilon^2 \mu U(Z) \frac{\partial W_M}{\partial X} = d_{WM} \left( \epsilon^2 \frac{\partial^2 W_M}{\partial X^2} + \frac{\partial^2 W_M}{\partial Z^2} \right) \qquad \text{for } Z \in [H_H, 1], \quad (2.19)$$

246 with

$$U(Z) = \frac{(Z - H_H)^2}{(1 - H_H)^2}.$$
(2.20)

247 The dimensionless parameters in Eqs. (2.16)–(2.20) are:

$$\mu = \frac{[u][t]}{L}, \quad \rho = [t]\nu_C N_0, \quad P = p[t],$$

$$(d_{CH}, d_{CM}, d_{WH}, d_{WM}) = \frac{[t]}{L^2} (D_{CH}, D_{CM}, D_{WH}, D_{WM}).$$
(2.21)

We provide a physical interpretation of these dimensionless parameters and their typical values in Table 3. The boundary and initial conditions, Eqs. (2.9)–(2.14), become:

$$-d_{CH}\frac{\partial C_H}{\partial X} = 0, \qquad -d_{WH}\frac{\partial W_H}{\partial X} = 0 \text{ at } X = 0, 1, \qquad (2.22)$$

$$\mu UC_M - d_{CM} \frac{\partial C_M}{\partial X} = \mu U, \quad \mu UW_M - d_{WM} \frac{\partial W_M}{\partial X} = 0 \text{ at } X = 0, \qquad (2.23)$$

$$-d_{CM}\frac{\partial C_M}{\partial X} = 0, \qquad -d_{WM}\frac{\partial W_M}{\partial X} = 0 \text{ at } X = 1, \qquad (2.24)$$

$$\frac{\partial C_H}{\partial Z} = \frac{\partial W_H}{\partial Z} = 0 \text{ at } Z = 0, \qquad (2.25)$$

$$\frac{\partial C_M}{\partial Z} = \frac{\partial W_M}{\partial Z} = 0 \text{ at } Z = 1, \qquad (2.26)$$

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$$C_M = C_H, \quad W_M = W_H \text{ at } Z = H_H,$$
 (2.27)

$$d_{CH}\frac{\partial C_H}{\partial Z} = d_{CM}\frac{\partial C_H}{\partial Z}, \quad d_{WH}\frac{\partial W_H}{\partial Z} = d_{WM}\frac{\partial W_M}{\partial Z} \text{ at } Z = H_H,$$
 (2.28)

$$C_H = 0, \quad C_M = 1, \quad W_H = W_M = 0 \text{ at } T = 0.$$
 (2.29)

248

## 249 2.2.4 Numerical solution of full model

We solve the full two-dimensional system, Eqs. (2.16)–(2.19) and (2.22)–(2.29), using the parameter values given in Table 3, via a finite-element method, using COMSOL Multiphysics® software. The results are checked to be independent of mesh size (results not shown). We plot the metabolite concentration profiles at dimensionless times T = 1, 3, 7, corresponding to one, three and seven days, in Figure 3. Note that we observe little variation in metabolite concentration in the vertical direction for the parameter values given in Table 3.

## 256 2.3 Model reduction

As discussed in Section 2.2.2, the different transport mechanisms in the system have associated timescales that can be grouped into either hours, days, or months. This is made explicit in the dimensionless system through the presence of the small parameter  $\epsilon$ . We propose a systematic model reduction, with the key advantage of reducing the complexity of the model while retaining the physical processes which dominate over the timescale of interest.

#### 262 2.3.1 Longwave approximation

Motivated by the long, thin geometry of the bioreactor, characterised by  $\epsilon \ll 1$ , and the lack of variation in Z compared to X revealed in Figure 3, we now systematically average Eqs. (2.16)–(2.19) and (2.22)–(2.29) in Z to derive the appropriate reduced lubrication model, referred to as the *longwave approximation*.

266 In the asymptotic analysis that follows, we consider the limit  $\epsilon \to 0$ , and assume all other dimensionless parameters remain  $\mathcal{O}(1)$  as  $\epsilon \to 0$ . This distinguished limit is consistent with the values of dimensionless 267 268 parameters given in Table 3, and assumes that diffusion in the vertical direction is the dominant transport 269 mechanism for the bioreactor geometry. We note that our choice of time scaling, [t] = 1 day, means that we 270 are investigating this system over the timescale of days. We could study the behaviour of this system over shorter timescales, and its transition to the timescale of days, if we systematically considered the timescale 271  $T = \mathcal{O}(\epsilon^2)$ . However, this will not be of fundamental importance to the problem we study here, and we do 272 273 not pursue this further.

274 We consider the following asymptotic expansions for the dependent variables:

$$f \sim f_0 + \epsilon^2 f_1 + \cdots$$
, as  $\epsilon \to 0$ , where  $f \in \{C_M, C_H, W_M, W_H\}$ . (2.30)

We note that the  $\mathcal{O}(\epsilon^2)$  size of the first-correction term is standard in lubrication-type models, and arises due to the size of the terms neglected in the leading-order problem. In the standard manner, we substitute Eq. (2.30) into the governing equations, Eqs. (2.16)–(2.19) and (2.22)–(2.29), and equate coefficients of  $\mathcal{O}(\epsilon^n)$ .

At leading order, the metabolite transport is given by

$$0 = \frac{\partial^2 f_{j0}}{\partial Z^2} \text{ where } f \in \{C, W\} \text{ and } j \in \{H, M\}.$$

$$(2.31)$$

Hence, we see that the leading-order mass transport is driven entirely by vertical diffusion, consistent withour discussion of timescales above.

Integrating Eq. (2.31) subject to the leading-order versions of the appropriate boundary conditions, Eqs. (2.25)–(2.28), we deduce that  $C_{H0}$ ,  $C_{M0}$ ,  $W_{H0}$ ,  $W_{M0}$  are independent of vertical position, Z. This is consistent with the numerical solutions seen in Figure 3. Given the continuity of concentration condition, Eq. (2.27), we deduce that

$$C_{H0}(T,X) = C_{M0}(T,X), \quad W_{H0}(T,X) = W_{M0}(T,X) \quad \text{for all } Z.$$
 (2.32)

285 However, the correct dependence of the metabolite profiles on T and X is currently undetermined.

To calculate this dependence, we proceed to  $\mathcal{O}(\epsilon^2)$  and derive an appropriate solvability condition. At  $\mathcal{O}(\epsilon^2)$ , the governing equations are

$$d_{CH}\frac{\partial^2 C_{H1}}{\partial Z^2} = \frac{\partial C_{H0}}{\partial T} - d_{CH}\frac{\partial^2 C_{H0}}{\partial X^2} + \rho C_{H0}e^{pT} \qquad \text{for } Z \in [0, H_H),$$
(2.33)

$$d_{WH}\frac{\partial^2 W_{H1}}{\partial Z^2} = \frac{\partial W_{H0}}{\partial T} - d_{WH}\frac{\partial^2 W_{H0}}{\partial X^2} - 2\rho C_{H0}e^{pT} \qquad \text{for } Z \in [0, H_H), \tag{2.34}$$

$$d_{CM}\frac{\partial^2 C_{H1}}{\partial Z^2} = \frac{\partial C_{M0}}{\partial T} + \mu U(Z)\frac{\partial C_{M0}}{\partial X} - d_{CM}\frac{\partial^2 C_{M0}}{\partial X^2} \qquad \text{for } Z \in (H_H, 1], \tag{2.35}$$

$$d_{WM}\frac{\partial^2 W_{H1}}{\partial Z^2} = \frac{\partial W_{M0}}{\partial T} + \mu U(Z)\frac{\partial W_{M0}}{\partial X} - d_{WM}\frac{\partial^2 W_{M0}}{\partial X^2} \qquad \text{for } Z \in (H_H, 1].$$
(2.36)

