Periodic, almost periodic, and not periodic at all problems

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Periodic differential operators

Periodic Schrödinger operator

$$H = -\Delta + V$$

with smooth periodic potential V = V(x), $x \in \mathbb{R}^d$. This means $V(x + \gamma) = V(x)$ for all $\gamma \in \Gamma$.

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If d = 1, then the number of gaps is almost always infinite.

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Bethe-Sommerfeld conjecture

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Bethe-Sommerfeld conjecture





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If $d \ge 2$, the number of gaps is always finite.

Proved: d = 2: V.Popov, M.Skriganov (1981)

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We can also consider the magnetic Schrödinger operator

 $H = (i\nabla + A)^2 + V$, where A = A(x) is a periodic vector-function with the same lattice of periods Γ . Then the Bethe-Sommerfeld conjecture was proved only for d = 2 (A. Mohamed, 1997).

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Important tool when working with periodic problems: Floquet-Bloch decomposition.





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Important tool when working with periodic problems: Floquet-Bloch decomposition.





$$H=\int_{\oplus}H(k)dk,$$

where $H(k) := (i\nabla + k)^2 + V$ acts in $L^2(\mathbb{R}^d/\Gamma)$, $k \in \mathbb{R}^d/\Gamma'$, and Γ' is the (analytical) dual to Γ . This means that

$$\sigma(H) = \bigcup_{k \in \mathbb{R}^d / \Gamma'} \sigma(H(k)).$$

The spectrum of H(k) consists of eigenvalues:

$$\sigma(H(k)) = \{\lambda_1(k) \leq \lambda_2(k) \leq \dots\}.$$

Now we can define

$$\ell_j := \cup_{k \in \mathbb{R}^d / \Gamma'} \lambda_j(k)$$

as the *n*-th spectral band, so that $\sigma(H) = \bigcup_j \ell_j$. Then for each λ we can define two functions:

$$m(\lambda) = \#\{j : \lambda \in \ell_j\}$$

(the multiplicity of overlapping) and

$$\zeta(\lambda) = \max_{j} \max\{t : [\lambda - t, \lambda + t] \subset \ell_j\}.$$

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(the overlapping function)

Theorem. (A.Sobolev,LP, 2001)

Let d = 2, 3, 4. Then for sufficiently large λ we have:

Dimension	$m(\lambda) \gg$	$\zeta(\lambda) \gg$
2	$\lambda^{rac{1}{4}}$	$\lambda^{rac{1}{4}}$
3	$\lambda^{rac{1}{2}}$	1
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Unfortunately, the method does not work for $d \ge 5!$

If we want to prove the conjecture, we need to study the eigenvalues of H(k). There are two types of eigenvalues of these operators: stable (corresponding to perturbations of simple eigenvalues, lying not too close to other eigenvalues) and unstable (corresponding to perturbations of a cluster of eigenvalues lying close together). It is relatively straightforward to compute stable eigenvalues with high precision. Unstable eigenvalues cause the main problem.

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Moreover, for large λ we have $m(\lambda) \geq 1$ and $\zeta(\lambda) \geq \lambda^{\frac{1-\alpha}{2}}$.

Theorem. (G.Barbatis, LP, 2009)

Let $d \ge 2$. Then the Bethe-Sommerfeld conjecture holds for operators $H = (-\Delta)^m + q$ with periodic pseudo-differential operators q of order smaller than 2m - 1.

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Let $d \ge 2$. Then the Bethe-Sommerfeld conjecture holds for operators $H = (-\Delta)^m + q$ with periodic pseudo-differential operators q of order smaller than 2m. In particular, this conjecture holds for periodic magnetic Schrödinger operators. Another line of research: the asymptotic behaviour of the (integrated) density of states. The density of states of $H = -\Delta + V$ can be defined by the formula

$$N(\lambda) = N(\lambda; H) := \lim_{L \to \infty} \frac{N(\lambda; H_D^{(L)})}{(2L)^d}.$$

Here, $H_D^{(L)}$ is the restriction of *H* to the cube $[-L, L]^d$ with the Dirichlet boundary conditions, and $N(\lambda; H_D^{(L)}) = \#\{\lambda_j(H_D^{(L)}) \le \lambda\}$ is the counting function of the discrete spectrum of $H_D^{(L)}$. More convenient definition:

$$N(\lambda) = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d/\Gamma'} N(\lambda, H(k)) dk$$

If we put $H_0 = -\Delta$, for positive λ we have $N(\lambda; H_0) = C_d \lambda^{d/2}$,

where

$$C_d = \frac{w_d}{(2\pi)^d}$$

and

$$w_d = \frac{\pi^{d/2}}{\Gamma(1+d/2)}$$

is a volume of the unit ball in \mathbb{R}^d .

There is a long-standing conjecture that the density of states of *H* enjoys the following asymptotic behaviour as $\lambda \to \infty$:

$$N(\lambda) \sim \lambda^{d/2} \Big(C_d + \sum_{j=1}^{\infty} e_j \lambda^{-j} \Big),$$
 (1)

meaning that for each $K \in \mathbb{N}$ one has

$$N(\lambda) = \lambda^{d/2} \left(C_d + \sum_{j=1}^{K} e_j \lambda^{-j} \right) + R_K(\lambda)$$
(2)

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with $R_{\kappa}(\lambda) = o(\lambda^{\frac{d}{2}-\kappa}).$

Formula (1) was proved in the case d = 1 by Shenk-Shubin (1987).

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If d = 2, formula (2) was proved by A.Sobolev (2005) with K = 2 (three terms) and $R(\lambda) = O(\lambda^{-6/5})$. Yu.Karpeshina (2000) has shown that formula (2) is valid with K = 1 (two terms) and $R(\lambda) = O(\lambda^{-\frac{1}{105}})$ when d = 3 and $R(\lambda) = O(\lambda^{\frac{d-3}{2}} \ln \lambda)$ when d > 3.

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Theorem. (R.Shterenberg,LP, 2008–2010)

Formula (1) holds in all dimensions.

Almost-periodic problems



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Let the potential V be quasi-periodic, i.e.

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where Θ is a finite set.

Almost nothing is known about the spectrum if $d \ge 2!$

We want to study the density of states of quasi-periodic operators. First, we need to impose additional condition: let $Z(\Theta)$ be the collection of all linear combination of elements from Θ with integer coefficients. Let $\theta_1, \ldots, \theta_d \in Z(\Theta)$. Then either $\{\theta_j\}$ are linearly independent, or $\sum_{j=1}^d n_j \theta_j = 0$, where $n_j \in \mathbb{Z}$ and not all n_j are zeros.

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Suppose now that V is almost-periodic, i.e. is a uniform limit of quasi-periodic functions. Then formula (1) still holds, if we impose additional diophantine-type conditions on V. These conditions are too complicated to write them here!

Periodic problems on manifolds

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Random problems

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