

# The spectral function of a first order system

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Will be studying the spectral problem for a system

$$Av = \lambda v$$

where  $A$  is an elliptic self-adjoint first order  $m \times m$  matrix pseudo-differential operator acting on columns of  $m$  complex-valued half-densities  $v$  over a connected compact  $n$ -dimensional manifold  $M$  without boundary. The operator is not necessarily semi-bounded.

Principal symbol  $A_1(x, \xi)$  is matrix-valued function on  $T^*M \setminus \{0\}$ .

The eigenvalues of the principal symbol are denoted  $h^{(j)}(x, \xi)$ . These  $h^{(j)}(x, \xi)$  are assumed to be simple.

## Objects of study

**Object 1.** Wave group (or unitary group)

$$U(t) := e^{-itA} = \sum_k e^{-it\lambda_k} v_k(x) \int_M [v_k(y)]^*(\cdot) dy$$

where  $\lambda_k$  and  $v_k$  are eigenvalues and eigenfunctions of operator  $A$ .

**Object 2.** Spectral function

$$e(\lambda, x, x) := \sum_{0 < \lambda_k < \lambda} \|v_k(x)\|^2.$$

**Object 3.** Counting function

$$N(\lambda) := \sum_{0 < \lambda_k < \lambda} 1 = \int_M e(\lambda, x, x) dx .$$

## Objectives

**Objective 1:** construct wave group explicitly in terms of oscillatory integrals, modulo integral operator with smooth kernel.

**Objectives 2 and 3:** derive, under appropriate assumptions on Hamiltonian trajectories, two-term asymptotics

$$e(\lambda, x, x) = a(x) \lambda^n + b(x) \lambda^{n-1} + o(\lambda^{n-1}),$$

$$N(\lambda) = a\lambda^n + b\lambda^{n-1} + o(\lambda^{n-1}),$$

as  $\lambda \rightarrow +\infty$ . Here I expect the real constants  $a$ ,  $b$  and real densities  $a(x)$ ,  $b(x)$  to be related as

$$a = \int_M a(x) dx, \quad b = \int_M b(x) dx.$$

## **Warning: doing microlocal analysis for systems is not easy**

Main contributors to the spectral theory of systems:

- 1 V.Ivrii, 1980, Soviet Math. Doklady.
- 2 V.Ivrii, 1982, Funct. Anal. Appl.
- 3 V.Ivrii, book, 1984, Springer Lecture Notes.
- 4 Yu.Safarov, DSc thesis, 1989, Steklov Mathematical Institute.
- 5 V.Ivrii, book, 1998, approx 800 pages, Springer.
- 6 V.Ivrii, future book. 2012? Approx 3000 pages?
- 7 O.Chervova, R.J.Downes and D.Vassiliev, *The spectral function of a first order system* (in preparation, approx 50 pages).

## Safarov's formula for the second asymptotic coefficient

$$b(x) = -n \sum_{j=1}^{m^+} \int_{h^{(j)}(x, \xi) < 1} \left( [v^{(j)}]^* A_{\text{sub}} v^{(j)} - \frac{i}{2} \{ [v^{(j)}]^*, A_1 - h^{(j)}, v^{(j)} \} \right) (x, \xi) d\bar{\xi}$$

where  $v^{(j)}(x, \xi)$  are eigenvectors of principal symbol,

$$\{P, R\} := P_{x^\alpha} R_{\xi_\alpha} - P_{\xi_\alpha} R_{x^\alpha}$$

is the Poisson bracket on matrix-functions and

$$\{P, Q, R\} := P_{x^\alpha} Q R_{\xi_\alpha} - P_{\xi_\alpha} Q R_{x^\alpha}$$

is its further generalisation.

## Testing $U(m)$ invariance of second asymptotic coefficient

Transform the operator as

$$A \mapsto RAR^*$$

where

$$R : M \rightarrow U(m)$$

is an arbitrary smooth unitary matrix-function. The coefficient  $b(x)$  should be unitarily invariant.