Integrating each equation over the vertical coordinate and applying the no flux conditions, Eqs. (2.25) and (2.26), at  $\mathcal{O}(\epsilon^2)$  yields:

$$\left. d_{CH} \frac{\partial C_{H1}}{\partial Z} \right|_{Z=H_H} = H_H \left( \frac{\partial C_{H0}}{\partial T} - d_{CH} \frac{\partial^2 C_{H0}}{\partial X^2} + \rho C_{H0} e^{PT} \right), \tag{2.37}$$

$$\left. d_{WH} \frac{\partial W_{H1}}{\partial Z} \right|_{Z=H_H} = H_H \left( \frac{\partial W_{H0}}{\partial T} - d_{WH} \frac{\partial^2 W_{H0}}{\partial X^2} - 2\rho C_{H0} e^{PT} \right), \tag{2.38}$$

$$-\left. d_{CM} \frac{\partial C_{M1}}{\partial Z} \right|_{Z=H_H} = \left(1 - H_H\right) \left( \frac{\partial C_{M0}}{\partial T} + \mu \bar{U} \frac{\partial C_{M0}}{\partial X} - d_{CM} \frac{\partial^2 C_{M0}}{\partial X^2} \right), \tag{2.39}$$

$$- d_{WM} \frac{\partial W_{M1}}{\partial Z} \bigg|_{Z=H_H} = (1 - H_H) \left( \frac{\partial W_{M0}}{\partial T} + \mu \bar{U} \frac{\partial W_{M0}}{\partial X} - d_{WM} \frac{\partial^2 W_{M0}}{\partial X^2} \right),$$
(2.40)

286 where the depth-averaged flow velocity,  $\overline{U}$  is given by:

$$\bar{U} = \frac{1}{1 - H_H} \int_{H_H}^1 U(Z) \, \mathrm{d}Z = \frac{1}{3}.$$
(2.41)

Recalling the continuity of flux condition, Eq. (2.28), and that  $C_{H0} = C_{M0}$  and  $W_{H0} = W_{M0}$ , we combine the above expressions for the glucose and lactate concentrations in the media and hydrogel to derive the longwave approximation:

$$\alpha \frac{\partial C_{M0}}{\partial T} + \beta \frac{\partial C_{M0}}{\partial X} = \delta_C \frac{\partial^2 C_{M0}}{\partial X^2} - \gamma C_{M0} e^{PT}, \qquad (2.42)$$

$$\alpha \frac{\partial W_{M0}}{\partial T} + \beta \frac{\partial W_{M0}}{\partial X} = \delta_W \frac{\partial^2 W_{M0}}{\partial X^2} + 2\gamma C_{M0} e^{PT}, \qquad (2.43)$$

287 where we have introduced the following parameters for ease of notation:

$$\theta = \frac{H_H}{1 - H_H}, \quad \alpha = 1 + \theta, \quad \beta = \mu \bar{U}, \quad \gamma = \theta \rho, \quad \delta_C = d_{CM} + \theta d_{CH}, \quad \delta_W = d_{WM} + \theta d_{WH}. \tag{2.44}$$

We derive the appropriate boundary and 'initial' conditions for Eqs. (2.42) and (2.43) in a similar manner, by integrating the leading over terms of Eqs. (2.22)–(2.24) and (2.29) over Z between 0 and 1. We solve Eqs. (2.42) and (2.43) subject to the following boundary and 'initial' conditions:

$$\beta C_{M0} - \delta_C \frac{\partial C_{M0}}{\partial X} = \beta \text{ at } X = 0, \qquad (2.45)$$

$$\beta W_{M0} - \delta_W \frac{\partial W_{M0}}{\partial X} = 0 \text{ at } X = 0, \qquad (2.46)$$

$$\frac{\partial C_{M0}}{\partial X} = \frac{\partial W_{M0}}{\partial X} = 0 \text{ at } X = 1, \qquad (2.47)$$

$$C_{M0} = \frac{1}{\alpha} \text{ and } W_{M0} = 0 \text{ at } T = 0 \text{ for } 0 \le X \le 1.$$
 (2.48)

The reason we refer to Eq. (2.48) as 'initial' conditions is because they actually represent asymptotic matching conditions with the earlier timescale problem we mentioned previously. This is the reason why there is a discontinuity in the boundary and 'initial' conditions as  $X, T \rightarrow 0$ . If it were of interest to understand this limit further, one could investigate this region using the scalings  $X = O(\epsilon), T = O(\epsilon^2)$ . Given that this asymptotic region does not affect any of our subsequent analysis, for brevity we do not pursue it further here.

Eqs. (2.42), (2.43) and (2.45)–(2.48) define the longwave approximation model. We will analyse this reduced system in more detail in Section 3. First, we derive a further reduction of the longwave approximation, by exploiting the separation in scales between horizontal diffusion and the remaining transport mechanisms, namely advection with the media flow, glucose consumption, and lactate production.

298 2.3.2 Sublimit of longwave approximation

From the typical parameter values given in Table 3, we note that the timescale of horizontal diffusion is significantly longer than the remaining transport mechanisms. Given that the longwave approximation derived in Section 2.3.1 is a distinguished asymptotic limit, we can include the separation of scales involved in horizontal diffusion by directly considering the sub-limit  $d_{CH}$ ,  $d_{CM}$ ,  $d_{WH}$ ,  $d_{WM} \rightarrow 0$ , corresponding to  $\delta_C$ ,  $\delta_W \rightarrow 0$  in Eqs. (2.42), (2.43) and (2.45)–(2.48). We refer to this as the *sublimit approximation*. This procedure results in the following governing equations for advection–dominated transport:

$$\alpha \frac{\partial C_{M0}}{\partial T} + \beta \frac{\partial C_{M0}}{\partial X} = -\gamma C_{M0} \exp(PT), \qquad (2.49)$$

$$\alpha \frac{\partial W_{M0}}{\partial T} + \beta \frac{\partial W_{M0}}{\partial X} = 2\gamma C_{M0} \exp(PT), \qquad (2.50)$$

with boundary and initial conditions

$$C_{M0} = 1, \quad W_{M0} = 0 \quad \text{at } X = 0,$$
 (2.51)

$$C_{M0} = \frac{1}{\alpha}, \quad W_{M0} = 0 \quad \text{at } T = 0.$$
 (2.52)

We note that the limit we have taken is singular in that the small parameters (diffusivities) pre-multiply the second-order spatial derivatives. As such, we have lost the ability to prescribe the outlet boundary conditions at X = 1, though we note that this boundary condition could be imposed through the analysis of an appropriate (weak) boundary layer near X = 1.

A benefit of this sublimit reduction is that we are able to construct analytic solutions for the glucose concentration, using the method of characteristics. The solution is split into two distinct regions: Region 1, given by  $0 < \beta T < \alpha X$ ; and Region 2, given by  $0 < \alpha X < \beta T$ :

$$C_{M0} = \begin{cases} \frac{1}{\alpha} \exp\left(\frac{\gamma}{\alpha P} \left(1 - e^{PT}\right)\right) & \text{for } 0 < \beta T < \alpha X, \qquad (2.53) \end{cases}$$

$$\left( \exp\left(\frac{\gamma}{\alpha P} \left( e^{-P\left(\frac{\alpha}{\beta}X - T\right)} - e^{PT} \right) \right) \qquad \text{for } 0 < \alpha X < \beta T, \qquad (2.54)$$

The solution (2.58)–(2.59) is discontinuous across the boundary separating the two regions,  $X = \beta T/\alpha$ , which we refer to as the *dividing characteristic*. The reason for this is that Region 1 is forced by the initial conditions whereas Region 2 is forced by the boundary conditions, and there is a discontinuity in these conditions near T = 0, X = 0 (which could be smoothed through an appropriate asymptotic analysis of the earlier timescale, as mentioned previously). As no information from the boundary condition propagates into Region 1, cells in Region 1 do not feel the effect of any replenishment by the flow. As such, we refer to Region 1 as the *unreplenished region* and Region 2 as the *replenished region*.

