

# Modelling the neutrino in terms of Cosserat elasticity

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## Subject of talk

Weyl equation (massless Dirac equation):

$$\sigma^\alpha_{\dot{a}b} \nabla_\alpha \xi^b = 0.$$

Here

- $\sigma^\alpha$ ,  $\alpha = 0, 1, 2, 3$ , are Pauli matrices,
- $\nabla_\alpha$  are covariant derivatives with respect to local coordinates  $x^\alpha$ ,  $\alpha = 0, 1, 2, 3$ , and
- $\xi^b$ ,  $b = 1, 2$ , is the unknown spinor field.

Will construct an alternative model without spinors, Pauli matrices or covariant derivatives.

## Describing a 3D deformable medium

(a) Classical elasticity: displacements only.

(b) Cosserat elasticity: displacements and rotations. See

E. Cosserat and F. Cosserat, *Théorie des Corps Déformables*, 1909. Available from Amazon.

I will assume that there are no displacements, only rotations.

To describe rotations of material points mathematically I attach to each geometric point a *coframe*.

A coframe  $\vartheta$  in 3D is a triplet  $\vartheta^j$ ,  $j = 1, 2, 3$ , of orthonormal covector fields. Each covector field  $\vartheta^j$  has a hidden tensor index:  $\vartheta^j = \vartheta^j_\alpha$ ,  $\alpha = 1, 2, 3$ .

Same in plain English: a coframe is a field of orthonormal bases.

**NB.** Coframe lives separately from local coordinates (not aligned with coordinate lines).

The coframe  $\vartheta$  is an unknown quantity (dynamical variable).

The other dynamical variable is a density  $\rho$ .

Choose potential energy from the condition of conformal invariance. Explicit formula:

$$P(x^0) = \int \|T^{\text{ax}}\|^2 \rho dx^1 dx^2 dx^3,$$

$$T^{\text{ax}} = \frac{1}{3}(\vartheta^1 \wedge d\vartheta^1 + \vartheta^2 \wedge d\vartheta^2 + \vartheta^3 \wedge d\vartheta^3).$$

Standard kinetic energy

$$K(x^0) = \int \|\dot{\vartheta}\|^2 \rho dx^1 dx^2 dx^3,$$

$$\dot{\vartheta} = \frac{1}{3}(\vartheta^1 \wedge \partial_0 \vartheta^1 + \vartheta^2 \wedge \partial_0 \vartheta^2 + \vartheta^3 \wedge \partial_0 \vartheta^3).$$

My action (variational functional)

$$S = \int (P(x^0) - K(x^0)) dx^0.$$

## **Difference with existing models (teleparallelism)**

1. I assume metric to be fixed, i.e. I do not vary metric.
2. My Lagrangian has never been considered.

## Solving Euler–Lagrange equations

Vary coframe and density to get Euler–Lagrange equations. Too complicated!

Switch to spinors:

coframe  $\vartheta$  and density  $\rho > 0$



nonvanishing spinor field  $\xi$  modulo sign

My Lagrangian density  $L(\xi)$  is a rational function of  $\xi$ ,  $\bar{\xi}$  and partial derivatives of  $\xi$ ,  $\bar{\xi}$ .

## Stationary solutions

Look first for stationary solutions

$$\xi(x^0, x^1, x^2, x^3) = e^{-i\varepsilon x^0} \eta(x^1, x^2, x^3), \quad \varepsilon \neq 0.$$

**Theorem 1** *In the stationary case my Euler–Lagrange equation is equivalent to a pair of Weyl equations.*

**Proof** Turns out my Lagrangian factorises as

$$L(\eta) = \frac{L_{\text{Weyl}}^+(\eta) L_{\text{Weyl}}^-(\eta)}{L_{\text{Weyl}}^+(\eta) - L_{\text{Weyl}}^-(\eta)}.$$

Result follows from factorisation.  $\square$

See <http://arxiv.org/abs/0902.1268>.



*Question 1.* Can I handle the non-stationary case, i.e. when time dependence is arbitrary? Yes, by means of perturbation theory. I look at perturbations of plane wave solutions. Again I get a pair of Weyl equations.

*Question 2.* Can I make my model relativistically invariant? Yes, by viewing  $(1 + 3)$ -dimensional spacetime as a Cosserat continuum. At perturbative level there is no difference between the nonrelativistic and relativistic models.

*Question 3.* Can I incorporate mass and external electromagnetic field into my model? Yes, by means of a Kaluza–Klein extension, i.e. by adding an extra coordinate  $x^4$ . See <http://arxiv.org/abs/0812.3948>.

**Theorem 2** *In the special case with no dependence on  $x^3$  (i.e. for electron in dimension  $1 + 2$ ) the massive version of my model is equivalent to the massive Dirac equation.*

## Summary

New mathematical model for fermions.

- Spacetime viewed as Cosserat continuum.
- Lagrangian chosen from condition of conformal invariance.
- Mass and electromagnetic field incorporated via Kaluza–Klein extension.