There's something missing in the integrand of Safarov's formula:

$$[v^{(j)}]^* A_{\text{sub}} v^{(j)} - \frac{i}{2} \{ [v^{(j)}]^*, A_1 - h^{(j)}, v^{(j)} \} + ? .$$

## The U(1) connection

Each eigenvector  $v^{(j)}(x, \xi)$  of the principal symbol is defined modulo a gauge transformation

$$v^{(j)} \mapsto e^{i\phi^{(j)}} v^{(j)}$$

where

$$\phi^{(j)} : T^*M \setminus \{0\} \rightarrow \mathbb{R}$$

is an arbitrary smooth function. This gives a U(1) connection characterised by a  $2n$ -component real covector potential  $(P_\alpha, Q^\beta)$  where

$$P_\alpha := i[v^{(j)}]^* v_{x^\alpha}^{(j)}, \quad Q^\beta := i[v^{(j)}]^* v_{\xi_\beta}^{(j)}.$$

The corresponding covariant derivative is

$$\nabla_\alpha := \partial/\partial x^\alpha + iP_\alpha, \quad \nabla^\beta := \partial/\partial \xi_\beta + iQ^\beta.$$

This definition ensures that  $\nabla v^{(j)}$  is orthogonal to  $v^{(j)}$ .



The curvature of our U(1) connection is

$$R := i \begin{pmatrix} \nabla_\alpha \nabla_\beta - \nabla_\beta \nabla_\alpha & \nabla_\alpha \nabla^\delta - \nabla^\delta \nabla_\alpha \\ \nabla^\gamma \nabla_\beta - \nabla_\beta \nabla^\gamma & \nabla^\gamma \nabla^\delta - \nabla^\delta \nabla^\gamma \end{pmatrix}.$$

This curvature is the exterior derivative of the covector potential.

The scalar curvature of our U(1) connection is

$$\nabla_\alpha \nabla^\alpha - \nabla^\alpha \nabla_\alpha = -i \{ [v^{(j)}]^*, v^{(j)} \}.$$

Recall: curly brackets denote Poisson bracket on matrix-functions

$$\{P, R\} := P_{x^\alpha} R_{\xi_\alpha} - P_{\xi_\alpha} R_{x^\alpha}.$$

## How curvature of U(1) connection manifests itself in microlocal analysis

$$U(t) \stackrel{\text{mod } C^\infty}{=} \sum_j U^{(j)}(t)$$

where phase function of each oscillatory integral  $U^{(j)}(t)$  is associated with corresponding Hamiltonian  $h^{(j)}(x, \xi)$ . New result:

$$\text{tr}[U^{(j)}(0)]_{\text{sub}} = -i\{[v^{(j)}]^*, v^{(j)}\}.$$

Note:  $\sum_j \{[v^{(j)}]^*, v^{(j)}\} = 0$ .

Correct integrand in formula for second asymptotic coefficient:

$$[v^{(j)}]^* A_{\text{sub}} v^{(j)} - \frac{i}{2} \{[v^{(j)}]^*, A_1 - h^{(j)}, v^{(j)}\} + \frac{i}{n-1} h^{(j)} \{[v^{(j)}]^*, v^{(j)}\}.$$

**Additional assumption 1:**  $m = 2$  and  $\text{tr } A_1 = 0$

In this case principal symbol has only two eigenvalues,  $h^\pm(x, \xi)$ , and two eigenvectors,  $v^\pm(x, \xi)$ . These are related as

$$h^- = -h^+, \quad v^- = \overline{\epsilon v^+}.$$

where

$$\epsilon := \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}.$$

The antilinear map  $v \mapsto C(v) := \epsilon \bar{v}$  is called *charge conjugation*.

Integrand in formula for second asymptotic coefficient simplifies:

$$[v^+]^* A_{\text{sub}} v^+ + \frac{n}{n-1} i h^+ \{ [v^+]^*, v^+ \}.$$

## **Additional assumption 2: the operator $A$ is differential**

In this case dimension of manifold can only be  $n = 2$  or  $n = 3$ .