Using the method of characteristics, we can write the lactate concentration as a single integral of known functions:

$$W_{M0}(S,\tau) = \int_0^\tau 2\gamma C_{M0}(T(S,\tau), X(S,\tau)) e^{PT(S,\tau)} \,\mathrm{d}\tau \quad \text{with } W_{M0} = 0 \text{ at } \tau = 0,$$
(2.55)

where we define the characteristic variables  $(S, \tau)$  as

$$S = \alpha X - \beta T \quad \text{and} \quad \tau = \begin{cases} \frac{T}{\alpha} & \text{for } \beta T < \alpha X, \\ Y \end{cases}$$
(2.56)

$$\left\{ \frac{X}{\beta} \quad \text{for } \alpha X < \beta T.$$
(2.57)

As outlined in the Supplementary Material, we can evaluate the integral in Eq. (2.55) to obtain the solution

$$W_{M0} = \begin{cases} \frac{2}{\alpha} \left( 1 - \exp\left(\frac{\gamma}{P\alpha} \left(1 - e^{PT}\right)\right) \right) & \text{for } 0 < \beta T < \alpha X, \end{cases}$$
(2.58)

$$W_{M0} = \left\{ 2 \left( 1 - \exp\left(\frac{\gamma}{P\alpha} \left( e^{-P\left(\frac{\alpha}{\beta}X - T\right)} - e^{PT} \right) \right) \right) \quad \text{for } 0 < \alpha X < \beta T.$$
 (2.59)

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We note that the quantity  $2C_{M0} + W_{M0}$  is conserved along the characteristics defined by  $dX/dT = \alpha/\beta$  (*i.e.* in the advective frame of reference). This means that the following relationships are satisfied between glucose and lactate concentrations:

$$2C_{M0} + W_{M0} = \frac{2}{\alpha} \quad \text{for } 0 < \beta T < \alpha X,$$
 (2.60)

$$2C_{M0} + W_{M0} = 2$$
 for  $0 < \alpha X < \beta T$ , (2.61)

312 where the differing constants are due to the 'initial' information on the characteristics arising from the 313 actual initial conditions for  $0 < \beta T < \alpha X$  (Region 1) and the replenishment boundary conditions for 314  $0 < \alpha X < \beta T$  (Region 2).

## 3 RESULTS

## 315 3.1 Model behaviour and comparison

We now discuss and compare results obtained from our reduced models and the full system. This will allow us to understand when each reduced model is a useful systematic reduction.

The longwave approximation model, Eqs. (2.42), (2.43) and (2.45)–(2.48), is solved numerically using the Chebfun toolbox in MATLAB. For the sublimit approaximation model, Eqs. (2.49)–(2.52), we obtain an analytical expression for the glucose concentration, and the lactate concentration is numerically computed from Eq. (2.50) subject to Eq. (2.51) with a Runge-Kutta method using the in-built ODE solver ode45 in MATLAB. For each numerical approach, we perform convergence tests to ensure the results are independent of mesh size (results not shown).

Computationally, there is a significant difference between the models: on a standard desktop, the full problem is solved in O(180s); the longwave approximation in O(20s); and the sublimit approximation in O(4s). That is, there is a nearly ten-fold speed-up in solving the longwave approximation compared to the full model, and the sublimit is five times quicker to solve than the longwave approximation. As we see later, rapid computation of solutions will allow us to perform parameter sensitivity analyses efficiently.

329 To present the model solutions over space and time, we average solutions of the full 2D model over 330 Z, to facilitate comparison with solutions of the reduced models (Figure 4). We see that the glucose concentration behaviour appears to be split into two approximate regions, divided by a straight line in 331 (X, T)-space that goes through the origin and reaches the end of the X-domain (X = 1) at  $T \approx 4$  (Figure 332 4A). In the lower-right region, the glucose concentration appears to be approximately constant in space, 333 334 and to decrease over time. However, in the upper-left region, there is a clear spatial dependence in the 335 glucose concentration, which appears to decrease in X until it reaches the lower-right region. The lactate 336 concentration behaviour appears to be split into the same two approximate regions (Figure 4B), though 337 the demarcation is less defined than for glucose. In the lower-right region, the lactate concentration also appears to be approximately constant in space, but now increases over time. In the upper-left region, the 338 lactate concentration appears to approximately increase in X until it reaches the lower-right region. To 339 compare these results with the reduced models, we also present solutions for the longwave approximation 340 (Figures 4C, 4D) and sublimit approximation (Figures 4E, 4F). We see that the longwave approximation 341 is an excellent approximation of the full system through the entire domain. The sublimit is also a good 342 approximation of the full model except in a small neighbourhood of the dividing characteristic,  $\alpha X = \beta T$ . 343 The sublimit solution is discontinuous across the dividing characteristic because it neglects horizontal 344 diffusion. Appropriate smoothing could be included in the sublimit by investigating a thin boundary layer 345

in the neighbourhood of this discontinuity in which diffusive effects are once again important. We also
note that the dividing characteristic is in approximately the same place as the boundary between regions
noted in the full model in Figures 4A and 4B. We investigate and interpret this observation below.

At this stage, we conclude that when information close to the dividing characteristic is of interest, the longwave approximation should be used instead of the sublimit approximation. If this information is not important, the sublimit approximation should be used since it is faster to solve than the full model and the longwave approximation, and it admits analytic solutions for glucose concentration.

We emphasize that our analytic solutions in the sublimit approximation allow us to understand 353 observations from the full numerical solutions. That is, we can use our analytic solutions from the 354 sublimit model to physically interpret our results and provide insight into the underlying physical system. 355 For example, the dividing characteristic ( $\alpha X = \beta T$ ) in the sublimit model represents the division between 356 information propagated from the initial and the boundary conditions. Physically, this means that the effect 357 of fresh media is only experienced at position X at time  $T = \alpha X/\beta$ . At earlier times, glucose delivery to 358 organoids at position X is due to the glucose initially present in the system. This allows us to determine 359 the *metabolite transit time*. That is, the average time taken for metabolite within the *fresh* media to traverse 360 the entire bioreactor 361

$$T^* = \frac{\alpha}{\beta} = \frac{1 + \frac{H_H}{1 - H_H}}{\mu \bar{U}} \approx 4.7 \text{ days.}$$
(3.1)

The above estimate is in good agreement with our observations of the full solution - that different model 362 solutions arise in the two regions on either side of the straight line through the origin that reaches X = 1 at 363  $T \approx 4$ . Hence, we now interpret this observation physically; the regions are separate according to whether 364 or not they have experienced fresh media. Since the media does not traverse the bioreactor with a constant 365 velocity, the metabolite transit time is not the same as the timescale associated with the maximum flow 366 velocity of the system, [t] = 25 hours. The relevant timescale is, therefore, not the one associated with the 367 experimentally imposed flow rate, but rather the metabolite transit timescale, which is associated with the 368 averaged velocity distribution of metabolite across the bioreactor. 369

Additionally, the analytic solution of our sublimit approximation provides insight into why the glucose and lactate concentration appear to be spatially-independent in the lower-right regions (Figure 4). In Region 1 (where  $0 < \beta T < \alpha X$ ), the analytical solutions for metabolite concentrations from the sublimit model are independent of the spatial coordinate. Region 1 is the non-replenished region, *i.e.* it is not replenished from the inlet and subsists on its initial conditions. Given spatially-uniform initial conditions, spatial effects are not seen in the concentration profiles until the wave of replenishment is experienced; this marks the onset of Region 2.

