### Many Particle Systems

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1st May

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#### Introduction

#### The Black Box

Master Equation Second Quantization

#### SDE

Its Form Solving the SDE

#### Results

Numerically

#### The Future

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#### The Classical Problem

- Classical rate equations fail in the limit of relatively small finite particle populations.
- Our investigations start with the reaction  $A + A \rightarrow A$ .
- ► We only have one species of particle whose population after some time T<sub>f</sub> will reach 1.

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#### The Rate Equation

#### For $A + A \rightarrow A$ we use the Smoluchowski equation.

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#### The Rate Equation

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#### The Rate Equation

- For  $A + A \rightarrow A$  we use the Smoluchowski equation.
- $\frac{dc}{dt} = -\kappa c^2$
- c is the mean particle population and  $\kappa$  is the rate coefficient
- In the "thermodynamic limit" i.e where populations are large and fluctuations are relatively small; (φ<sup>2</sup>) ≈ (φ)<sup>2</sup>, c = (φ).

Master Equation Second Quantization

#### The Master Equation

First we must write a master equation for this reaction.

 <sup>dP(N,t)</sup>/<sub>dt</sub> = <sup>κ</sup>/<sub>2V</sub>[(N+1)NP(N+1,t) − N(N-1)P(N,t)]

Master Equation Second Quantization

### Second Quantization

 We create a Hilbert Space, with the usual notion of annihilation and creation operators.

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Master Equation Second Quantization

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Master Equation Second Quantization

# Second Quantization

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- We define a wave equation to describe our system for any given time and write down the evolution operator, the Hamiltonian.
- We can then formulate an observable and its expectation value, this is the population and its expected value.
- Finally using path integral formalism we are able to extract a Stochastic Differential Equation (SDE).

Its Form Solving the SDE

### Our SDE from the black box

$$\blacktriangleright \partial_t \bar{\Phi}(t) = -\bar{\kappa} \bar{\Phi}^2(t) + i\eta \sqrt{2\bar{\kappa}} \bar{\Phi}$$

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Its Form Solving the SDE

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$$\Phi(t) \in \mathbb{C}$$
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Its Form Solving the SDE

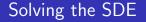
#### Our SDE from the black box

$$\blacktriangleright \partial_t \bar{\Phi}(t) = -\bar{\kappa} \bar{\Phi}^2(t) + i\eta \sqrt{2\bar{\kappa}} \bar{\Phi}$$

- The 
   η is a "white noise".
- ▶  $\Phi(t) \in \mathbb{C}$  .
- Where we are interested in  $\langle \Phi(t) \rangle$ .

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Its Form Solving the SDE

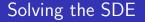


There are two main methods;

► Analytically: Here we use Itô calculus to construct a solution.

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Its Form Solving the SDE



There are two main methods;

- ► Analytically: Here we use Itô calculus to construct a solution.
- Numerically: Using a sequence of random numbers to create a Wiener process and implementing it in an iterative method.
   We then average over a large number of runs.

Its Form Solving the SDE

### Analytically

Using Itô calculus gives us

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Its Form Solving the SDE

# Analytically

 Using Itô calculus gives us
 Φ(t) = Φ<sub>0</sub>exp(κt+i√2κW(t)) 1+κΦ<sub>0</sub> ∫<sub>0</sub><sup>t</sup> exp(κs+i√2κW(s))ds



Its Form Solving the SDE

# Numerically

A Wiener process can be defined as such, W(0) = 0 w.p.1, E(W(t)) = 0 and finally Var(W(t) − W(s)) = t − s

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Its Form Solving the SDE

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- We can construct something very similar by dividing our time interval into equal length sub intervals. <sup>t<sub>i</sub></sup>/<sub>N</sub>

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- We then generate a list of gaussian random numbers,  $\{X_1, X_2, ..., X_N\}$  with  $\mu = 0$ .
- ► Then we sum these numbers such that  $S_N(t_n^{(n)}) = (X_1 + X_2 + ... + X_n)\sqrt{\Delta t}$

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- We then generate a list of gaussian random numbers,  $\{X_1, X_2, ..., X_N\}$  with  $\mu = 0$ .
- ► Then we sum these numbers such that  $S_N(t_n^{(n)}) = (X_1 + X_2 + ... + X_n)\sqrt{\Delta t}$
- ▶ It has been shown that  $Var(S_N(t) S_N(s)) \rightarrow t s$  as  $N \rightarrow \infty$

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Its Form Solving the SDE

#### Iterative Method

#### > This can then be used to solve the following form of our SDE

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Its Form Solving the SDE

#### Iterative Method

- This can then be used to solve the following form of our SDE
- $d\Phi(t) = -\kappa \Phi^2(t) dt + i \sqrt{(2\kappa)} \Phi(t) dW$
- Where  $dW = \eta dt$ .

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# Single Projection of $\operatorname{Re}(\Phi(t))$

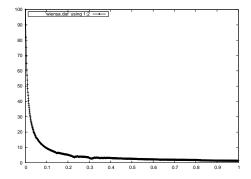


Figure: The real part of a single realisation of the solution to the SDE

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# Single Projection of $Im(\Phi(t))$

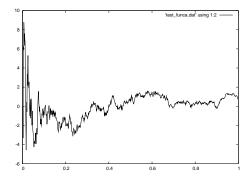


Figure: The real part of a single realisation of the solution to the SDE

Numerically

# $\operatorname{Re}(\langle \Phi(t) \rangle)$

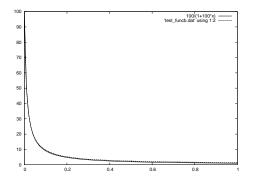


Figure: The real part of a single realisation of the solution to the SDE

Numerically

# $\operatorname{Im}(\langle \Phi(t) \rangle)$

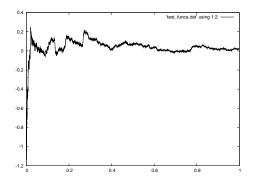


Figure: The real part of a single realisation of the solution to the  $SDE = -9 \circ c$ 



What I would like to move onto next;

 To calculate a reaction where a single "super" dust particle forms.

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What I would like to move onto next;

- To calculate a reaction where a single "super" dust particle forms.
- ▶ Is it possible to show that  $Im(\langle \Phi(t) \rangle) = 0$  for all time?