The subprincipal symbol  $A_{\text{sub}}$  does not depend on the dual variable  $\xi$  (momentum) and is a function of  $x$  (position) only.

We acquire a geometric object, the metric. The determinant of the principal symbol is a negative definite quadratic form

$$\det A_1(x, \xi) = -g^{\alpha\beta} \xi_\alpha \xi_\beta$$

and coefficients  $g^{\alpha\beta}(x)$ ,  $\alpha, \beta = 1, \dots, n$ , can be interpreted as components of a (contravariant) Riemannian metric. Our Hamiltonian (positive eigenvalue of principal symbol) takes the form

$$h^+(x, \xi) = \sqrt{g^{\alpha\beta}(x) \xi_\alpha \xi_\beta}$$

and Hamiltonian trajectories become geodesics.

### **Additional assumption 3: $n = 3$**

In this case the manifold  $M$  is bound to be parallelizable (and, hence, orientable). Established by examining principal symbol.

Furthermore, we acquire a topological invariant

$$\mathbf{c} := -\frac{i}{2} \sqrt{\det g_{\alpha\beta}} \operatorname{tr} \left( (A_1)_{\xi_1} (A_1)_{\xi_2} (A_1)_{\xi_3} \right).$$

The number  $\mathbf{c}$  can take only two values,  $+1$  or  $-1$ . It describes the orientation of the principal symbol  $A_1(x, \xi)$  relative to the chosen orientation of local coordinates.

Most importantly, we acquire a new differential geometric object, namely, a *teleparallel connection*.

## The teleparallel connection

Define an affine connection as follows. Suppose we have a covector  $\eta$  “based” at the point  $y \in M$  and we want to construct a “parallel” covector  $\xi$  “based” at the point  $x \in M$ . This is done by solving the linear system of equations

$$A_1(x, \xi) = A_1(y, \eta).$$

The teleparallel connection coefficients  $\Gamma^\alpha_{\beta\gamma}(x)$  can be written down explicitly in terms of the principal symbol and this allows us to define yet another geometric object — the *torsion tensor*

$$T^\alpha_{\beta\gamma} := \Gamma^\alpha_{\beta\gamma} - \Gamma^\alpha_{\gamma\beta}.$$

The teleparallel connection has zero curvature and nonzero torsion. It is the opposite of the Levi-Civita connection.

More convenient to work with  $T^*\alpha_\beta := \frac{1}{2} T^{\alpha\gamma\delta} \varepsilon_{\gamma\delta\beta} \sqrt{\det g_{\mu\nu}}$ .

*Lemma 1* The scalar curvature of the U(1) connection is expressed via torsion of the teleparallel connection as

$$-i\{[v^+]^*, v^+\}(x, \xi) = \frac{\mathbf{c}}{2} \frac{\overset{*}{T}{}^{\alpha\beta}(x) \xi_\alpha \xi_\beta}{(g^{\mu\nu}(x) \xi_\mu \xi_\nu)^{3/2}}.$$

Lemma 1 allows us to evaluate the integral in the formula for the second asymptotic coefficient of the spectral function, giving

$$b(x) = \frac{1}{8\pi^2} \left( (\mathbf{c} \operatorname{tr} \overset{*}{T} - 2 \operatorname{tr} A_{\text{sub}}) \sqrt{\det g_{\alpha\beta}} \right) (x).$$

Note that the two traces appearing in this formula have a different meaning:  $\operatorname{tr} \overset{*}{T}$  is the trace of a  $3 \times 3$  tensor whereas  $\operatorname{tr} A_{\text{sub}}$  is the trace of a  $2 \times 2$  matrix.

## The massless Dirac operator

Imagine a single neutrino living in a compact 3-dimensional universe. I want to find the spectrum of energy levels this neutrino can occupy. This leads to an eigenvalue problem

$$Av = \lambda v$$

where  $A$  is the massless Dirac operator, a particular  $2 \times 2$  matrix first order differential operator with trace-free principal symbol.