To quantitatively compare the model predictions, we consider the following time-dependent variables: *minimum glucose concentration*,  $C_{\min}(T) = \min_{X}(C(X,T))$ ; *maximum lactate concentration*,  $W_{\max}(T) = \max_{X}(W(X,T))$ ; *spatial position of maximum lactate concentration*,  $X_{\max}(T)$ , where  $W(X_{\max},T) = W_{\max}(T)$ ; and the *lactate concentration at outlet*, W(X = 1, T). We emphasize that Eq. (2.32) allows us to denote the metabolite concentrations  $C_{M0} = C_{H0} = C$  and  $W_{M0} = W_{H0} = W$  for ease of notation.

In Figure 5A, we plot the minimum glucose concentration,  $C_{\min}(T)$ , against time for our two reduced onedimensional models and the Z-averaged full model and compare these values to the predicted minimum glucose concentration in hydrogel, which is found using the full two-dimensional model. We see that the predicted minimum glucose from each model reduction generally agrees well with the minimum glucose within the hydrogel from the full model. The only exceptions to this are around 4–5 days, where the sublimit model disagrees slightly with the other models, and for early times (< 1 day). The first of these is due to the dividing characteristic being important for this metric around 4.7 days, as discussed above. The second is due to our choice of timescale in deriving the reduced model. That is, our reduced models focus over the timescale of days and neglect the initial transient behaviour in the system, as mentioned previously.

Similar plots showing how the maximum lactate concentration,  $W_{\max}(T)$ , changes over time are presented 392 in Figure 5B. Again, the Z-averaged full model and the longwave approximation are in good agreement 393 with the predicted value within the hydrogel. Given that there is initially no lactate in the system, this metric 394 395 avoids the issue with the early time transient behaviour that occurs for the minimum glucose concentration metric. The sublimit approximation systematically overestimates the lactate concentration, though we note 396 397 that this is preferable to underestimation, given the detrimental effects of high lactate concentrations. The 398 overestimation arises because the sublimit approximation neglects the removal effect of lactate transport through horizontal diffusion over the dividing characteristic. 399

We compare the position at which the maximum lactate concentration occurs,  $X_{max}(T)$ , in Figure 5C. We see that  $X_{max}$  is increasing in time, which is consistent with advection being the dominant transport mechanism over the timescale of days (Table 2), as the lactate produced is advected towards the outlet by the media. As seen in Figures 5A and 5B, the sublimit approximation agrees less well with the full model than the longwave approximation, which has excellent agreement.

405 It is infeasible to obtain experimental data for maximum lactate concentrations, which we would need to validate our model. Therefore, we consider the lactate concentration at the media outlet, W(X = 1, T), 406 which is measurable empirically, in Figure 5D. We compare the reduced models to the Z-averaged full 407 solution, the average concentration within the hydrogel at the outlet, and the maximum value in the 408 hydrogel (which are all obtained from numerical solutions to the full 2D system). We find that the lactate 409 concentration at the media outlet is very similar to the maximum lactate concentration within the hydrogel 410 and can, therefore, be used as a proxy for it. The sublimit is a good prediction of the outlet and maximum 411 lactate concentrations at 4 days and earlier, but overestimates the maximum concentration within the 412 hydrogel at 5 days and later. This is again due to the dividing characteristic, and its exit from the domain at 413 4.7 days. 414

# 415 3.2 Bioreactor characterisation

In this section, we start by exploiting our reduced modelling approach to *characterise* the conditions within the bioreactor. We show how the metabolite concentrations depend on the bioreactor operating parameters such as the inlet flow rate and cell seeding density, and the characteristics of the cells, such as the rates of cell proliferation and glucose consumption. Armed with this insight, we then show how the operating parameters can be selected to ensure the biochemical environment within the bioreactor promotes cell growth.

We investigate and quantify the metabolite behaviour by introducing the following time-dependent metrics. We previously defined the *maximum lactate concentration*,  $W_{max}(T)$ , as

$$W_{\max}(T) = \max_{X} (W(X,T)).$$
 (3.2)

We now introduce the cumulative glucose conversion, Q(T), as

$$\boldsymbol{Q}(T) = \frac{\text{glucose consumed}}{\text{glucose supplied}} = \frac{\int_0^T \int_0^1 \gamma C \exp(PT) \, \mathrm{d}X \mathrm{d}T}{\int_0^T (1 - H_H) \mu \bar{U} \, \mathrm{d}T}.$$
(3.3)

In general, it is desirable to choose operating parameters that ensure high glucose conversion, so the maximum amount of glucose supplied to the bioreactor is utilised by the cells, and resource wastage is minimised. However, high glucose conversion will also cause high lactate levels, and lactate concentrations above a critical tolerance,  $W_{tol}$ , can adversely affect organoid growth. To assess this, we define a point X to be *uninhabitable* if  $W(T, X) > W_{tol}$ . We use the metric *proportion of domain which is uninhabitable*,  $P_U(T)$ , defined as

$$\mathbf{P}_{U}(T) = \int_{0}^{1} H\left[W(T, X) - W_{\text{tol}}\right] \, \mathrm{d}X, \tag{3.4}$$

428 where *H* is the Heaviside function. In general, it is desirable to choose operating parameters such that  $P_U$ 429 is minimised for the duration of the bioreactor run. In addition to the time-dependent metrics, it is also 430 helpful to quantify the *time at which intolerable lactate levels are first experienced*, which we refer to as 431 the *turn–off time*, and define as

$$T_{\text{off}} = \min(T) \text{ for } T \in \{T : W(X, T) \ge W_{\text{tol}}\}.$$
 (3.5)

432 In general, it is desirable to choose operating parameters such that  $T_{\text{off}}$  is larger than the duration of the 433 bioreactor run.

There is a trade–off between high glucose conversion and minimising the fraction of the domain which is *uninhabitable*. We show how the mathematical model can be used to identify parameter regimes which strike a balance between promoting glucose conversion and facilitating waste removal in Section 3.2.2.

In addition to the metrics we have introduced to assess metabolite distribution, an important cell–specific metric is the *glucose consumption rate per cell*. In our model, the glucose consumption rate per cell is proportional to the glucose concentration and, thus, we can use results such as Figure 4C to understand the spatial variation in glucose consumption rate per cell. We see that cells nearer the inlet have higher rates of glucose uptake than those closer to the media outlet, and this spatial heterogeneity could lead to spatial variation in cell growth within the physical system.

#### 443 3.2.1 Characterising model behaviour for different organoid lines

Organoid lines differ in many ways including, but not limited to, proliferation rate, glucose consumption 444 rate, the maximum lactate concentration cells can tolerate without affecting cell properties, and minimum 445 glucose level needed for cellular proliferation. To understand the metabolic environment experienced by 446 different organoid lines within the bioreactor, we perform a discrete parameter sensitivity analysis in which 447 we vary the rates of proliferation, P, and glucose consumption per cell,  $\rho$ , for the bioreactor operating 448 regime specified in Table 3. We consider organoid lines whose proliferation rates take the values P = 1/6449 and P = 1, which we refer to as low and high proliferation, respectively, and whose glucose consumption 450 rates take values  $\rho = 0.027$  and  $\rho = 2.7$ , referred to as low and high consumption, respectively. We 451 consider five different organoid lines: (i) with P = 1/6 and  $\rho = 0.027$ ; (ii) with P = 1/6 and  $\rho = 2.7$ ; 452 (*iii*) with P = 1 and  $\rho = 0.027$ ; (*iv*) with P = 1 and  $\rho = 2.7$ ; and the *typical organoid line* considered in 453 Figure 4 (v) with P = 1/3 and  $\rho = 0.27$ . In Figure 6, we plot the metabolite concentration profiles C and 454

455 W for these four organoid lines, (i-iv), expanded under an operating regime which does not otherwise 456 differ. The same results for organoid line (v) are shown in Figures 4C and 4D.

In Figures 6A and 6E, we show organoid line (*i*), cells with low proliferation and low glucose uptake rates. The lactate levels are very low throughout the bioreactor domain and the domain remains within tolerable lactate concentrations for the entire experiment. The glucose concentration in the replenished region is high and remains close to its inlet value, C = 1, so the media flow supplies significantly more glucose into the system than is consumed by the cells. The glucose concentration becomes increasingly homogeneous as time evolves, and consequently the rate of glucose consumption per cell becomes more spatially homogeneous across the bioreactor as time evolves.

We consider organoid line (*ii*), with low proliferation and high glucose uptake rates, in Figures 6B and 6F. We see that this larger uptake rate means that the lactate concentration quickly increases and the majority of the region becomes intolerable, even for slowly proliferating cells. While cells close to the inlet still have reasonably high glucose and low lactate levels, resulting in the rate of glucose uptake per cell being high at the inlet, this quickly decreases as one moves into the bioreactor.

For rapidly proliferating cells with a low rate of glucose uptake (organoid line (*iii*)) Figures 6C and 469 470 6G, we see the environment is tolerable until around day 4 of the experiment. At this point, there are approximately 55 times more cells within the hydrogel than at the start of the experiment. This suggests 471 472 that the selected operating conditions provide tolerable conditions and allow reasonable rate of glucose 473 consumption per cell up to a critical number of cells, but beyond this critical number, the low glucose concentration means the cells have a very low rate of glucose consumption. The lactate concentration 474 475 is reasonably spatially homogeneous, which suggests that all cells will be subject to a similar metabolic 476 environment and therefore be affected by lactate to a similar degree.

477 Finally, we consider cells with high proliferation and high uptake (organoid line (*iv*)), in Figures 6D and 6H. The glucose concentration within the bioreactor decays very quickly over the course of a day, 478 479 and it is never replenished sufficiently by the media flow. As such, the glucose consumption per cell is 480 consistently small away from the inlet region. In the same vein, the lactate concentration quickly increases to above the tolerable level over the course of a day. In contrast to the low proliferation organoid line (*ii*)) 481 (Figures 6B and 6F), the maximum lactate concentration for organoid line (iv) occurs close to the inlet 482 483 rather than in the middle of the bioreactor. This is because the rapid expansion of cells means that lactate is produced very quickly throughout the bioreactor, and so is maximised in the location where glucose is 484 mainly consumed. This indicates that the media flow is too slow to facilitate significant waste removal for 485 486 this organoid line. We note that our cell growth model is not dependent on metabolite concentration, so the cell proliferation rate is unaffected when the metabolic environment is harsh. This limitation is most 487 prominent for the high proliferation and high uptake organoid line, where the cells continue to proliferate 488 489 exponentially in the presence of no glucose and high lactate levels.

Using the metrics we introduced above, we now quantify the behaviour of the bioreactor environment during cell culture for each of the five organoid lines. In Figure 7, we plot the total glucose conversion, Q(T) (Eq. (3.3)), maximum lactate concentration,  $W_{max}(T)$  (Eq. (3.2)), and proportion of uninhabitable domain,  $P_U(T)$  (Eq. (3.4)) (strongly related to the turn-off time), for each of the five organoid lines.

The glucose conversion generically increases over time, as the cells grow. However, the shape of this increase over time varies significantly between the different organoid lines. While solely considering the standard case (organoid line (*v*), given by parameters in Table 3) would suggest that the glucose conversion is approximately linear in time, the additional organoid lines show that this behaviour is not universal.

Cells with high rates of glucose consumption (organoid lines (*ii*) and (*iv*)) have a sharp increase in glucose 498 conversion over the first two days before plateauing. For low rates of glucose consumption, the shape of the 499 glucose conversion curve strongly depends on the cell proliferation rate. For low proliferation (organoid 500 line i), the conversion is low throughout and appears linear. However, for high proliferation (organoid 501 line (*iii*)), the curve has an S-shape. That is, the conversion starts off low, then rapidly increases before 502 plateauing. This rapid increase is linked to the increase in the number of cells in the bioreactor for organoid 503 line (iii), and so we would expect organoid line (i) to exhibit a similar S-shape if the experiment went on 504 for longer. 505

We show the maximum lactate concentration in Figure 7B, where the red line represents  $W = W_{tol}$ , to 506 understand which of these organoid lines are growing in tolerable environments. This graph is qualitatively 507 very similar to that of the glucose conversion, Figure 7A. For the value of  $W_{tol}$  we use, we see that the 508 maximum lactate concentration reaches the tolerated level within 1 day for high uptake cells (organoid lines 509 510 (ii) and (iv)). In comparison, the standard case (organoid line (v)) reaches the maximum tolerated level approximately halfway through the experiment. For the low uptake organoid lines, the proliferation rate 511 again makes a significant difference. For high proliferation (organoid line (iii)), the maximum tolerated 512 level is again reached approximately halfway through the experiment, whereas for low proliferation (cell 513 514 line (i)) the lactate never reaches harmful levels.

We examine the time at which the lactate concentration equals the tolerated lactate concentration in 515 Figure 7C, a graph showing the time-dependent proportion of the domain which is uninhabitable,  $P_U(T)$ , 516 for each organoid line. Notably, we see that as soon as some of the domain becomes uninhabitable, the rest 517 of the domain follows over a short timescale. This can be explained through the insight gained from our 518 sublimit approximation. That is, as Region 1 ( $\alpha X > \beta T > 0$ ) has yet to experience replenishment from 519 the inlet, the lactate concentration in this region is approximately spatially homogeneous, and an increase 520 above the tolerable level will quickly be experienced in a large part of the domain. The turn-off time  $T_{off}$ 521 (Eq. (3.5)) can also be determined from Figure 7C – it is the first time at which  $P_U(T)$  is non-zero. We see 522 that the high glucose consumption organoid lines ((ii) and (iv)) have much smaller turn-off times than 523 the other organoid lines. The lactate concentration for organoid line (i) does not reach  $W_{tol}$  during the 524 experiment, so the turn-off time is larger than the run time of the experiment. 525

There is a trade-off between promoting: (1) high glucose conversion, to ensure resources are not wasted; (2) high glucose consumption rate per cell, to ensure cells absorb sufficient glucose to proliferate; and (3) increasing the turn-off time, to ensure the lactate concentrations within the bioreactor remain tolerable everywhere throughout the experiment. Our model framework allows for efficient quantification of all these metrics. By determining how these metrics vary with bioreactor operating parameters, we can then identify operating conditions that enhance cell growth. We illustrate this in the next section.

532 3.2.2 Determining operating conditions for a given organoid line

In this subsection, we focus on the standard organoid line (v), with proliferation rate and glucose consumption rate given in Table 3. This is the organoid line with a "medium" rate of glucose consumption per cell, and a doubling time of three days. The current operating conditions lead to lactate concentrations above the tolerated level for half of the experimental run time, suggesting that these operating conditions are sub–optimal.