Note: the explicit formula for the massless Dirac operator in curved space is very complicated.

Now suppose I am a spectral analyst and I want to know whether my differential operator  $A$  is a massless Dirac operator, without having to learn the differential geometry of spinors.



*Theorem 1* An operator is a massless Dirac operator iff

- a) subprincipal symbol is proportional to the identity matrix, and
- b) second asymptotic coefficient of the spectral function is zero.

Note: the conditions of Theorem 1 are invariant under transformations of the operator  $A \mapsto RAR^*$  where  $R : M \rightarrow \text{SU}(2)$ .

Note: Theorem 1 does not feel the topological differences between different parallelizations of the manifold  $M$ .

Bottom line: the basic differential geometry of spinors is encoded within the microlocal analysis of PDEs.

**More radical approach:  
poor man's way of studying nonlinear PDEs**

Working on a compact 3-manifold, consider eigenvalue problem

$$Av = \lambda wv$$

where  $A$  is a  $2 \times 2$  matrix first order partial differential operator with  $\text{tr } A_1 = 0$  and  $A_{\text{sub}} = 0$ , and  $w : M \rightarrow (0, +\infty)$  is a scalar weight function. Note that this eigenvalue problem is invariant under the gauge transformation

$$A \mapsto fAf, \quad w \mapsto f^2w$$

where  $f : M \rightarrow (0, +\infty)$ .

View coefficients as dynamical variables and consider the map

**principal symbol of operator and scalar weight  $\mapsto$   
second asymptotic coefficient of counting function.**

*Theorem 2 (to be published in a separate paper)* The map

**principal symbol of operator and scalar weight  $\mapsto$   
second asymptotic coefficient of counting function.**

is equivalent to the map

**metric and spinor field  $\mapsto$  static massless Dirac action.**

### *Explanation*

Metric is the determinant of the principal symbol.

A  $2 \times 2$  trace-free Hermitian matrix is not fully defined by its determinant. The remaining degrees of freedom are called “spinor”.

Gauge invariance of the eigenvalue problem  $Av = \lambda wv$  manifests itself as the conformal invariance of the massless Dirac action.

## Four fundamental equations of theoretical physics

- 1 Maxwell's equations. Describe electromagnetism and photons.
- 2 Dirac equation. Describes electrons and positrons.
- 3 Massless Dirac equation. Describes\* neutrinos and antineutrinos.
- 4 Linearized Einstein field equations of general relativity. Describe gravity.

All four contain the same physical constant — the speed of light.

\*OK, I know that neutrinos actually have a small mass.

## Accepted explanation: theory of relativity

God is a geometer. He created a 4-dimensional world parametrized by coordinates  $x^0, x^1, x^2, x^3$  (here  $x^0$  is time), in which distances are measured in a funny way:

$$\text{distance}^2 = -c^2(dx^0)^2 + (dx^1)^2 + (dx^2)^2 + (dx^3)^2$$

where  $c$  is the speed of light.

Without the term  $-c^2(dx^0)^2$  this would be Pythagoras' theorem. Funny way of measuring distances is called *Minkowski metric*.

Having decided to use the Minkowski metric, God then wrote down the main equations of theoretical physics using **only geometric constructions**, i.e. using concepts such as connection, curvature etc. This way all equations have the same physical constant — the speed of light — encoded in them.

## Alternative explanation

God is an analyst. He created a 3-dimensional Euclidean world, then added (absolute) time and wrote down one system of non-linear PDEs which describes phenomena in this world. This system of PDEs has different solutions which we interpret as electromagnetism, gravity, electrons, neutrinos etc. The reason the same physical constant — the speed of light — manifests itself in all physical phenomena is because we are looking at different solutions of the **same** system of PDEs.

I cannot write down the unifying system of field equations but I can extract some geometric constructs of theoretical physics from systems of PDEs, simply by performing microlocal analysis.

Action plan: spend next 100 years studying systems of nonlinear hyperbolic PDEs, hoping to find soliton-type solutions.