We now determine how the metrics depend on the inlet flow rate for this organoid line, and show how this leads to the identification of flow rates that enhance cell growth. We focus on flow rate as this is an experimental parameter that is easily varied. We investigate flow rates over two order of magnitudes,  $[u] \in [1 \times 10^{-7}, 1 \times 10^{-5}] \text{ m s}^{-1}$ , all within the range of the peristaltic pump used in the CXP1 protocol. 542 In Figure 8, we show how the metrics vary with inlet flow rate. To illustrate the dependence of the metrics 543 on flow rate, we first present time-dependent results for five different flow rates. The glucose conversion monotonically increases in time (Figure 8A), due to the increasing number of cells causing an increased 544 glucose consumption. The effect of increasing flow rate is to decrease the glucose conversion. This is 545 546 because stronger flows correspond to feeding more glucose into the system over a given time period as well as the media spending less time within the bioreactor, so there is less time for the glucose to be consumed 547 548 by the cells. However, we also note that the conversion is relatively insensitive to flow rate: increasing the flow by two orders of magnitude only decreases the conversion by a factor of around six. 549

550 While the time-dependent maximum lactate concentration within the domain monotonically increases for a given flow rate, the effect of varying the flow rate is non-monotonic (Figure 8B). For a given 551 run time of the experiment, there is a flow rate that maximises the maximal lactate concentration. We 552 emphasize that this flow rate will depend on the experimental run time. The reason for there being a flow 553 rate which maximises the maximal lactate concentration (the 'worst' flow rate, in some sense) is due to 554 two competing factors. Firstly, the rate of glucose consumption per cell, and therefore the rate of lactate 555 production, increases with increasing flow rate. Secondly, for slower flow rates the media is not able to 556 advect sufficient quantities of lactate out of the bioreactor to maintain a tolerable lactate level. These two 557 factors combine to produce a worst possible flow rate for a given experimental run time. We also note that 558 up until approximately one day (T = 1), the maximum lactate concentration is the same for all the flow 559 rates considered. This reflects the fact that there is a lag in the production of lactate, and that the lactate 560 production is initially set by the initial conditions rather than the operating regime of the bioreactor. 561

562 In Figure 8C, we plot the proportion of the domain which is uninhabitable against time, for the five 563 different flow rates considered. In general, a lower flow rate corresponds to a sharper increase in the 564 uninhabitable proportion once initially triggered. This is because more of the domain is in the nonreplenished Region 1 for lower flow rates, and the metabolite concentrations are approximately spatially 565 566 independent in Region 1, for reasons discussed above. In addition, we note that a large enough flow rate 567 can ensure that none of the domain becomes uninhabitable for the duration of the experimental run, as we see for a flow rate of  $1 \times 10^{-5} \text{m s}^{-1}$ . However, we also note that increasing the flow rate can have 568 an unwanted effect on the turn-off time. From Figure 8C, we see that increasing the flow rate slightly 569 570 decreases the turn-off time, up to a point. As noted above, for large enough flow rates the system never exhibits intolerable lactate concentrations. 571

We now consider a more finely refined investigation of the effect of flow rate of the system metrics. In Figure 9, we consider the effect of flow rate both on the glucose conversion at day 7 (Figure 9A) and on the turn–off time (Figure 9B).

We see that the relationship between glucose conversion at 7 days and media flow velocity is 575 monotonically decreasing, and the rate of decrease is larger for flows faster than  $[u] = 10^{-6} \mathrm{m \, s^{-1}}$ 576 (Figure 9A). However, as noted above, the turn-off time is not monotonic in the flow rate (see also 577 Figure 8C). We see that there is a minimal turn–off time when the flow is approximately  $2 \times 10^{-6} \text{m s}^{-1}$ . 578 This is the *worst* possible flow rate from the point of view of ensuring the domain remains tolerable for 579 as long as possible. For flow rates below this, the bioreactor is *transport-limited*, either by insufficient 580 glucose delivery to cells or by insufficient waste removal from the bioreactor. For flow rates above this, 581 the turn-off time is *proliferation-limited*, where the rate at which the cell population is growing sets the 582 timescale at which lactate is produced. 583

An advantage of our mathematical modelling framework is that we have been able to easily explore a wide range of parameter values, in this case the flow rate, and explore the nonlinear effects of varying experimental parameters. For example, an experimentalist may start with a slow flow rate of  $10^{-7}$ m s<sup>-1</sup> and conduct a set of experiments over which they increased the flow. Over an order of magnitude increase in flow, they would see no improvement in turn–off time, and therefore might be discouraged from increasing the flow any further. In such a scenario, they would miss finding the flow rate values required for turn–off times greater than 4 days.

The "optimal" operating conditions for the bioreactor will determine glucose and lactate concentrations which (1) yield a specified value for glucose conversion; (2) maintain a glucose consumption rate per cell which is sufficient for cellular proliferation; and (3) predict a turn–off time which is greater than the run time of the experiment. The specific values and relative importance of each of these requirements will depend on the user. Our model reduction facilitates rapid calculation of each metric. Hence, our work could be combined with an optimisation algorithm, with user–specified cost functions, to produce an efficient framework that can identify the bioreactor operating conditions that optimise for growth of organoids.

## 4 **DISCUSSION**

We have presented an unsteady, two-dimensional model of metabolite transport that predicts metabolite 598 concentrations within the CXP1 bioreactor system. We used an asymptotic analysis to systematically 599 derive two reduced models which exploit the extreme spatial and temporal parameter ratios in the system. 600 601 Our model predicts the spatiotemporal distribution of the metabolic environment within the bioreactor, information which is challenging to obtain experimentally. Both reduced models are one-dimensional 602 in space; the *longwave approximation* comprises two coupled reaction-advection-diffusion equations, 603 whereas the sublimit approximation comprises two coupled reaction-advection equations. Our systematic 604 analysis allows us to relate parameters in the reduced models to geometric and operating parameters of the 605 CXP1 system, such as the ratio between the depth of the hydrogel and media layers, and the fluid flux over 606 the hydrogel. We have shown that both reduced models provide good approximations of the full model for 607 most physically relevant parameter regimes. The longwave approximation is an excellent representation 608 throughout the entire domain, whereas the sublimit approximation is a good representation everywhere 609 apart from one specific line in space-time that we are able to calculate. 610

Although the above may appear to suggest that the sublimit approximation is not useful, it does have additional benefits over the longwave approximation. A notable benefit is that it admits analytic solutions in the entire domain. Interpreting these analytic results, and understanding why they are discontinuous across the specific line in space–time, provides insight into the underlying physical system. We find that the specific line in space-time is a dividing characteristic in the (hyperbolic) sublimit approximation we derive. We are able to infer that this line divides the domain into two regions, depending on whether or not the effect of replenishment from the inlet has been experienced.

The flow of media through the bioreactor has the dual function of delivering nutrients to, and removing 618 waste from, the growing organoids. As such, the inlet flow rate needs to be chosen carefully. The systematic 619 reduction we have performed yields models that are easier to solve numerically than the full model. More 620 importantly, they provide insight into the behaviour of the full model, particularly the dominant transport 621 mechanisms. This systematic reduction has enabled us to efficiently characterise the experimental parameter 622 space for given cell characteristics. One key outcome from this analysis is our prediction of a 'worst-case' 623 flow rate that minimises the turn-off time (the time when intolerable lactate concentrations first occur), 624 Eq. (3.5). Our model reduction has allowed us to understand why this minimum arises: for higher flow 625

rates, the lactate is washed away more quickly (the bioreactor is in a proliferation–limited regime), for
lower flow rates the lactate is produced more slowly since glucose is not delivered quickly enough (the
bioreactor is in a transport–limited regime).

To understand how outcomes change as the control parameters are varied, we introduced the following time-dependent metrics which characterise bioreactor performance:

- *Glucose conversion* is the ratio between the total amounts of consumed and supplied glucose. It is desirable to minimise the amount of resources, *e.g.* glucose, required for bioreactor operation, which corresponds to maximising glucose consumption.
- *Maximum lactate concentration* within the bioreactor represents the worst metabolic environment experienced by the cells. High lactate concentrations have a detrimental effect on cells (Romero-Garcia et al., 2016), and therefore an ideal bioreactor operating regime would have low maximum lactate concentrations.
- *Proportion of uninhabitable domain* is the fraction of the domain where the lactate concentrations
   exceeds the maximum tolerated level for the specific organoid line. An operating regime is improved if
   the proportion of the domain which is uninhabitable decreases, and an 'ideal' operating regime would
   maintain lactate levels below the maximum tolerable level for the entire experiment.
- *Turn-off time* is the time at which lactate concentration first reaches levels which are intolerable for
   the cells. To optimise operating conditions, the turn-off time should be increased. Ideally, the turn-off
   time should exceed the run time of the experiment.

645 Different bioreactor operating conditions will yield different values of these metrics. The relative importance 646 of each metric will depend on the particular organoid line being investigated and the specific user 647 requirements. Our work provides a framework for efficiently determining desirable bioreactor operating 648 conditions for given cell properties.

In this study, we performed a systematic model reduction to study metabolite transport within the CXP1 649 bioreactor, whose geometry differs significantly from other bioreactors, such as hollow fibre or perfusion 650 bioreactors. An important insight gained from our model reduction is the identification of the transport 651 mechanisms that are dominant on our timescale of interest. We performed model reductions in two ways: 652 (1) we exploited the slender geometry of the system, to obtain the *longwave approximation*; and then (2) we 653 exploited the separation of timescales of the physical processes in play, to derive the sublimit approximation. 654 655 By systematically reducing our original model (Eqs. (2.16)–(2.19) and (2.22)–(2.29)), we have simplified a two-dimensional parabolic PDE system first to a one-dimensional parabolic PDE system (the longwave 656 approximation), and then to a one-dimensional hyperbolic PDE system (the sublimit approximation). A 657 658 significant advantage of this approach is the analytical tractability of the sublimit approximation. As a result, we can construct explicit expressions for the metabolite concentrations across the entire bioreactor 659 that reveal both the spatiotemporal-dependence and the dependence on the control parameters, e.g. flow 660 rate, of the metabolite concentrations in the bioreactor. We have shown that the reduced models serve 661 as excellent approximations of the full system and are much easier to solve numerically. We have also 662 663 identified the small region of space-time where the assumptions required for the validity of sublimit model break down. 664

665 There are a number of interesting possible extensions to this work. For example, the optimal operating 666 conditions are likely to change during the course of organoid growth. Future modelling work could 667 predict how, and when, operating conditions should change to account for this growth. While we have 668 considered steady flows, it would be straightforward to extend our framework to examine more complex 669 flow behaviours, such as oscillating flows, or three-dimensional effects. The potential use of unsteady 670 flows will be of particular interest when minimisation of spatial variation in metabolite concentrations 671 across the bioreactor is important, as we have seen that steady flows with little spatial variation in 672 metabolite concentration also have very low conversion (see Figures 6A, 6E, 7A). The ability to change the 673 mathematical flow model when predicting the metabolite concentrations is particularly useful because it 674 can be done in advance of engineering the prototype bioreactors needed to test the system experimentally.

In this work, we considered a spatially constant cell density, with growth rates independent of the local 675 biochemical environment. Future modelling work will represent individual organoids as small, localised 676 regions within the hydrogel where glucose consumption and lactate production occur, and regulate organoid 677 growth. We will use a mathematical homogenisation approach (see e.g. Dalwadi et al. (2018); Dalwadi and 678 King (2020); Sanz-Herrera et al. (2008); Shipley et al. (2009)) to systematically average the behaviour over 679 680 the microscale to obtain a macroscale governing equation for the hydrogel layer with effective glucose consumption, lactate production and organoid growth terms. This in turn will increase our understanding 681 of the relationship between the bioreactor operating parameters and the mean and variation in organoid 682 683 size, ultimately facilitating optimisation of the bioreactor operating conditions to minimise organoid size variation. 684

The mathematical modelling approach developed in this paper provides a framework for establishing 685 how organoid viability can be improved by varying bioreactor operating conditions. The framework has 686 the flexibility to consider different organoid lines, via characterisation of their proliferation and nutrient 687 consumption rates and their tolerance to the presence of waste metabolite. Our work has the potential to 688 improve the quality and reproducibility of bioreactor-expanded organoid output. We intend our theoretical 689 framework to be used to scale-up the production of viable organoids, contributing to overall organoid 690 technology development, and enabling organoids to be exploited as a powerful tool for accelerating drug 691 692 discovery and testing.

# **CONFLICT OF INTEREST STATEMENT**

693 MJE is a co-founder of Cellesce.

# **AUTHOR CONTRIBUTIONS**

694 All authors designed the research; MAE performed the research; all the authors wrote the paper.

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# DATA AVAILABILITY STATEMENT

699 The datasets generated for this study can be found in the GitHub repository https://github.com/ 700 meredithellis/A\_reduced-order\_model\_for\_organoid\_expansion.

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# TABLES

**Table 1.** Definitions of dimensional model parameters, together with typical values. Where no citation is given, parameters are taken from the CPX1 set–up.

parameter	definition	typical value	
$D_{CH}$	diffusivity of glucose in hydrogel	$6.0 \times 10^{-10} \text{m}^2 \text{ s}^{-1}$ (Suhaimi et al., 2015)	
$D_{CM}$	diffusivity of glucose in media	$6.0 \times 10^{-10} \text{m}^2 \text{ s}^{-1}$ (Suhaimi and Das, 2016)	
$D_{WH}$	diffusivity of lactate in hydrogel	$1.2 \times 10^{-9} \text{m}^2 \text{ s}^{-1}$ (Zhou et al., 2008)	
$D_{WM}$	diffusivity of lactate in media	$1.4 \times 10^{-9} \text{m}^2 \text{ s}^{-1}$ (Shipley et al., 2011)	
$c_{-\infty}$	glucose concentration in upstream reservoir	$0.36 \text{ mol m}^{-2}$	
[u]	maximum velocity of media flow	$1 \times 10^{-6} \mathrm{m  s^{-1}}$	
L	length of bioreactor	$9 \times 10^{-2} \mathrm{m}$	
$h_H$	height of hydrogel layer	$1 \times 10^{-3} \mathrm{m}$	
$h_M$	combined height of hydrogel and media	$3 \times 10^{-3} \mathrm{m}$	
$N_0$	initial cell seeding density	$2.7 \times 10^{10} \text{cell m}^{-2}$ to $4 \times 10^{10} \text{cell m}^{-2}$	
p	proliferation rate	$3.9 \times 10^{-6} \mathrm{s}^{-1}$	
$ u_C $	rate of glucose consumption per unit cell density	$9.4 \times 10^{-17} \mathrm{m}^2  \mathrm{cell}^{-1}  \mathrm{s}^{-1}$	

**Table 2.** Timescale groupings of the various physical processes present in the CXP1 bioreactor. We use "x" and "z" to denote "vertical" and "horizontal", respectively. The timescale for each process is the value such that the each dimensionless parameter grouping, defined in Eq. (2.21) as the ratio of the timescale of interest to the timescale of the physical process, is equal to one.

	physical process	timescale
$\mathcal{O}(\text{hour})$	z diffusion glucose in hydrogel	$\frac{\epsilon^2 L^2}{D_{CH}} = 1.5 \times 10^4 \text{s} = 4.2 \text{h}$
	z diffusion glucose in media	$\frac{\epsilon^2 L^2}{D_{CM}} = 1.5 \times 10^4 \text{s} = 4.2 \text{h}$
	z diffusion lactate in hydrogel	$\frac{\epsilon^2 L^2}{D_{WH}} = 7500 \text{s} = 2.1 \text{h}$
	z diffusion lactate in media	$\frac{\epsilon^2 L^2}{D_{WM}} = 6400 \text{s} = 1.8 \text{h}$
$\mathcal{O}(\mathrm{day})$	flow	$\frac{L}{[u]} = 9 \times 10^4 \text{s} = 25 \text{h}$
	glucose consumption	$\frac{1}{\nu_C N_0} = 2.7 \times 10^5 - 4 \times 10^5 \text{s} = 74 - 110 \text{h}$
	lactate production	$\frac{1}{2\nu_C N_0} = 1.3 \times 10^5 - 2.0 \times 10^5 \text{s} = 37 - 55\text{h}$
	cell proliferation	$\frac{1}{p} = 2.6 \times 10^5 \text{s} = 72 \text{h}$
$\mathcal{O}(\text{month})$	x diffusion glucose in hydrogel	$\frac{L^2}{D_{CH}} = 1.4 \times 10^7 \text{s} = 3800 \text{h}$
	x diffusion glucose in media	$\frac{L^2}{D_{CM}} = 1.4 \times 10^7 \text{s} = 3800 \text{h}$
	x diffusion lactate in hydrogel	$\frac{L^2}{D_{WH}} = 6.8 \times 10^6 \text{s} = 1900 \text{h}$
	x diffusion lactate in media	$\frac{L^2}{D_{WM}} = 5.8 \times 10^6 \text{s} = 1600 \text{h}$

**Table 3.** Definitions of non–dimensionalised model parameters with their typical values. For the simulations in this paper, we take  $\rho = 0.27$  unless otherwise stated.

parameter	definition	typical value
$\epsilon$	ratio of vertical to horizontal lengthscales	1/30
$d_{CH}$	ratio of timescale of interest to timescale of diffusion of glucose in hydrogel	$6.4 \times 10^{-3}$
$d_{CM}$	ratio of timescale of interest to timescale of diffusion of glucose in media	$6.4 \times 10^{-3}$
$d_{WH}$	ratio of timescale of interest to timescale of diffusion of lactate in hydrogel	$1.28 \times 10^{-2}$
$d_{WM}$	ratio of timescale of interest to timescale of diffusion of lactate in media	$1.49 \times 10^{-2}$
$\mu$	ratio of timescale of interest to timescale of flow	0.96
ρ	ratio of timescale of interest to that of glucose consumption per cell	0.22 - 0.32
P	ratio of timescale of interest to timescale of cellular proliferation	1/3
$H_H$	ratio of hydrogel height to the combined height of hydrogel and media layers	1/3
$W_{\rm tol}$	dimensionless maximum tolerated lactate concentration	0.7

# FIGURES



**Figure 1.** (Top) Schematic of 'CXP1' bioreactor (Ellis et al., 2019). (Bottom) Two-dimensional crosssection of the bioreactor, with arrows indicating the half-Poiseuille flow profile. Blue is media, yellow is hydrogel, grey is organoid biomass. The glucose concentrations within the media and hydrogel are given by  $c_M$  and  $c_H$ , respectively. Similarly, the lactate concentrations within hydrogel and media are denoted  $w_M$  and  $w_H$ , respectively. (Bottom right) Example of colorectal cancer organoid. Confocal image using 20X objective of Cell Insight Cx7. Organoid stained for nuclear (blue) and cytoskeletal (red) markers for imaging. Scalebar 50µm. Reproduced with permission from Cellesce.



**Figure 2.** Schematic of the boundary conditions for the media (blue) and hydrogel (yellow) layers for eqs. (2.1), (2.2), (2.6) and (2.7). At the media-hydrogel interface, we impose continuity of concentration and flux. At the air–media interface and at the impermeable hashed boundaries, we impose no flux. The black arrows indicate the half-Poiseuille flow profile.



**Figure 3.** Metabolite concentrations at 1, 3, and 7 days into a typical simulation. The horizontal lines at Z = 1/3 represents the media–hydrogel interface. (**Top**) Glucose distribution C(X, Z, T) at (**A**) T = 1, (**B**) T = 3, (**C**) T = 7. (**Bottom**) Lactate distribution W(X, Z, T) at (**D**) T = 1, (**E**) T = 3, (**F**) T = 7. Parameter values: see Table 3.



**Figure 4.** Results showing how the glucose (**A**, **C**, **E**) and lactate (**B**, **D**, **F**) concentrations change over time during a typical simulation. (**A**, **B**) Results from *Z*-averaged full model. (**C**, **D**) Longwave approximation. (**E**, **F**) Sublimit of longwave approximation, where the upper left and lower right regions are the *replenished* and *unreplenished* regions, respectively. Parameter values: see Table 3.



**Figure 5.** Comparison of outputs from the different mathematical models and their evolution in time: (A) minimum glucose concentration,  $C_{\min}(T)$ ; (B) maximum lactate concentration,  $W_{\max}(T)$ ; (C) spatial position of maximum lactate concentration,  $X_{\max}(T)$  s.t.  $W(X_{\max},T) = W_{\max}(T)$ ; (D) lactate concentration at outlet of bioreactor, W(X = 1, T). The red points represent the values predicted in the hydrogel region of the full 2D model. Parameter values: see Table 3.



**Figure 6.** Evolution of glucose (**left grid**) and lactate (**right grid**) concentration profiles over the duration of a typical experiment for different organoid lines under the same operating conditions. The rates of cell proliferation rates and glucose consumption per cell are: (**A**), (**E**) organoid line (*i*), p = 1/6,  $\rho = 0.027$ ; (**B**), (**F**) organoid line (*ii*), p = 1/6,  $\rho = 2.7$ ; (**C**), (**G**) organoid line (*iii*), p = 1,  $\rho = 0.027$ ; (**D**), (**H**) organoid line (*iv*), p = 1,  $\rho = 2.7$ . The other parameters used are given in Table 3.



**Figure 7.** Comparison of (A) glucose conversion Q, Eq. (3.3), (B) maximum lactate concentration  $W_{\text{max}}(T)$ , Eq. (3.2), where the red line represents the maximum tolerated lactate concentration,  $W = W_{\text{tol}}$ , and (C) proportion of domain which is uninhabitable at time T,  $P_U$ , Eq. (3.4), for different organoid lines cultured within the bioreactor under the same operating conditions. The proliferation rates and rate of glucose consumption per cell for each organoid line are: (i) p = 1/6,  $\rho = 0.027$ , (ii) p = 1/6,  $\rho = 2.7$ , (iii) p = 1,  $\rho = 0.027$ , (iv) p = 1,  $\rho = 2.7$ , and (v) p = 1/3,  $\rho = 0.27$ . The other parameters used are given in Table 3. The line styles correspond to rate of cellular proliferation: solid, P = 1/6; dashed, P = 1/3; and dotted, P = 1. The line colours correspond to rate of glucose consumption per cell density: purple,  $\rho = 0.027$ ; blue,  $\rho = 2.7$ ; and green,  $\rho = 2.7$ .



**Figure 8.** Results for a specific organoid line within the CXP1 bioreactor showing the evolution of: (A) glucose conversion Q, Eq. (3.3), (B) maximum lactate concentration  $W_{\max}(T)$ , Eq. (3.2), where the red line respresents the maximum tolerated lactate concentration,  $W = W_{tol}$ , and (C) proportion of domain which is uninhabitable at time T,  $P_U$ , Eq. (3.4), against time for five different flow rates. For  $[u] = 10^{-5} \text{m s}^{-1}$ , the value of  $P_U$  is zero. The peak flow velocities  $[u] \in \{10^{-7}, 5 \times 10^{-7}, 10^{-6}, 2 \times 10^{-6}, 10^{-5}\} \text{m s}^{-1}$  used correspond to the dimensionless flow velocity parameter  $\mu \in \{0.096, 0.48, 0.96, 1.92, 9.6\}$ , respectively. Remaining parameter values: see Table 3.



Figure 9. (A) Glucose conversion Q, Eq. (3.3), at time T = 7 and (B) turn-off time  $T_{\text{off}}$  (the time when intolerable lactate levels first experienced) for the CXP1 bioreactor varying with flow rate, for a given organoid line. Peak flow velocities  $[u] \in [10^{-7}, 10^{-5}] \text{ m s}^{-1}$  correspond to dimensionless flow rate,  $\mu$ , in the range  $\mu \in [0.096, 9.6]$  and the other parameter values are given in Table 3.